A History of defining jets

Grégory Soyez

CERN

In collaboration with Gavin Salam, Matteo Cacciari and Juan Rojo

 V^e rencontre SPP/IPhT — CEA/Saclay — November 17 2009

Plan

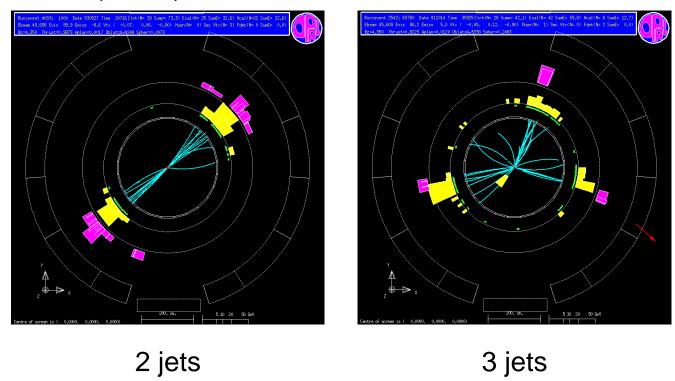
- The past/present: Jet algorithms and jet definitions
 - basic ideas: why jets? recombinations and cones
 - failures of the 20th-century cone algorithms
 - new algorithms in the 21st-century without the failures

- The present/future: how to better use the tools we have?
 - UE sensitivity
 - keywords: jet areas, subtraction, filtering
 - boosted taggers (Higgs, top)
 - keywords: same as above + subjets

Motivation: why jets

Collinear divergence ⇒ QCD produces "jetty" showers

Example: LEP (OPAL) events



"Jets" \equiv bunch of collimated particles \cong hard partons

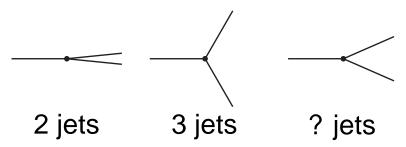
Motivation: why jets

Collinear divergence ⇒ QCD produces "jetty" showers

"Jets" \equiv bunch of collimated particles \cong hard partons

BUT

- a "parton" is an ambiguous concept (NLO)
- "collinear" has some arbitraryness



Motivation: why jets

Collinear divergence ⇒ QCD produces "jetty" showers

"Jets" \equiv bunch of collimated particles \cong hard partons

In practice: use of a jet definition

jet
particles $\{p_i\}$ definition

Jet algorithm: the recipe (insufficient!)
Jet definition: algorithm + the parameters

20th century jet algorithms

Recombination:

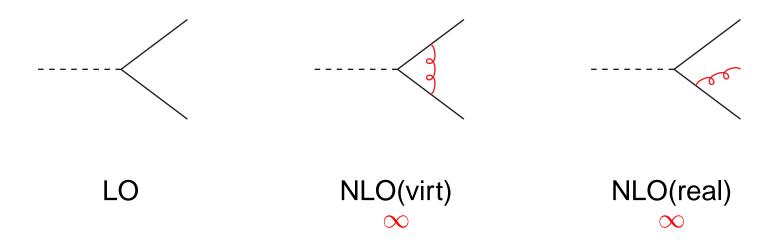
- k_t algorithm
- Cambridge/Aachen alg.

Cone:

- CDF JetClu
- CDF MidPoint
- D0 (run II) Cone
- PxCone
- ATLAS Cone
- CMS Iterative Cone
- PyCell/CellJet
- GetJet

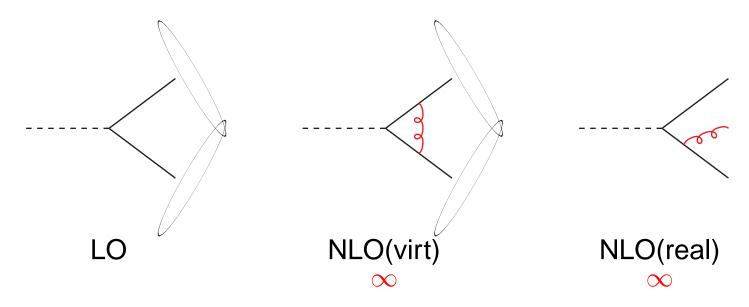
How has that list changed recently?

Ingredient: QCD soft and collinear divergencies



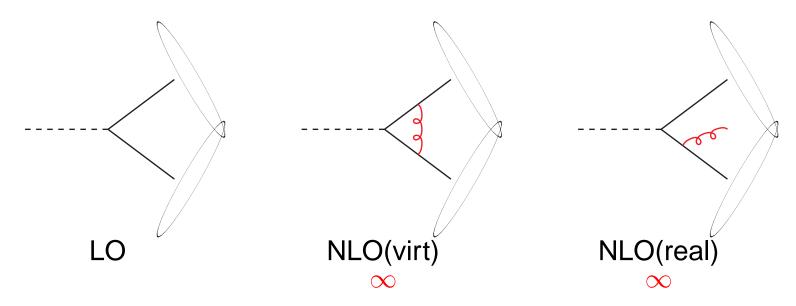
 \bullet ∞ (from soft gluons) cancel (inclusive x-section)

Ingredient: QCD soft and collinear divergencies



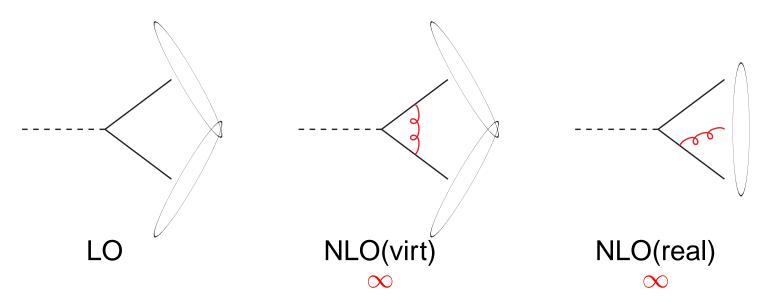
- Consider an extra (NLO) soft gluon
- Assume LO gives 2 jets \Rightarrow NLO(virt) gives 2 jets

Ingredient: QCD soft and collinear divergencies



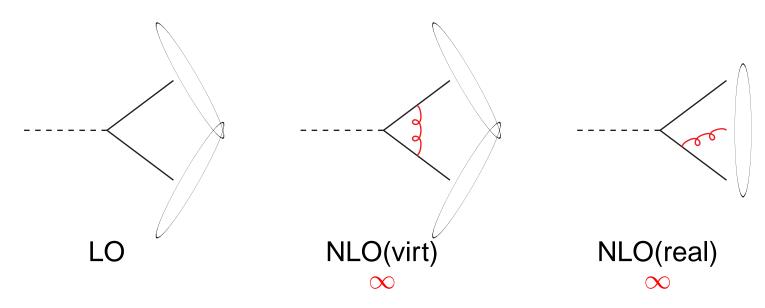
- Consider an extra (NLO) soft gluon
- Assume LO gives 2 jets \Rightarrow NLO(virt) gives 2 jets
- NLO(real) gives 2 jets $\Rightarrow \infty$ cancel \Rightarrow finite jet cross-section

Ingredient: QCD soft and collinear divergencies



- Consider an extra (NLO) soft gluon
- Assume LO gives 2 jets \Rightarrow NLO(virt) gives 2 jets
- NLO(real) gives 2 jets $\Rightarrow \infty$ cancel \Rightarrow finite jet cross-section NLO(real) gives 1 jets $\Rightarrow \infty$ do not cancel \Rightarrow infinite jet x-section

Ingredient: QCD soft and collinear divergencies

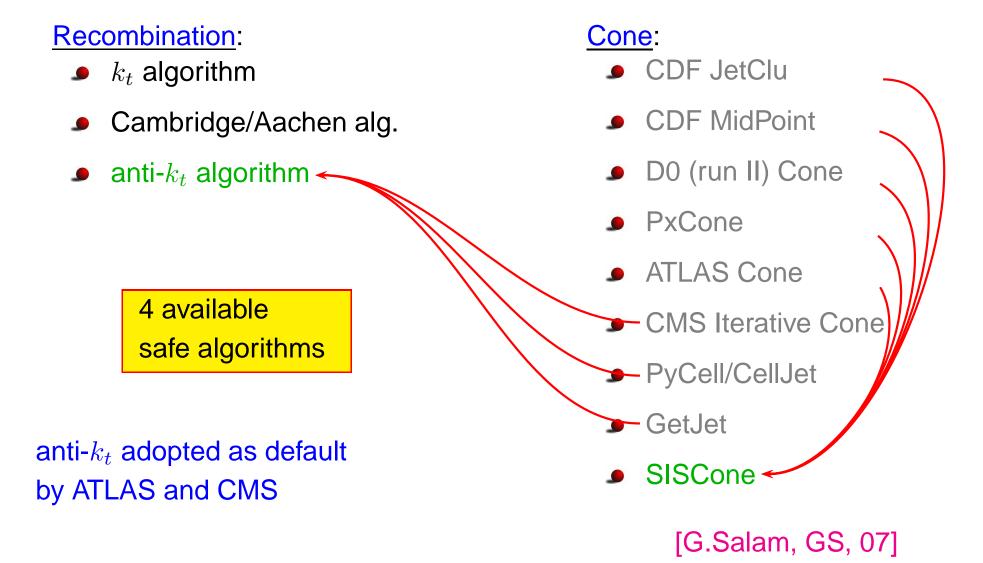


For pQCD to make sense, the (hard) jets should not change when

- one has a soft emission i.e. adds a very soft gluon
- one has a collinear splitting i.e. replaces one parton by two at the same place (η, ϕ)

