

Optimising jet finding in *pp* **and** *AA* **collisions**

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C. Buttar *et al.*, arXiv:0803.0678 M. Cacciari, J. Rojo, G.P. Salam, G. Soyez, in prep.

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Foreword: why jets?

introducing the basic terms

Part 0: a solid toolkit

jet definitions meeting the fundamental requirements

- Part 1: "best" jet finder in pp collisions
 - sample processes to study
 - figure of merit for quality measure
 - results (1)
 - analytic insight

Part 2: when pileup enters the game

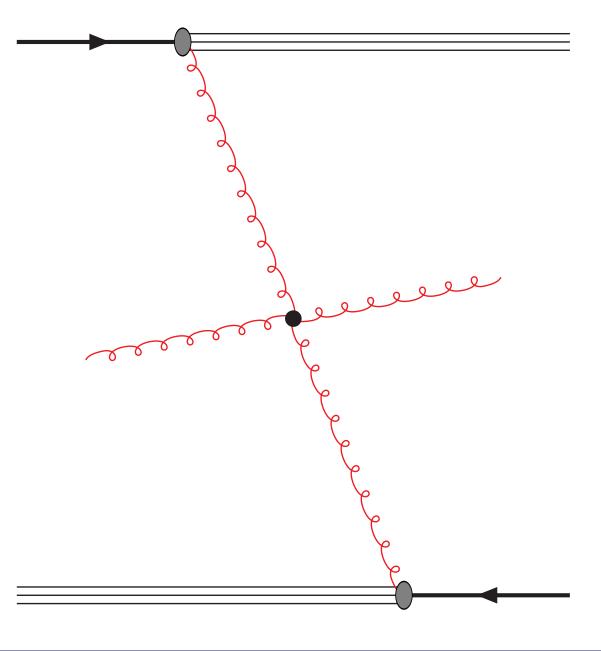
- subtracting soft background using jet areas
- results (2)
- Part 3: heavy-ion background
 - subtraction subtleties
 - results (3)



Foreword: why jets?

optimising jet-finding – p. 3/34

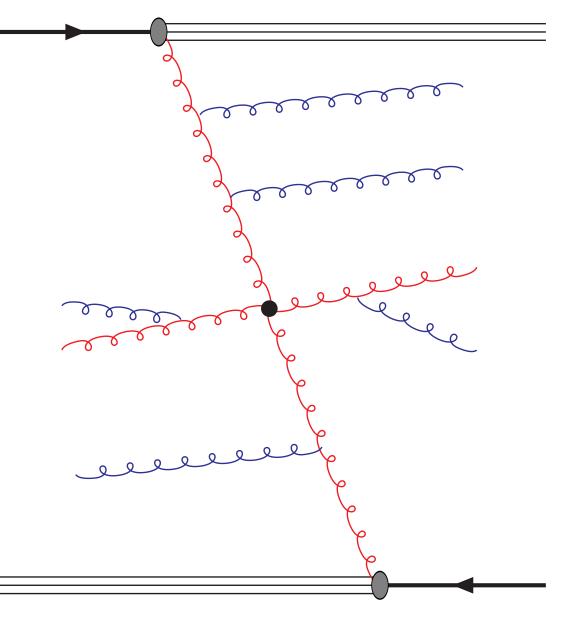




Perturbative level

Hard scattering $2 \rightarrow n$ computed exactly at $\mathcal{O}(\alpha_s^p)$ $gg \rightarrow gg, gg \rightarrow ggg,$ $gg \rightarrow gggg,$ $gg \rightarrow H \rightarrow b\bar{b},$ $gg \rightarrow t\bar{t} \rightarrow \mu\nu_{\mu}b\bar{b}q\bar{q},$ $gg \rightarrow Z' \rightarrow q\bar{q}, \dots$