[SNOWMASS Accords, Fermilab, 1990]

21st century jet algorithms



[M.Cacciari, G.Salam, GS, 08]

Inside the black box

• Recombination: recombine closest pair until d > R

$$d_{ij} = \min\left(k_{t,i}^{2p}, k_{t,j}^{2p}\right) \left(\Delta y_{ij}^2 + \Delta \phi_{ij}^2\right)$$

- $\underline{k_t}$ (p = 1): the closest to QCD matches the soft and collinear divergences of QCD
- Cambridge/Aachem (p = 0): the simplest close to k_t in many respects
- $\frac{\text{anti-}k_t}{\text{calibration advantages, a safe CMS iterative cone}}$
- Cone: find a direction of energy flow stable cone = the total momentum points towards the centre of the cone + split-merge for the overlap
 - SISCone: reduced sensitivity to the UE
 a safe CDF JetClu/Midpoint, D0 MidPoint, ATLAS Cone

Inside the black box

• Recombination: recombine closest pair until d > R

$$d_{ij} = \min\left(k_{t,i}^{2p}, k_{t,j}^{2p}\right) \left(\Delta y_{ij}^2 + \Delta \phi_{ij}^2\right)$$

• $\underline{k_t}$ (p=1): the closest to QCD

matches the soft and collinear divergences of OCD

```
#-----
# FastJet release 2.4
# Written by M. Cacciari, G.P. Salam and G. Soyez
# http://www.fastjet.fr
#-----
```

All those algorithms (and much more) implemented (efficiently) in FastJet

points towards the centre of the cone + split-merge for the overlap

SISCone: reduced sensitivity to the UE a safe CDF JetClu/Midpoint, D0 MidPoint, ATLAS Cone ntum

We (finally) have a good set of tools

Can we do better?

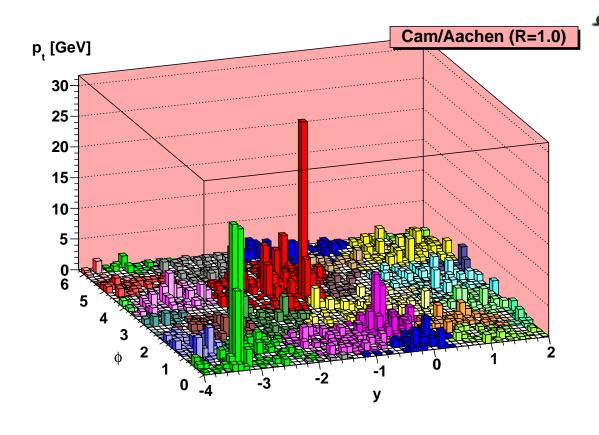
A growing list

Many ideas and applications:

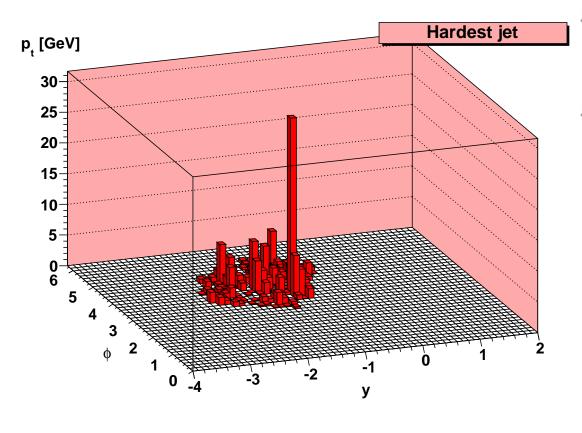
- Handle soft backgrounds
 - ✓ tune the algorithm, filter
 ex.: kinematic dijet reconstruction
 - ✓ jet areas and background subtraction
 → UE, pileup, heavy-ion background subtraction
- Tag boosted objects (mostly for discovery purposes)
 - ✓ subjet analysis $\longrightarrow H \to b\bar{b}, t \to bq\bar{q}, \tilde{\chi}_0^1 \to qqq, \dots$

I will only give a brief overview of what can be done

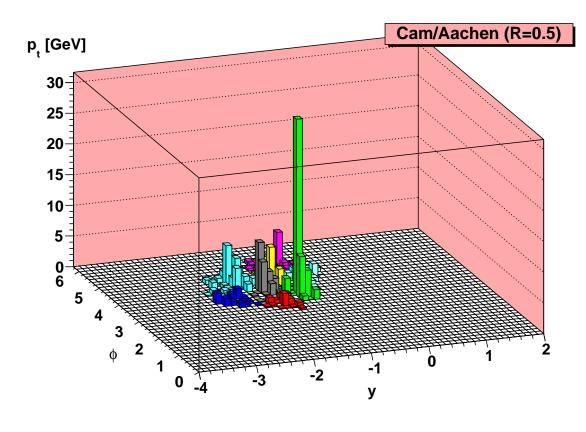
New idea #1: filtering



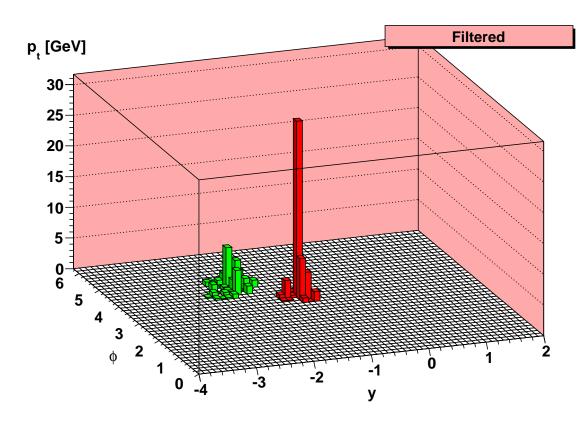
cluster withCambridge/Aachen(R)



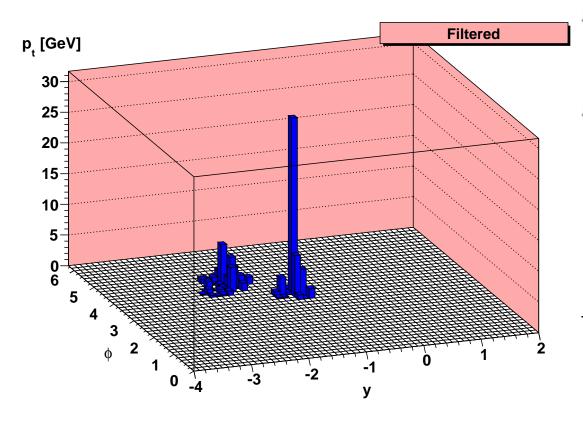
- cluster with Cambridge/Aachen(R)
- for each jet



- cluster with Cambridge/Aachen(R)
- for each jet
 - recluster with Cambridge/Aachen(R/2)



- cluster with Cambridge/Aachen(R)
- for each jet
 - recluster with Cambridge/Aachen(R/2)
 - keep the 2 hardest subjets



- cluster with Cambridge/Aachen(R)
- for each jet
 - recluster with Cambridge/Aachen(R/2)
 - keep the 2 hardest subjets

Idea:

- √ keep perturb. radiation
- √ remove UE
- ullet Proven useful for boosted jet H o bar b tagging

[J.Butterworth, A.Davison, M.Rubin, G.Salam, 08]

Proven useful for kinematic reconstructions

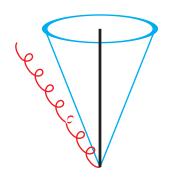
[M.Cacciari, J.Rojo, G.Salam, GS, 08]

Application: optimise dijet kinematic reconstruction

The pQCD-vs-UE paradox

Competition between

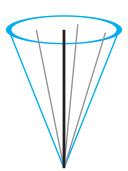
catching perturbative radiation



Out-of-cone radiation:

$$\langle \delta p_t \rangle \propto -\int_R \frac{d\theta}{\theta} \sim -\log(1/R)$$

not catching soft background radiation (underlying event)



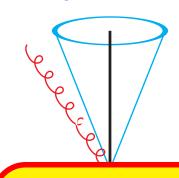
 $\langle \delta p_t \rangle \sim$ Soft contents \propto jet area $\sim R^2$

the coefficients depend on the algorithm

The pQCD-vs-UE paradox

Competition between

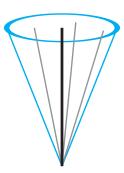
catching perturbative radiation



Out-of-cone radiation:

$$\langle \delta p_t \rangle \propto -\int_R \frac{d\theta}{\theta} \sim -\log(1/R)$$

What is the optimal jet definition (algo+R!)?



 $\langle \delta p_t \rangle \sim$ Soft contents \propto jet area $\sim R^2$

the coefficients depend on the algorithm

Optimisation: dijet reconstruction

Example process to illustrate various effects:

$$Z' \rightarrow q \bar{q} \rightarrow 2$$
 jets

- $M_{Z'}$ can be varied (between 100 GeV and 4 TeV)
- Also valid for $H \rightarrow gg$ to study gluon jets
- Reconstruction method:
 - get the 2 hardest jets: j_1 and j_2
 - reconstruct the Z': $m_{Z'} = (j_1 + j_2)^2$

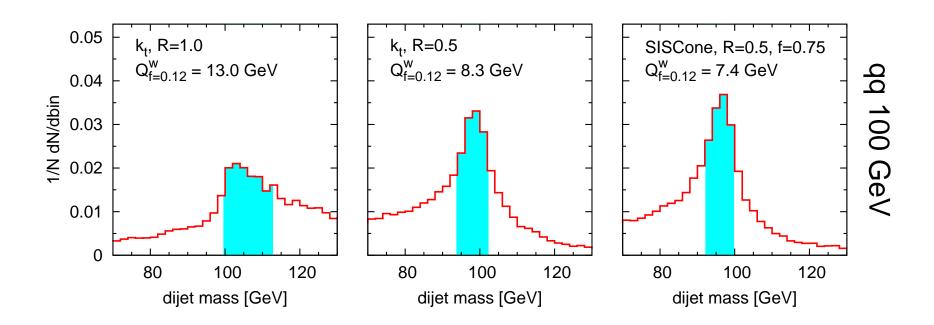
Look how the mass peak is reconstructed

• Also $t\bar{t}$ with full hadronic decay for multijet tests

Optimisation: quality measure

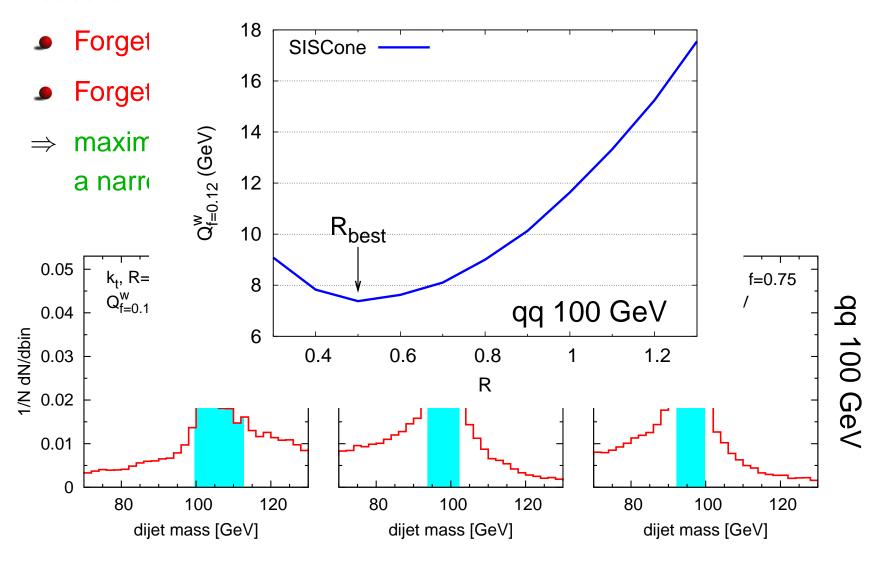
Measure of the jet reconstruction efficiency:

- Forget about measures related to parton-jet matching
- Forget about fits depending on the shape of the peak
- \Rightarrow maximise the signal over background ratio (S/\sqrt{B}) a narrower peak is better.



Optimisation: quality measure

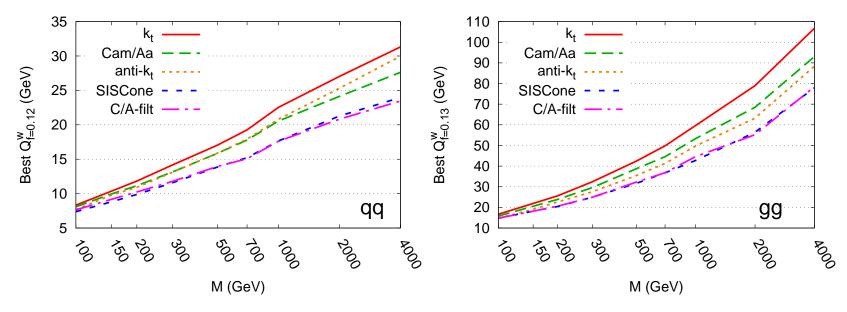
Measure of the jet reconstruction efficiency:



Optimisation: best definition

[M.Cacciari, J.Rojo, G.Salam, GS, 08]

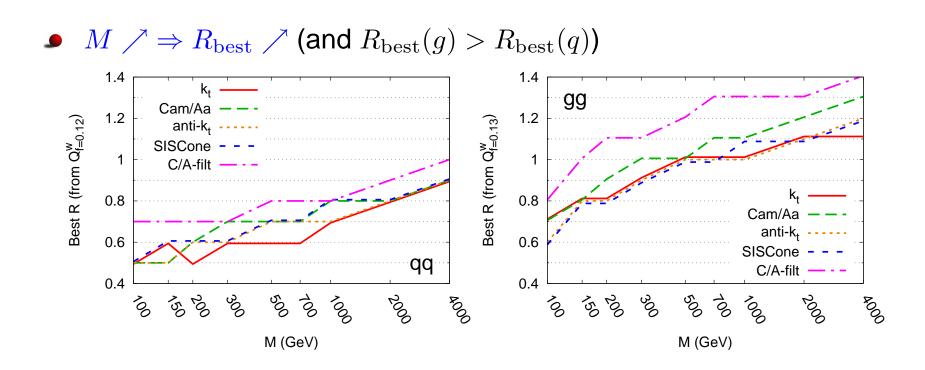
• SISCone and C/A+filt. do slightly better than k_t , C/A or anti- k_t



Optimisation: best definition

[M.Cacciari, J.Rojo, G.Salam, GS, 08]

• SISCone and C/A+filt. do slightly better than k_t , C/A or anti- k_t



Optimisation: quality measure ~ luminosity

Assuming a constant background,

quality measure — effective luminosity ratio

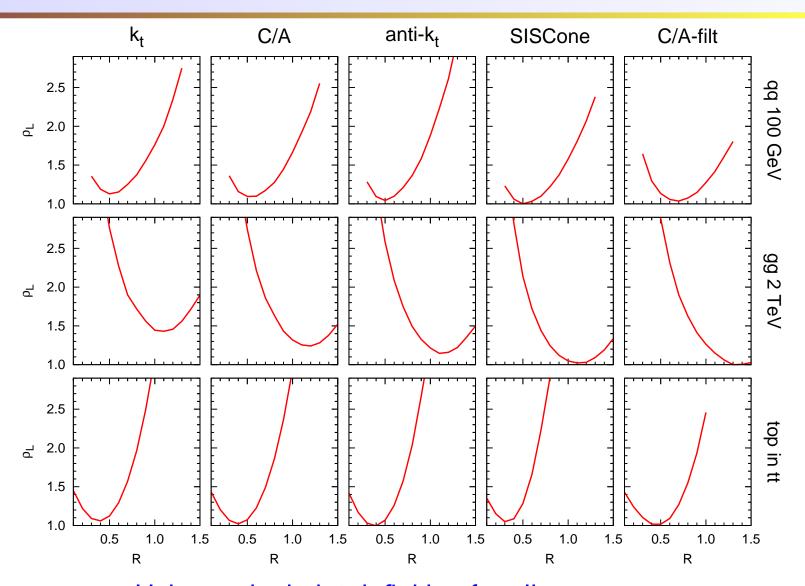
$$\rho_{\mathcal{L}}(\mathrm{JD}_2/\mathrm{JD}_1) = \frac{\mathcal{L} \text{ needed with } \mathrm{JD}_2}{\mathcal{L} \text{ needed with } \mathrm{JD}_1} = \frac{Q_{f=z}^w(\mathrm{JD}_2)}{Q_{f=z}^w(\mathrm{JD}_1)}$$

e.g.
$$\rho_{\mathcal{L}}(JD_2/JD_1) = 2$$

 $\Leftrightarrow \mathrm{JD}_2$ requires 2 times the integrated luminosity of JD_1 to achieve the same discriminative power.