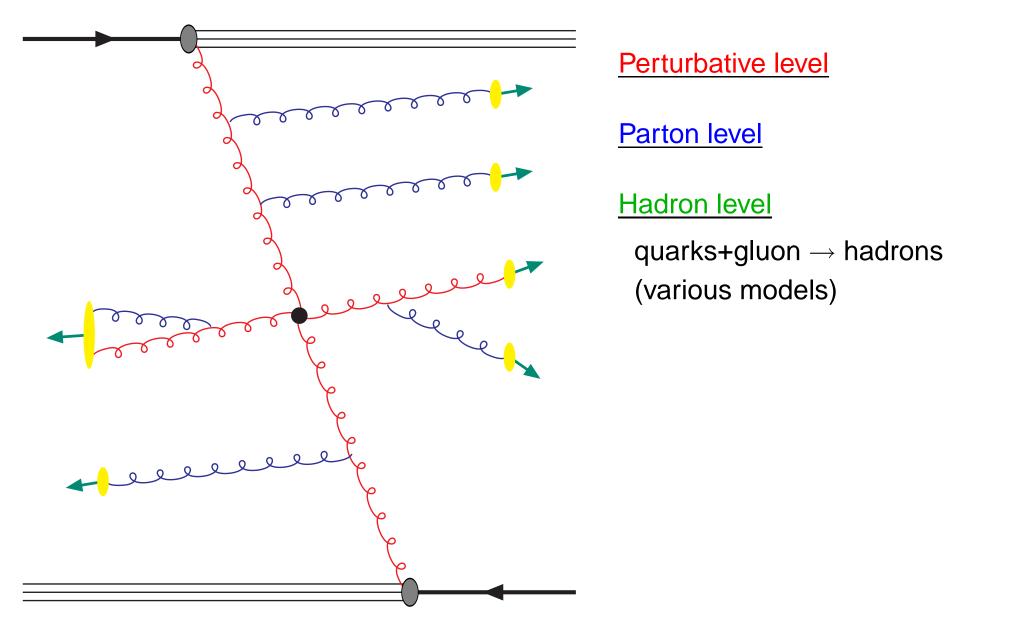


Perturbative level

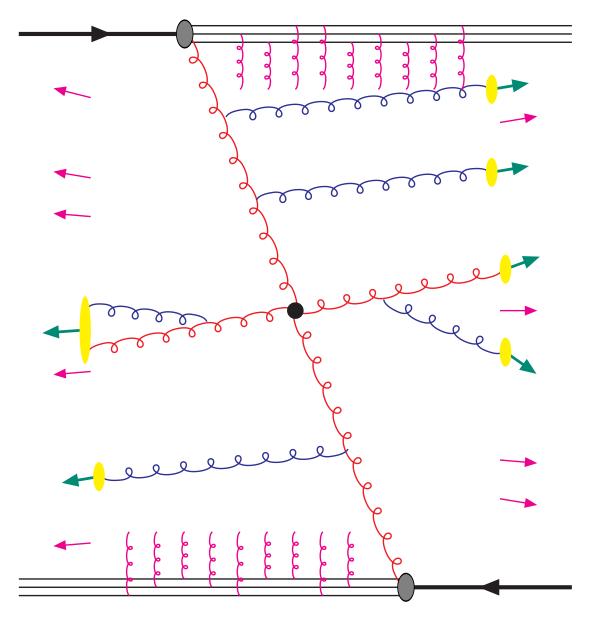
Parton level

pprox collinear divergences resummation $\sum_i lpha_s^i \log^i(p_t^2/\mu^2)$









Perturbative level

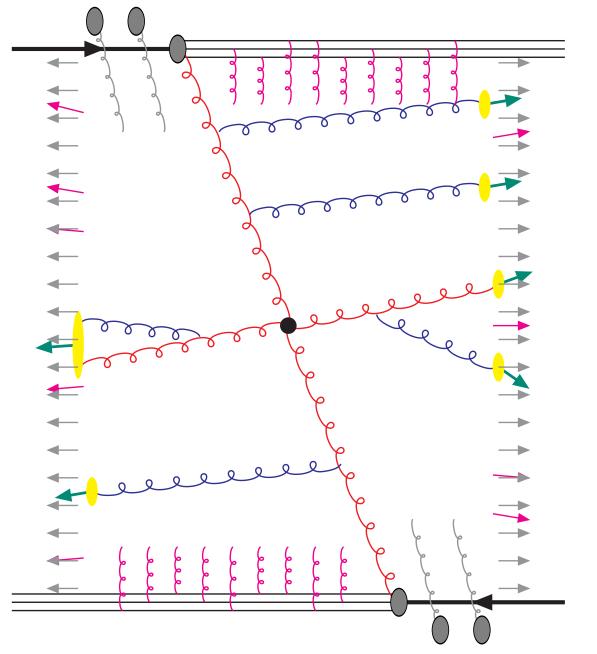
Parton level

Hadron level

+ Underlying event

Multiple interactions from beam remnants \Rightarrow soft background





Perturbative level

Parton level

Hadron level

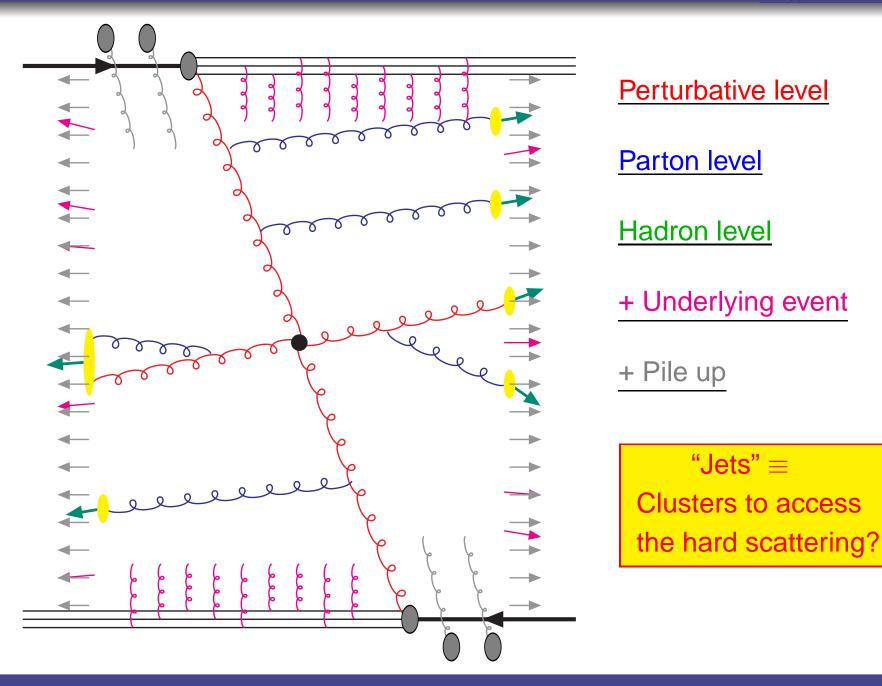
+ Underlying event

+ Pile up

additionnal pp interactions

- \Rightarrow soft background
 - pprox uniform







Part 0: a solid toolkit

optimising jet-finding – p. 5/34



SNOWMASS accords, Tevatron 1990 (i.e. old!):

Several important properties that should be met by a jet definition are [3]:

- 1. Simple to implement in an experimental analysis;
- 2. Simple to implement in the theoretical calculation;
- 3. Defined at any order of perturbation theory;
- 4. Yields finite cross section at any order of perturbation theory;
- 5. Yields a cross section that is relatively insensitive to hadronization.

i.e. usable by theoreticians (*e.g.* finite perturbative results) and experimentalists (*e.g.* fast enough, not much UE sensitivity)