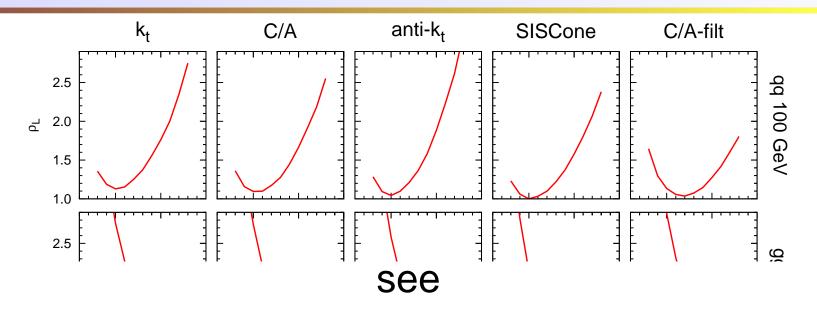
Note: results cross-checked with 2 different definitions of the quality measure

Optimisation: consequences

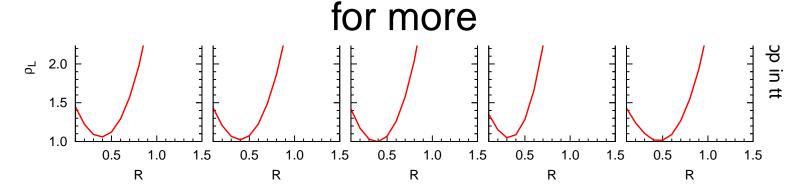


Using a single jet definition for all processes may cost a factor ~ 2 in time for early discoveries at the LHC

Optimisation: consequences



http://quality.fastjet.fr



Using a single jet definition for all processes may cost a factor ~ 2 in time for early discoveries at the LHC

New idea #2: jet area and soft background subtraction

Jet areas

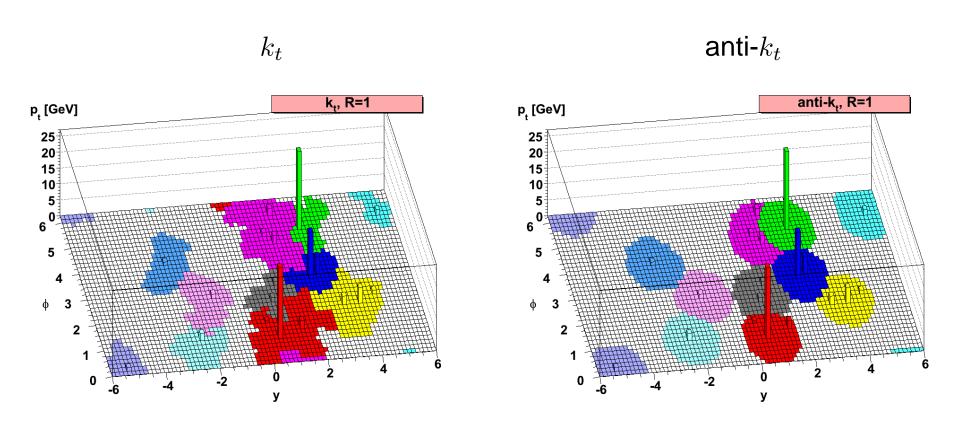
[M.Cacciari, G.Salam, GS, 08]

Area \equiv region where the jet catches soft particles

- Recipe: add infinitely soft particles (aka ghosts) and see in which jet they are clustered
- 2 methods:
 - Passive area: add one ghost at a time and repeat many times
 - Active area: add a set of ghosts and cluster once
- <u>Idea</u>: ghost ≈ background particle
 - \Rightarrow active area \approx uniform background passive area \approx pointlike background
- Notes:
 - passive = active for large multiplicities
 - require an IR-safe algorithm!
 - generic/universal definition (e.g. independent of a calorimeter)

Jet area: examples

Example 1: active area for a simple event



Note: analytic control

Example 2: perturbative expansion of areas (at order α_s)

$$\langle \mathcal{A}(p_t, R) \rangle = \mathcal{A}_0 + \frac{C_{F,A}}{b_0 \pi} \pi R^2 \frac{d}{d} \log \left(\frac{\alpha_s(\mathbf{Q}_0)}{\alpha_s(Rp_t)} \right)$$

• area $\neq \pi R^2$, area \neq const.

coefficients computable

ble	$\mathcal{A}_0/(\pi R^2)$		d	
DIC	passive	active	passive	active
k_t	1	0.81	0.56	0.52
Cam/Aachen	1	0.81	0.08	0.08
anti- k_t	1	1	0	0
SISCone	1	1/4	-0.06	0.12

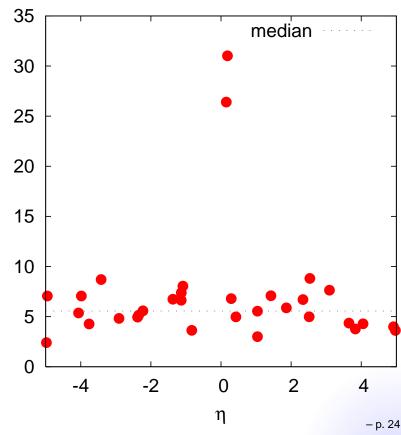
• $Q_0 \equiv IR$ regulator \propto background density

Pileup subtraction (for uniform backgrounds)

Basic idea: [M.Cacciari, G.Salam, 08]

$$p_{t, \text{subtracted}} = p_{t, \text{jet}} - \rho_{\text{pileup}} \times \text{Area}_{\text{jet}}$$

- Jet area: [M.Cacciari, G.Salam, G.S., 08]
 - region where the jet catches infinitely soft particles (active/passive)
 - analytic control and understanding in pQCD
- Pileup density per unit area: $\rho_{
 m pileup}$ e.g. estimated from the median of $p_{t,
 m jet}/{
 m Area}_{
 m jet}$

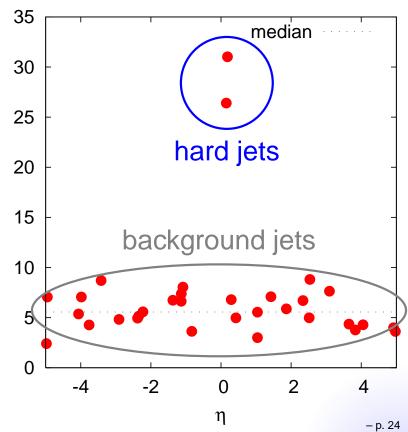


Pileup subtraction (for uniform backgrounds)

Basic idea: [M.Cacciari, G.Salam, 08]

$$p_{t, \text{subtracted}} = p_{t, \text{jet}} - \rho_{\text{pileup}} \times \text{Area}_{\text{jet}}$$

- Jet area: [M.Cacciari, G.Salam, G.S., 08]
 - region where the jet catches infinitely soft particles (active/passive)
 - analytic control and understanding in pQCD
- Pileup density per unit area: $\rho_{
 m pileup}$ e.g. estimated from the median of $p_{t,
 m jet}/{
 m Area}_{
 m jet}$



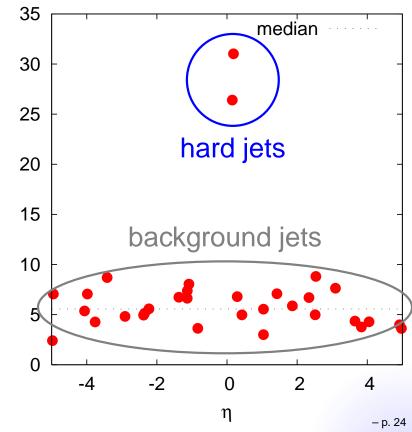
Pileup subtraction (for uniform backgrounds)

Basic idea: [M.Cacciari, G.Salam, 08]

$$p_{t, \text{subtracted}} = p_{t, \text{jet}} - \rho_{\text{pileup}} \times \text{Area}_{\text{jet}}$$

- Jet area: [M.Cacciari, G.Salam, G.S., 08]
 - region where the jet catches infinitely soft particles (active/passive)
 - analytic control and understanding in pQCD
- Pileup density per unit area: ρ_{pileup} e.g. estimated from the median of $p_{t,\mathrm{jet}}/\mathrm{Area}_{\mathrm{jet}}$

implemented in FastJet on an event-by-event basis



Effect on dijet reconstruction

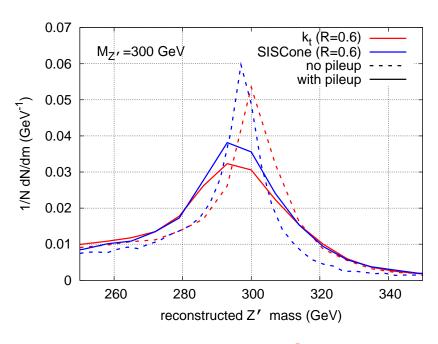
Pileup unsubtracted

0.07 $k_t (R=0.6)$ M_{7} , =300 GeV SISCone (R=0.6) 0.06 no pileup with pileup 0.05 $1/N \, dN/dm \, (GeV^{-1})$ 0.04 0.03 0.02 0.01 0 280 340 260 300 320 reconstructed Z' mass (GeV)

width = 29.5 GeV

width = 21.0 GeV

pileup subtracted



width = 21.0 GeV

width = 17.7 GeV

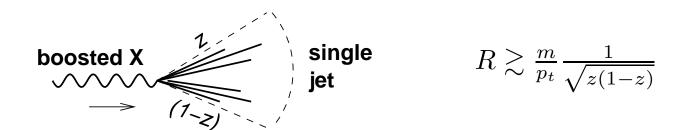
- √ position reasonnable
- √ dispersion reduced (thanks to the event-by-event approach)
- √ used by STAR for the first jet analysis in heavy-ions

New idea #3: Subjets help tagging boosted objects

Example: boosted Higgs

[J.Butterworth, A.Davison, M.Rubin, G.Salam,08]

- $H \rightarrow b\bar{b}$: dominant decay for small M_H but large backgrounds
- boosted H (HW, HZ): many advantages (e.g. no $t\bar{t}$ background), main problem: small cross-section
- boosted particle: decay products in the same jet



Note: other similar examples:

- boosted top
- $t\bar{t}H$
- $\tilde{\chi}^0 \to qqq$

Example: boosted Higgs

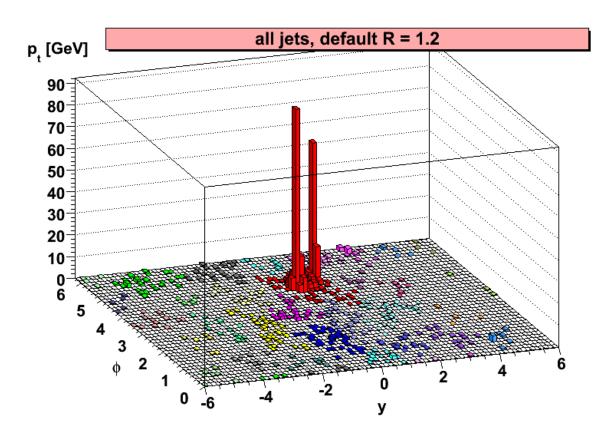
[J.Butterworth, A.Davison, M.Rubin, G.Salam,08]



Method: start with a hard (C/A, radius R) jet j

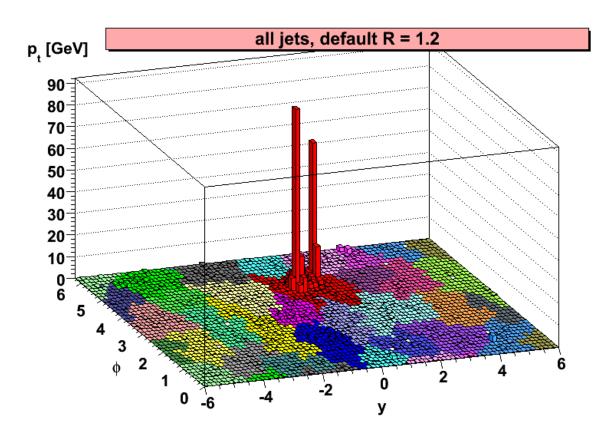
- **①** Undo the last clustering $\rightarrow j_1, j_2$
- If $\max(m_1, m_2) < 0.67m$, we have a mass drop, else back to 1 idea: find the 2 b-jets, dynamically find R_{bb}
- Require symmetric splitting $y_{12} \approx \frac{\min(z_1, z_2)}{\max(z_1, z_2)} > 0.09$, else go to 1 idea: remove QCD asymmetric splittings
- Require 2 b taggings
- Filter *i.e.* uncluster down to $R_{\rm filt}=R/3$, keep the 3 hardest subjets idea: keep "hard" QCD radiations, reduce UE

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



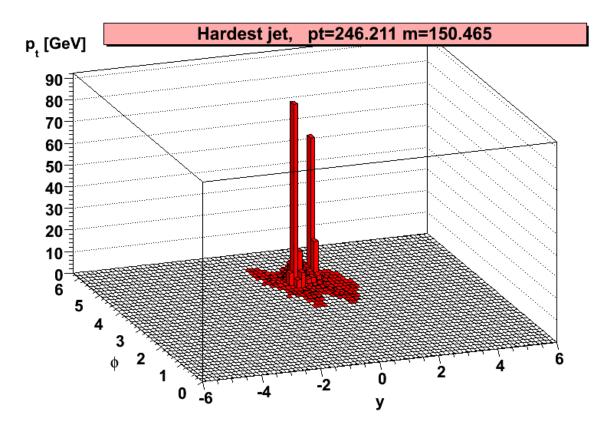
Cluster C/A, R=1.2

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3

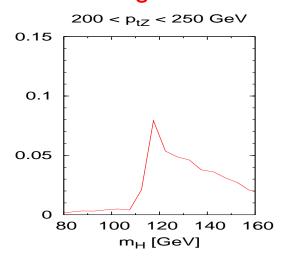


Show jets more clearly

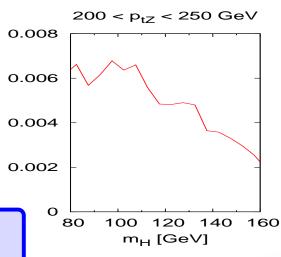
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



HZ Signal

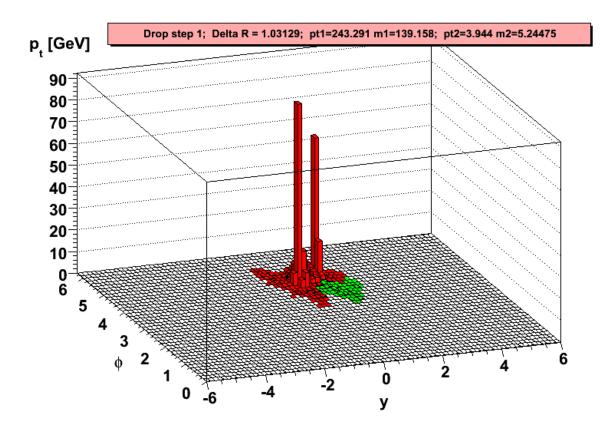


Zbb Background

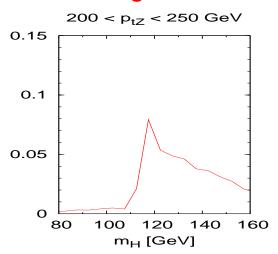


Hardest jet (m = 150 GeV)

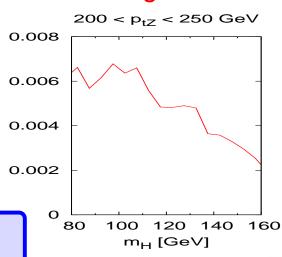
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



HZ Signal

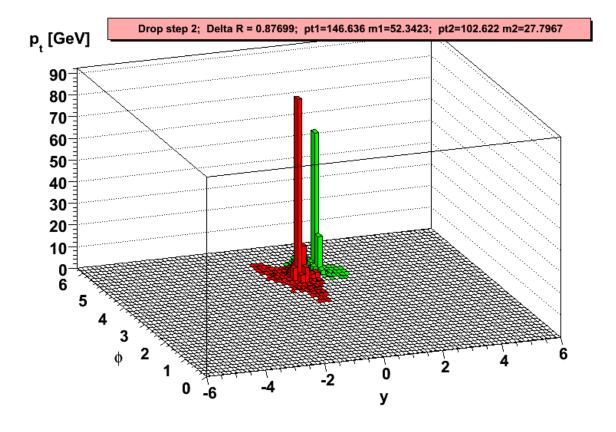


Zbb Background

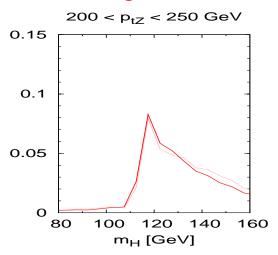


Split: $\frac{\max(m_1, m_2)}{m} = 0.92$, repeat (m = 150 GeV)

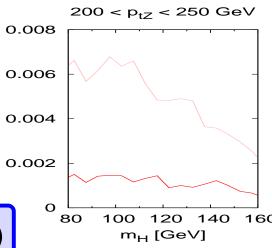
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



HZ Signal

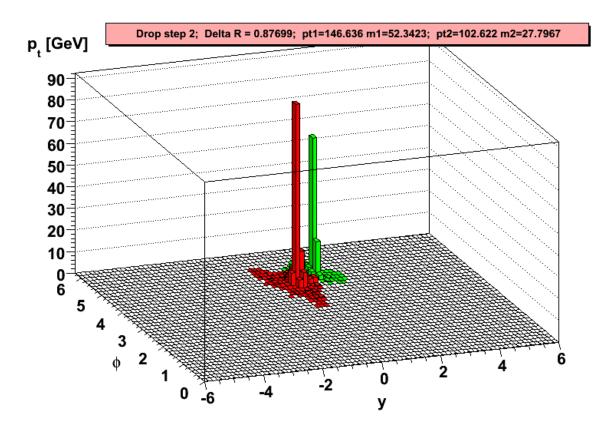


Zbb Background

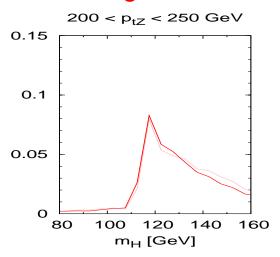


Split: $\frac{\max(m_1, m_2)}{m} = 0.37$, mass drop (m = 139 GeV)

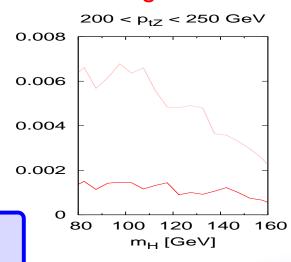
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



HZ Signal

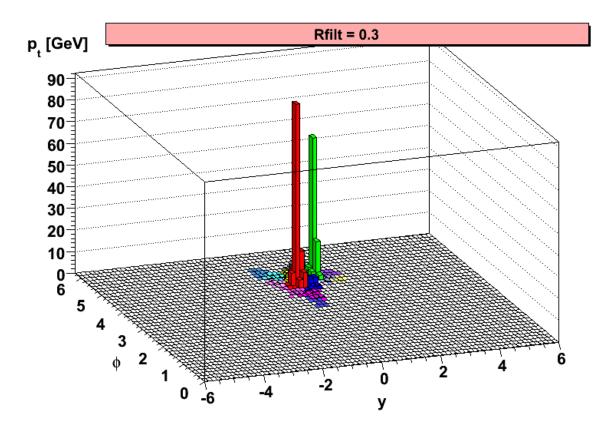


Zbb Background

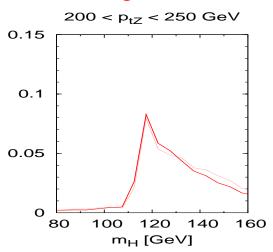


Split: $y_{12} = 0.7$, 2 *b* tags \Rightarrow OK (m = 139 GeV)