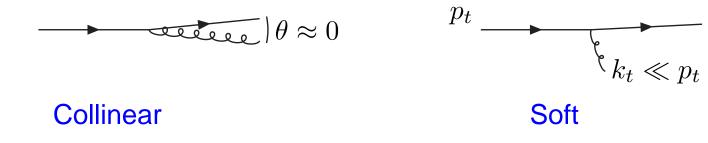
QCD divergences



QCD probability for gluon bremsstrahlung at angle θ and \perp -mom. k_t :

$$dP \propto \alpha_s \, \frac{d\theta}{\theta} \, \frac{dk_t}{k_t}$$

Two divergences:



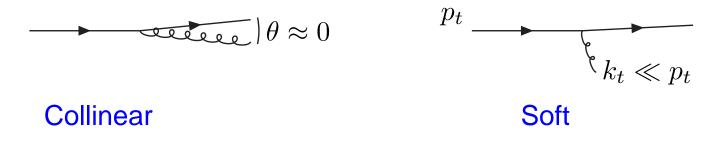
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Two divergences:



For pQCD to make sense, the (hard) jets should not change when

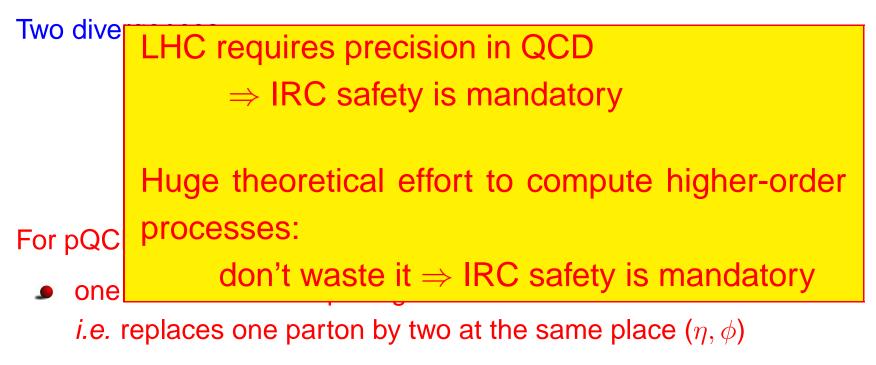
- one has a collinear splitting
 i.e. replaces one parton by two at the same place (η, φ)
- one has a soft emission *i.e.* adds a very soft gluon

QCD divergences



QCD probability for gluon bremsstrahlung at angle θ and \perp -mom. k_t :

$$dP \propto \alpha_s \, \frac{d\theta}{\theta} \, \frac{dk_t}{k_t}$$



one has a soft emission *i.e.* adds a very soft gluon



Recombination:

- k_t algorithm
- Cambridge/Aachen alg.

Cone:

- CDF JetClu
- CDF MidPoint
- D0 (run II) Cone
- PxCone
- ATLAS Cone
- CMS Iterative Cone
- PyCell/CellJet
- GetJet



Many important recent developments between 2006 and 2008:

• Speeding up the k_t and Cam/Aachen algorithms using computational-geometry techniques: $\mathcal{O}(N^3) \rightarrow \mathcal{O}(N \log N)$

[M. Cacciari, G. Salam, 06]

- Existing cones are infrared and/or collinear unsafe
 - Cones with split-merge replaced by SISCone

the only infrared-and-collinear-safe cone geometry technques \Rightarrow no cost in time ($\mathcal{O}(N^3) \rightarrow \mathcal{O}(N^2 \log N)$)

[G. Salam, G.S., 07]

• Cones with progressive removal \longrightarrow anti- k_t

fast recombiniton-type algorithm hard jets are circular as for cones with progressive removal

[M. Cacciari, G. Salam, G.S., 08]



Recombination:

- k_t algorithm
- Cambridge/Aachen alg.
- anti- k_t algorithm \leftarrow

Cone:

- CDF JetClu
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- GetJet
- SISCone

4 available safe algorithms

All part of FastJet

[M.Cacciari,G.Salam,G.S.]