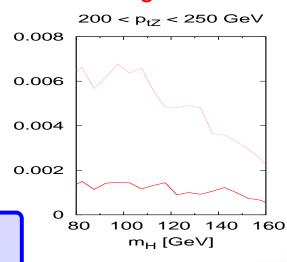
Herwig 6.510 + Jimmy 4.31 + FastJet 2.3



HZ Signal

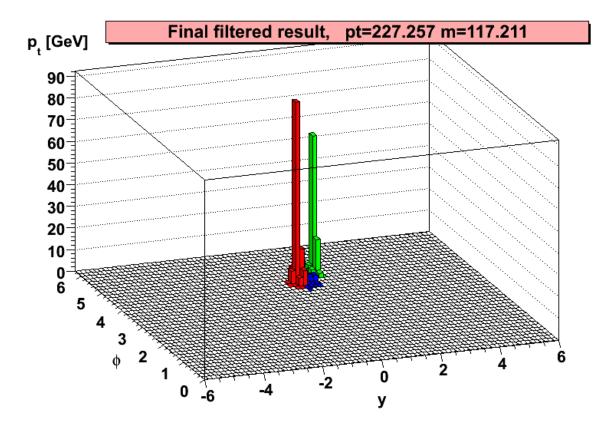


Zbb Background



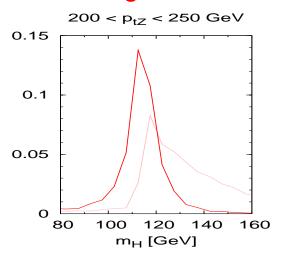
Re-cluster: $R_{\rm filt} = 0.3$

Herwig 6.510 + Jimmy 4.31 + FastJet 2.3

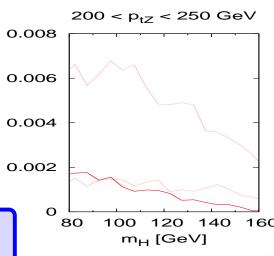


Filter: keep 3 hardets (m = 117 GeV)

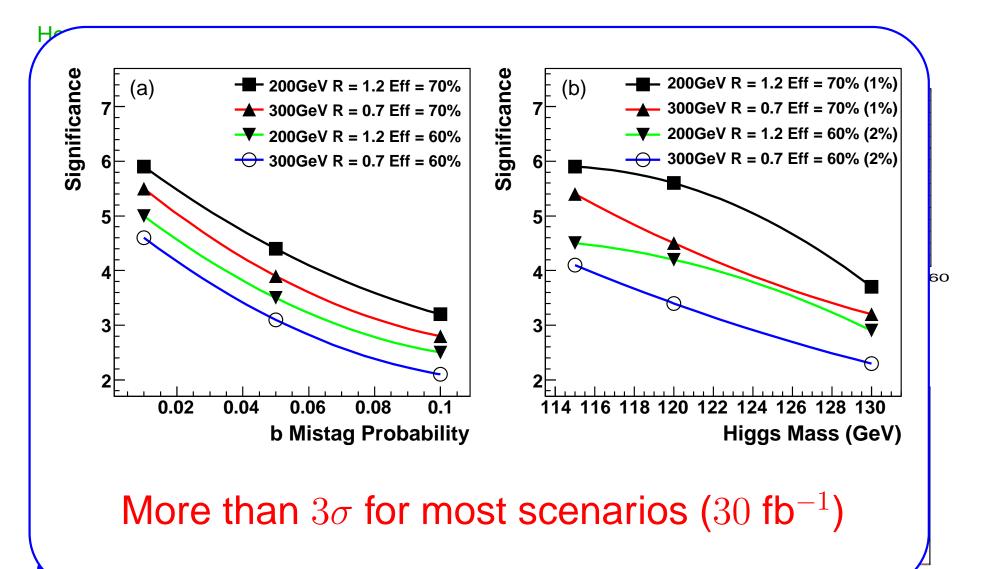
HZ Signal



Zbb Background



Titel. Reep 3 Haruels (711)

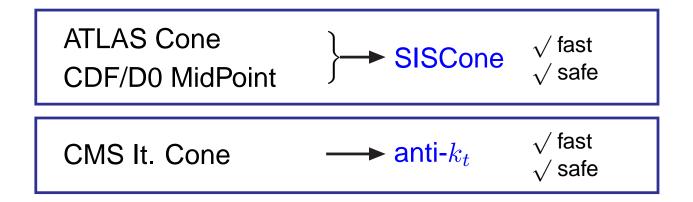


- p. 28

Summary (1)

Message #1:

Use infrared-and-collinear-safe algorithms



Mandatory to benefit fully from pQCD multileg/loop calculations at the LHC

Summary (2)

Message #2:

correct tools ⇒ new ideas, new concepts ⇒ new generation of jet definitions

- jet areas → pileup and HI background subtraction
- jet substructure improves reconstruction (Higgs, top, SUSY, ...)

Message #3:

keep some flexibility in the jet definition choice

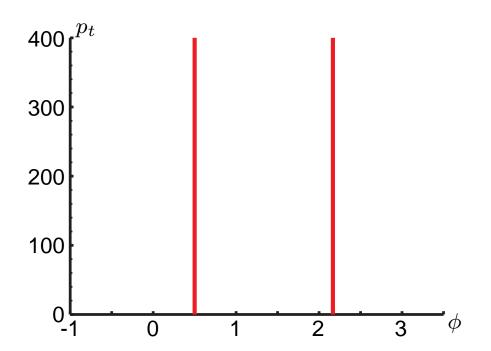
● optimisation → luminosity gains for LHC searches

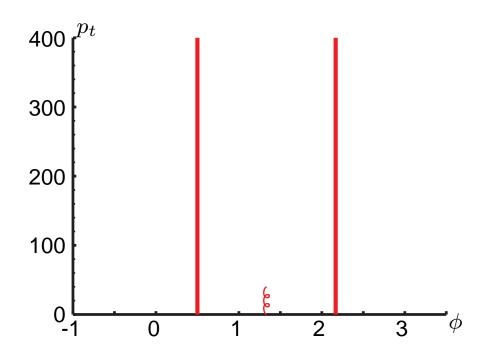
Future perspectives

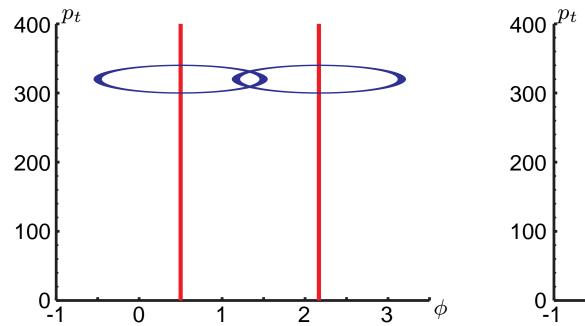
This is not the end of the story:

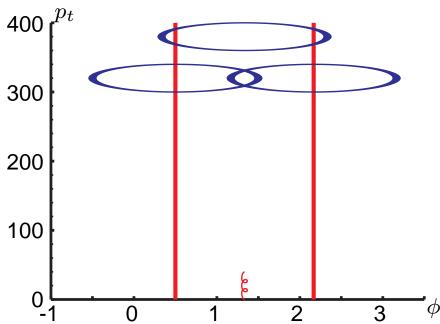
- Study the parameters used when filtering
- ullet Analytic understanding of the optimisation of R
- Optimisation for other processes