More refined clustering ("2nd generation of algorithms")

Cambridge+Filtering algorithm:

- Cluster with Aachen/Cambridge and radius *R*
- For each jet, recluster it with Aachen/Cambridge and radius R_{sub} keep only n_{sub} hardest sub-jets of the initial jet

Aim: remove the soft background

Properties:

• Proven to improve jet reconstruction, in $H \rightarrow b\bar{b}$

[J.Butterworth, A.Davison, M.Rubin, G.Salam, 08]

- Additional parameters that deserve appropriate studies
- We will use the simplest choice: $R_{sub} = R/2$, $n_{sub} = 2$



Part 1: "best" jet finder in pp collisions

sample processes to study



We analyse 3 processes:

• $Z' \rightarrow q\bar{q} \rightarrow 2$ jets: (ficticious narrow Z')

simple environment: identify 2 jets and reconstruct $M_{Z'}$ source of monochromatic quark jets scale dependence: mass of the Z' between 100 GeV and 4 TeV

- H → gg → 2 jets: (ficticious narrow Higgs)
 simple environment: identify 2 jets and reconstruct M_H source of monochromatic gluon jets
 scale dependence: mass of the Higgs between 100 GeV and 4 TeV
- $t\bar{t} \rightarrow W^+ b W^- \bar{b} \rightarrow q\bar{q} b q\bar{q} \bar{b} \rightarrow 6$ jets:

complex environment: identify 6 jets and reconstruct 2 top balance between reconstruction efficiency and identification

with

- the 5 IRC-safe algorithms: k_t , Cambridge, anti- k_t , SISCone, Cam+filtering
- jet radius varied between 0.1 and 1.5

figure of merit for quality measure



We need a measure of the jet reconstruction efficiency

- Forget about measures related to parton-jet matching
- use the reconstructed mass peak

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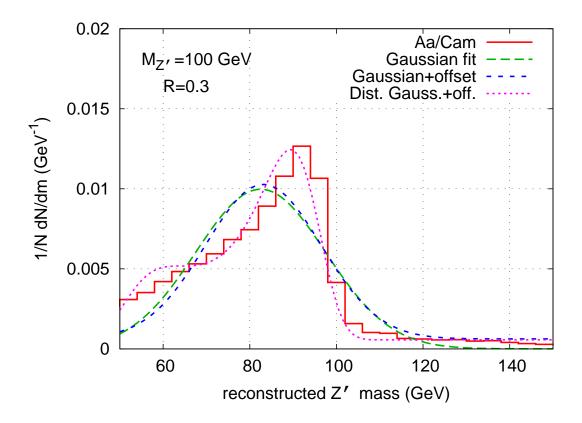


figure of merit for quality measure



We need a measure of the jet reconstruction efficiency

- Forget about measures related to parton-jet matching
- use the reconstructed mass peak
- Forget about fits depending on the shape of the peak
- we shall maximise the signal over background ratio (S/\sqrt{B}) :

 $Q_{f=z}^{w}(JA,R) =$ minimal width of a window containing a fraction f = z of the events



 $Q_{f=z}^{w}(JA, R) =$ minimal width of a window containing a fraction f = z of the events

- it intuitively does what it should
- for a constant background,

$$\frac{Q_{f=z}^{w}(JA_{1},R_{1})}{Q_{f=z}^{w}(JA_{2},R_{2})} = \frac{B_{JA_{1},R_{1}}}{B_{JA_{2},R_{2}}} = \left[\frac{(S/\sqrt{B})_{JA_{1},R_{1}}}{(S/\sqrt{B})_{JA_{2},R_{2}}}\right]^{-2}$$