backup slides

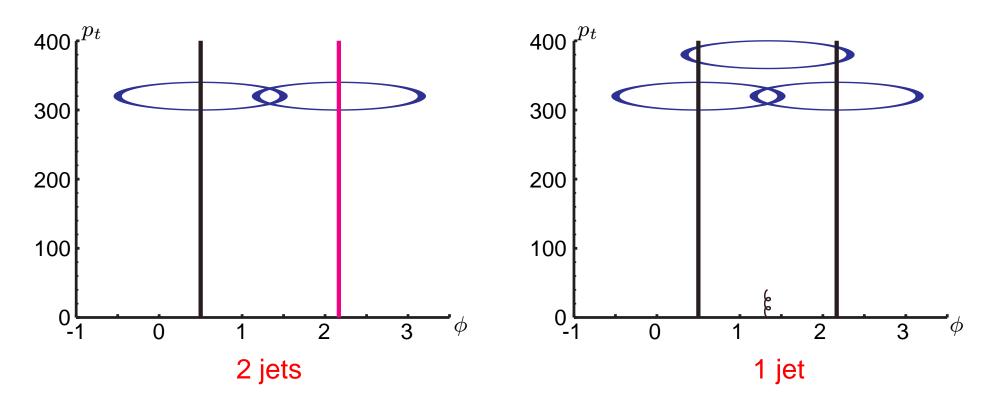






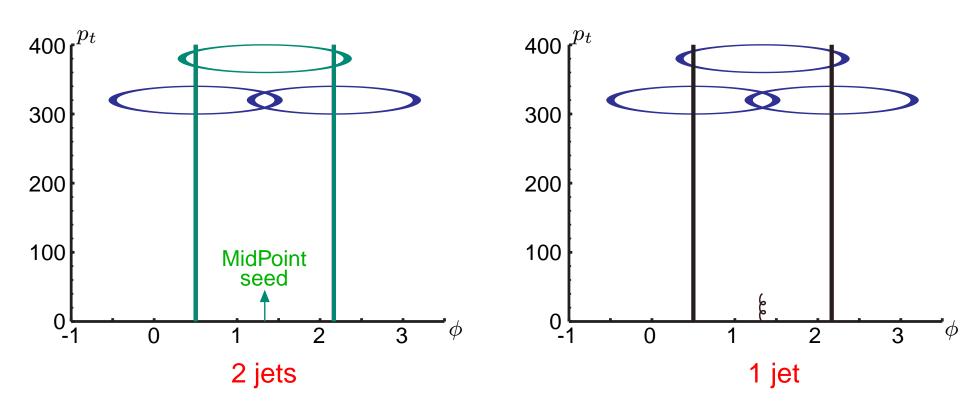


Stable cones found



A soft gluon changed the number of jets

⇒ IR unsafety of JetClu and the ATLAS Cone

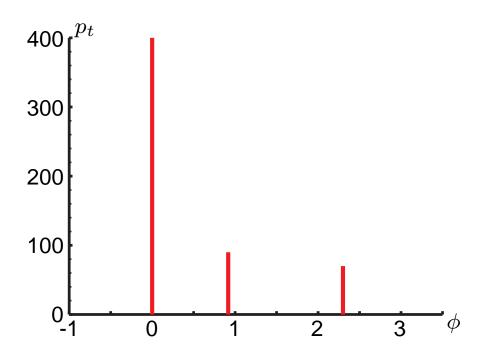


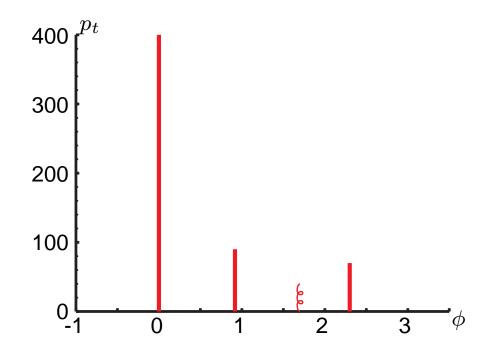
A soft gluon changed the number of jets

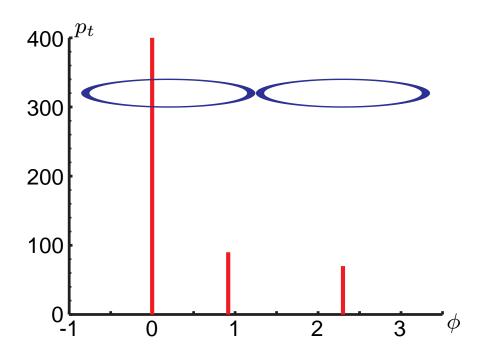
⇒ IR unsafety of JetClu and the ATLAS Cone

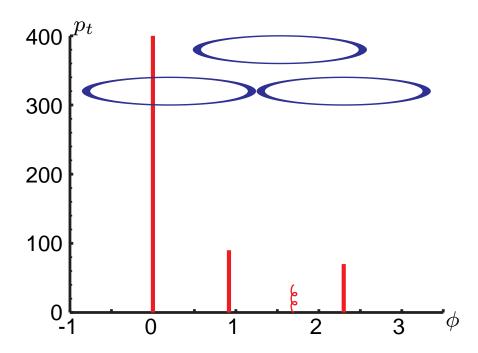
Fixed by MidPoint

[Blazey et al., 00]

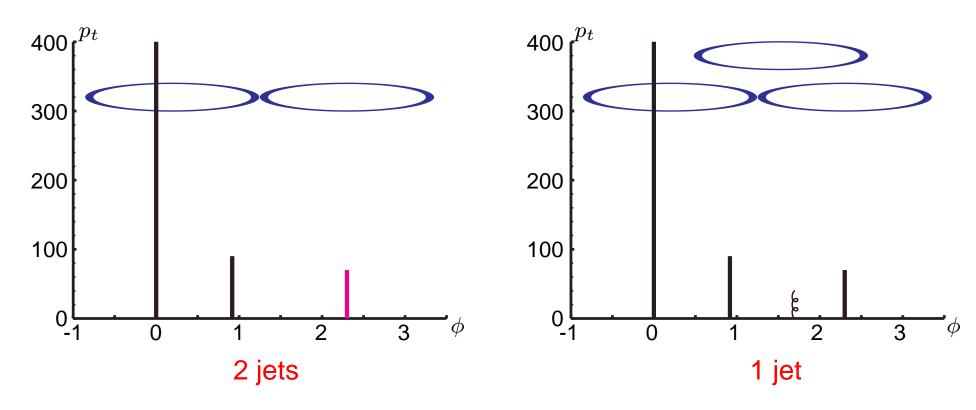






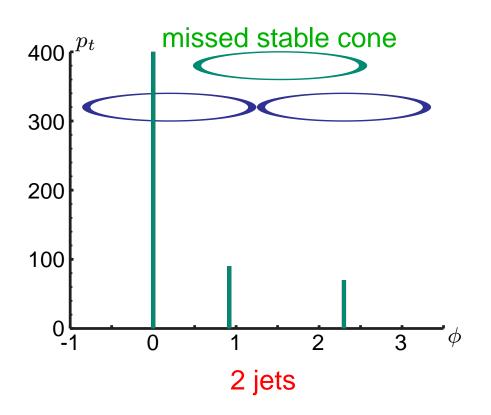


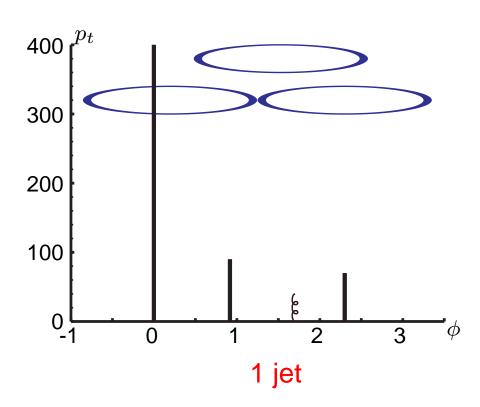
Stable cones found



A soft gluon changed the number of jets

 \Rightarrow IR unsafety of MidPoint (1 order in α_s later than JetClu)





Solution: be sure to find all stable cones

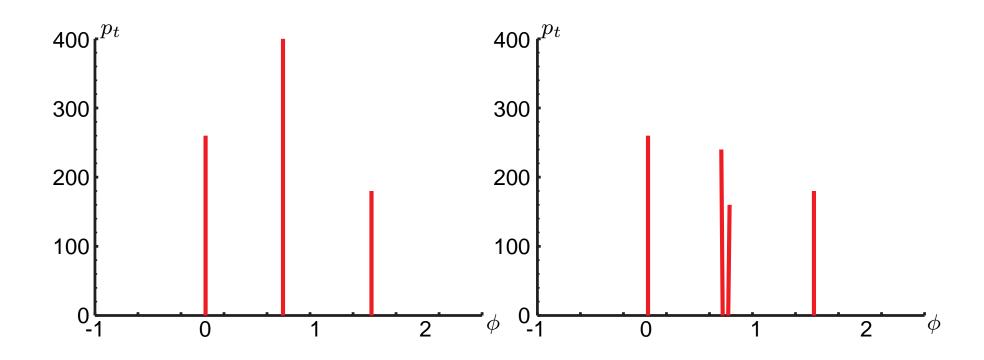
SISCone: Seedless Infrared-Safe Cone algorithm

http://projects.hepforge.org/siscone

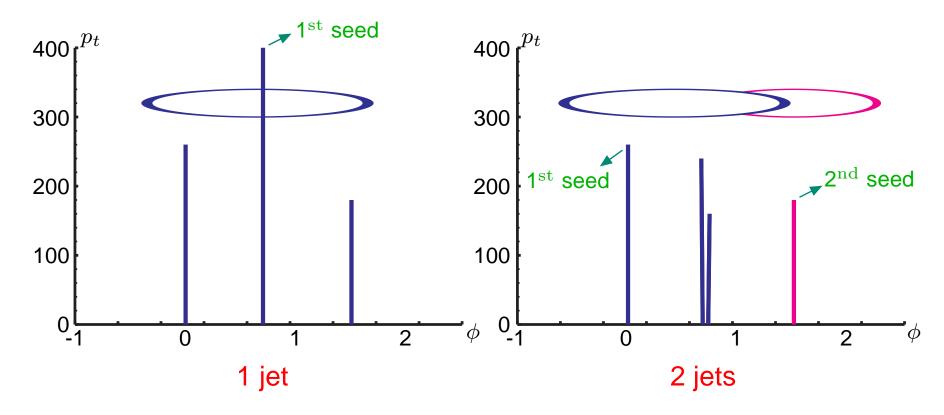
[G.Salam, G.S., 07]

Idea: enumerate enclosures by enumerating pairs of particles

Collinear (un)safety? the CMS iterative cone



Collinear (un)safety? the CMS iterative cone



A colinear splitting changed the number of jets

⇒ Collinear unsafety of the CMS iterative cone

Anti- k_t

Come back to recombination-type algorithms:

$$d_{ij} = \min(k_{t,i}^{2p}, k_{t,j}^{2p}) \left(\Delta \phi_{ij}^2 + \Delta \eta_{ij}^2\right)$$

- p=1: k_t algorithm
- p = 0: Aachen/Cambridge algorithm



Come back to recombination-type algorithms:

$$d_{ij} = \min(k_{t,i}^{2p}, k_{t,j}^{2p}) \left(\Delta \phi_{ij}^2 + \Delta \eta_{ij}^2\right)$$

- p=1: k_t algorithm
- p = 0: Aachen/Cambridge algorithm
- p = -1: anti- k_t algorithm [M.Cacciari, G.Salam, G.S., 08]

Why should that be related to the iterative cone?!?