smaller width \equiv better signal-to-background ratio

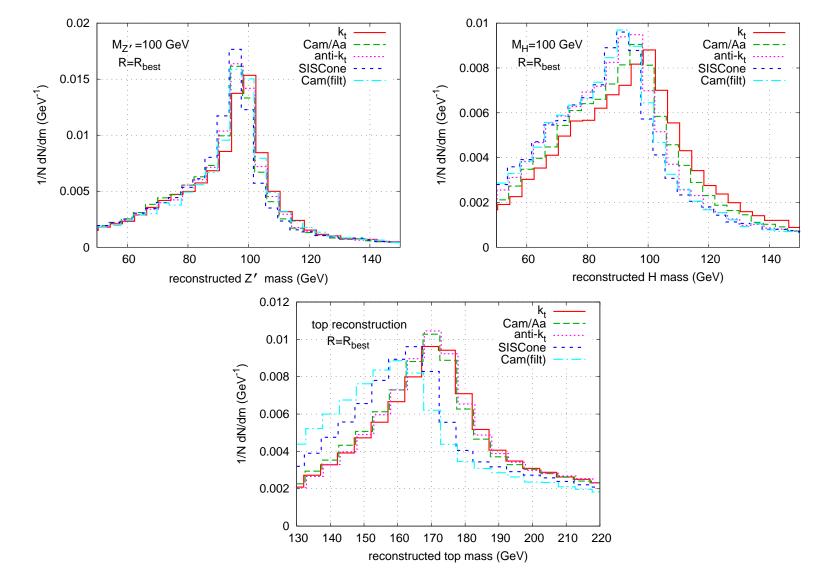
we can associate an effective luminosity ratio

$$\rho_{\mathcal{L}} = \frac{\mathcal{L}_1}{\mathcal{L}_2} = \left[\frac{(S/\sqrt{B})_{JA_1,R_1}}{(S/\sqrt{B})_{JA_2,R_2}}\right]^2 = \frac{Q_{f=z}^w(JA_2,R_2)}{Q_{f=z}^w(JA_1,R_1)}$$

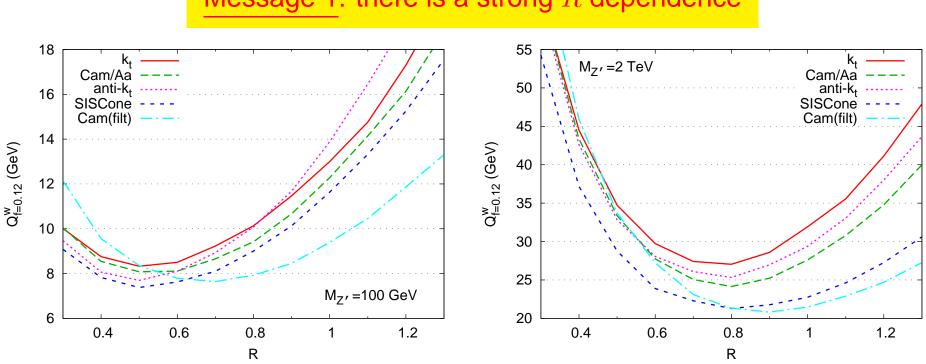
e.g. if $Q_{f=z}^w(JA_2, R_2) = 2Q_{f=z}^w(JA_1, R_1)$, (JA_2, R_2) will need twice the luminosity of (JA_1, R_1) to achieve the same discriminative power.



we see peaks...







Message 1: there is a strong *R* dependence

At 100 GeV,

using R = 0.8 instead of R = 0.5 means a discr. power loss of 20% ($\rho_{\mathcal{L}} \approx 0.8$)

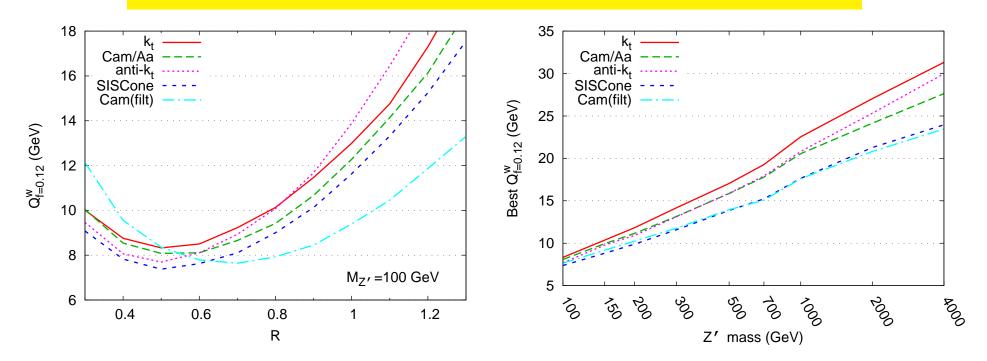
R = 1 40% ($\rho_{\mathcal{L}} \approx 0.6$)

At 2 TeV,

using R = 0.5 instead of R = 0.8 means a discr. power loss of 20% ($\rho_{\mathcal{L}} \approx 0.8$)



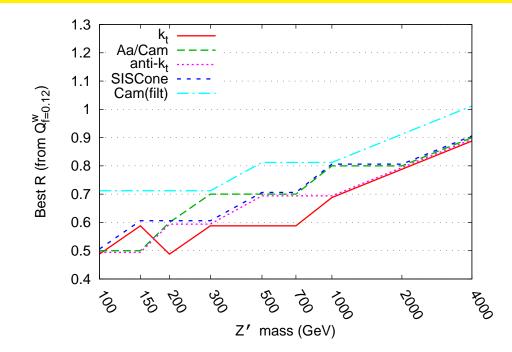
Message 2: SISCone and Cam+filt do a slightly better job



Using k_t instead of SISCone means a discr. power loss of 15% at 100 GeV ($\rho_{\mathcal{L}} \approx 0.85$) 20% at 2 TeV ($\rho_{\mathcal{L}} \approx 0.8$)



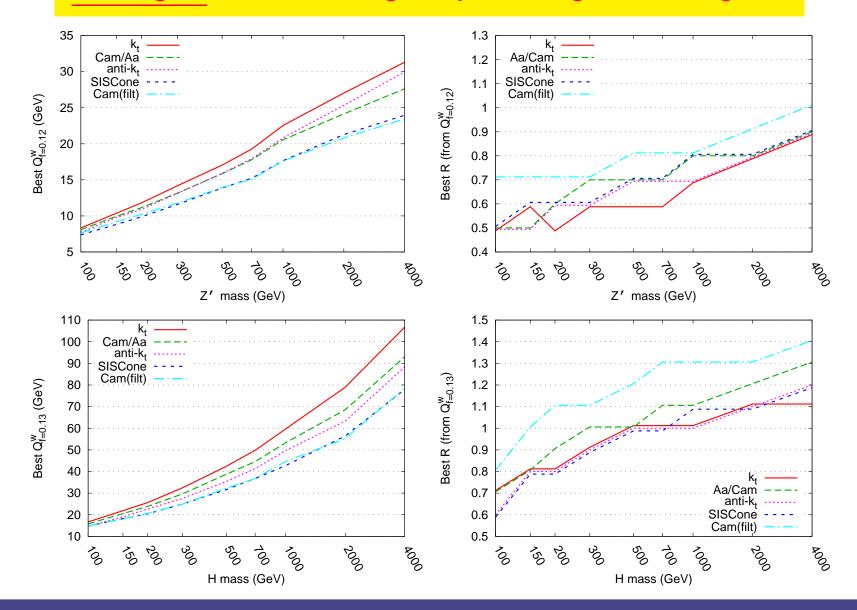
Message 3: The parameters vary with the scale



The preferred value for R increase with the mass scale (typically like log(M))