- "large $k_t \Rightarrow$ small distance"

 i.e. hard partons "eat" everything up to a distance Ri.e. circular/regular jets, jet borders unmodified by soft radiation
- infrared and collinear safe

Take e.g. the MidPoint cone

2 particles 3 particles 4 particles 4 particles 4 particles
$$\alpha_s^2 \times \ldots + \alpha_s^3 \times \ldots + \alpha_s^4 \times \ldots + \alpha_s^5 \times \ldots + \ldots$$

ullet QCD expansion (one $lpha_s$ can be replaced by $lpha_{
m EW}$)

Take e.g. the MidPoint cone

2 particles 3 particles 4 particles 4 particles 4 particles + 1 soft
$$\overbrace{\alpha_s^2 \times \ldots}^2 + \overbrace{\alpha_s^3 \times \ldots}^3 + \overbrace{\alpha_s^4 \times \ldots}^4 + \overbrace{\alpha_s^5 \times \log(p_t/\Lambda_{\rm QCD}) \ldots}^5 + \ldots$$

- QCD expansion (one α_s can be replaced by $\alpha_{\rm EW}$)
- ullet IRC unsafety (regulated at the hadronic scale $\sim \Lambda_{
 m QCD}$)

Take e.g. the MidPoint cone

2 particles 3 particles 4 particles 4 particles 4 particles + 1 soft
$$\overbrace{\alpha_s^2 \times \ldots}^2 + \overbrace{\alpha_s^3 \times \ldots}^3 + \underbrace{\alpha_s^4 \times \ldots}^4 + \underbrace{\alpha_s^5 \times \log(p_t/\Lambda_{\rm QCD}) \ldots}_{\rm cannot \ be \ trusted} + \ldots$$

- QCD expansion (one α_s can be replaced by α_{EW})
- ullet IRC unsafety (regulated at the hadronic scale $\sim \Lambda_{
 m QCD}$)
- $\alpha_s \log(p_t/\Lambda_{\rm QCD}) \sim 1$
- ullet last meaningful order = $lpha_{f s}^{f 3}$ or $lpha_{
 m EW}lpha_{f s}^{f 2}$

Take e.g. the MidPoint cone

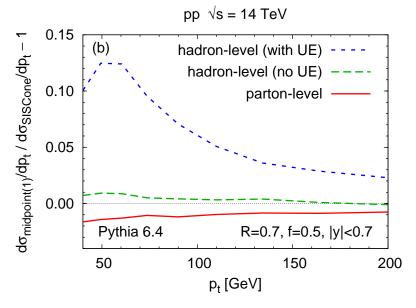
2 particles 3 particles 4 particles 4 particles 4 particles + 1 soft
$$\overbrace{\alpha_s^2 \times \ldots}^2 + \overbrace{\alpha_s^3 \times \ldots}^3 + \underbrace{\alpha_s^4 \times \ldots}^4 + \underbrace{\alpha_s^5 \times \log(p_t/\Lambda_{\rm QCD}) \ldots}_{\rm cannot \ be \ trusted} + \ldots$$

- QCD expansion (one α_s can be replaced by $\alpha_{\rm EW}$)
- ullet IRC unsafety (regulated at the hadronic scale $\sim \Lambda_{
 m QCD}$)
- $\alpha_s \log(p_t/\Lambda_{\rm QCD}) \sim 1$
- ullet last meaningful order = $lpha_{f s}^{f 3}$ or $lpha_{
 m EW}lpha_{f s}^{f 2}$
- same argument for the Iterative Cone
- 1 order worse for JetClu or the ATLAS cone

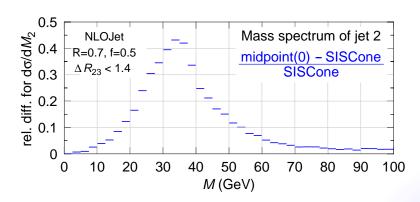
Physical impact

	Last meaningful order		
Observable	MidPoint/CMS	JetClu/ATLAS	
Inclusive jet cross sect.	NLO	LO (NLOJet: NLO)	
3 jet cross section	LO	none (NLOJet: NLO)	
W/Z/H + 2 jet x-sect.	LO	none (MCFM: NLO)	
jet masses in 3 jets	none	none (NLOJet: LO)	

Example: (Midpoint-SISCone)/SISCone



- Incl. cross-section: a few %
- Masses in 3-jet events: $\sim 45\%$



Physical impact

	Last meaningful order		
Observable	MidPoint/CMS	JetClu/ATLAS	
Inclusive jet cross sect.	NLO	LO (NLOJet: NLO)	
3 jet cross section	LO	none (NLOJet: NLO)	
W/Z/H + 2 jet x-sect.	LO	none (MCFM: NLO)	
jet masses in 3 jets	none	none (NLOJet: LO)	



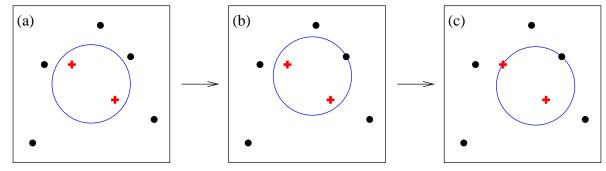
Huge effort ($\sim 50~\text{M}{\equiv}$) to compute processes in pQCD

- Solution: use a seedless approach, find ALL stable cones
- Naive approach: check stability of each subset of particle

- Solution: use a seedless approach, find ALL stable cones
- Naive approach: check stability of each subset of particle Complexity is $\mathcal{O}\left(N2^N\right)$
 - \Rightarrow definitely unrealistic: 10^{17} years for N=100
- Midpoint complexity: $\mathcal{O}\left(N^3\right)$

- Solution: use a seedless approach, find ALL stable cones
- Midpoint complexity: $\mathcal{O}\left(N^3\right)$

<u>Idea</u>: use geometric arguments



- Each enclosure can be moved (in any dir.) until it touches a point
- ... then rotated until it touches a second one
- ⇒ Enumerate all pairs of particles
 with 2 circle orientations and 4 possible inclusion/exclusion
 → find all enclosures

- Solution: use a seedless approach, find ALL stable cones
- Midpoint complexity: $\mathcal{O}\left(N^3\right)$

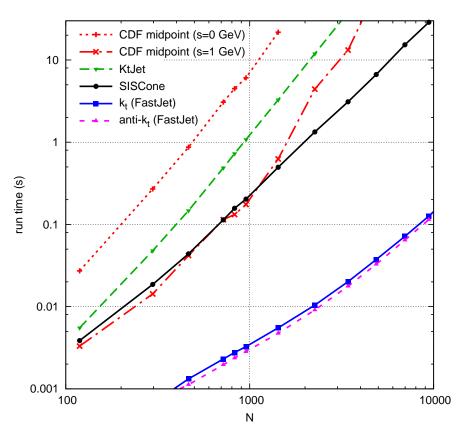
<u>Idea</u>: use geometric arguments

- ⇒ Enumerate all pairs of particles with 2 circle orientations and 4 possible inclusion/exclusion
- → find all enclosures
- Complexity: $\mathcal{O}\left(N^3\right)$, with improvements: $\mathcal{O}\left(N^2\log(N)\right)$

— C++ implementation: Seedless Infrared-Safe Cone algorithm (SISCone) G.Salam, G.S., JHEP 04 (2007) 086; http://projects.hepforge.org/siscone

NB.: also available from FastJet [M.Cacciari, G.Salam, G.S.]; http://www.fastjet.fr

Algorithm timings



Recombination algorithms very fast

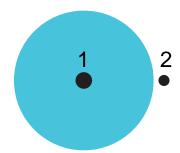
[M. Cacciari, G. Salam, 06]

 SISCone not slower than Midpoint (even with a 1 GeV seed threshold)

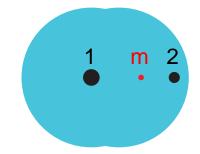
A technical point: Back-reaction

Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
- Modify the hard scattering (back-reaction)
 - can be pointlike or diffuse
 - gain:

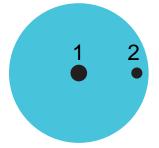


no medium: $p_t = p_{t1}$

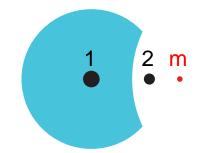


medium: $p_t = p_{t1} + p_{t2} + p_{tm}$

OSS:



no medium: $p_t = p_{t1} + p_{t2}$



medium: $p_t = p_{t1} + p_{tm}$

A technical point: Back-reaction

Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
- Modify the hard scattering (back-reaction)
 - can be pointlike or diffuse
 - tractable analytically (similar to areas)
 - $k_t \gtrsim$ Cambridge > SISCone \gg anti- k_t