Message 4: same for the gluon jets, though with a larger R



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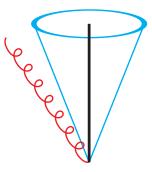
BNL, USA, August 15th 2008

optimising jet-finding – p. 20/34



Competition between

catching perturbative radiation



Out-of-cone radiation:

$$\sim \int_R \frac{d\theta}{\theta} \sim \log(1/R)$$

not catching soft background radiation

Soft contents \sim jet area $\sim R^2$

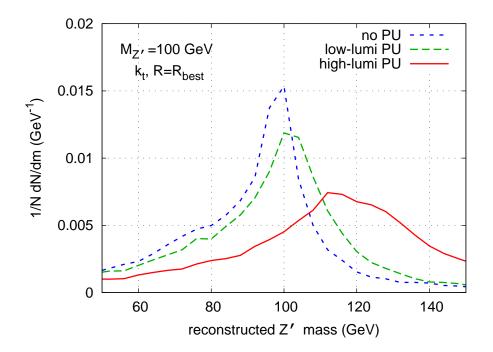
more detailed computation in progress...



Part 2: when pileup enters the game



Pileup \approx uniform soft background that shifts jets to higher p_t



... that needs to be subtracted!





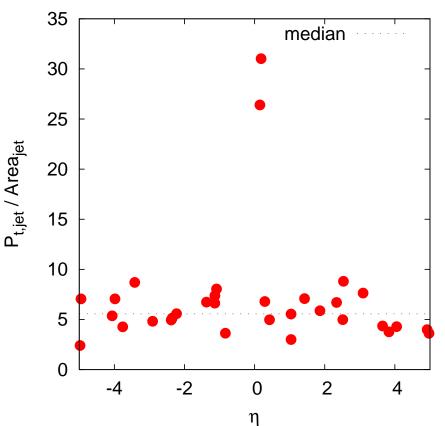
Basic idea:

$$p_{t,\text{subtracted}} = p_{t,\text{jet}} - \rho_{\text{pileup}} \times \text{Area}_{\text{jet}}$$

Jet area:

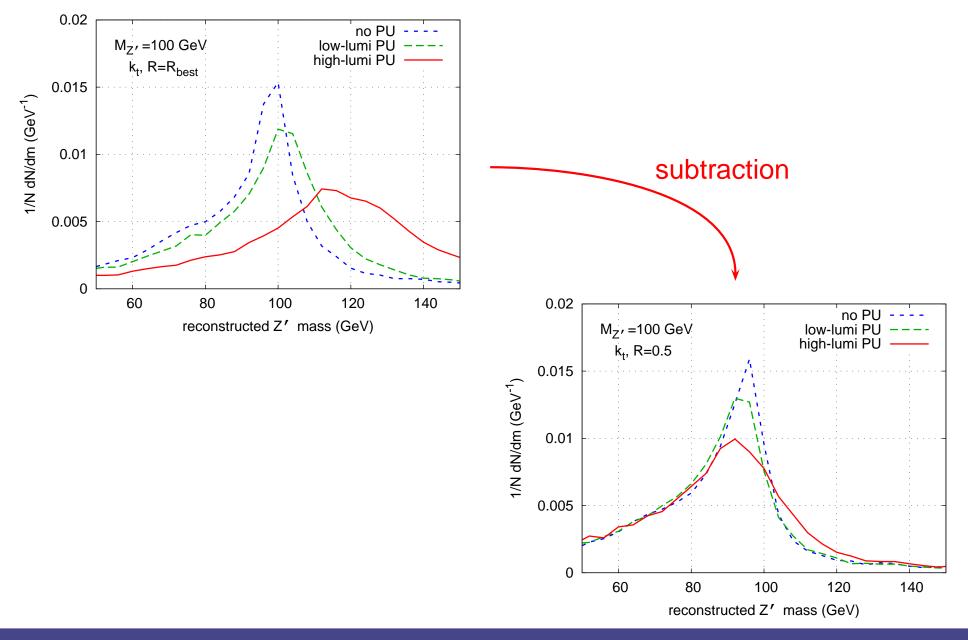
- region where the jet catches infinitely soft particles
- tractable analytically
- Pileup density per unit area: ρ_{pileup} e.g. estimated from the median of $p_{t,\text{jet}}/\text{Area}_{\text{jet}}$

implemented in FastJet on an event-by-event basis



Subtraction recipe





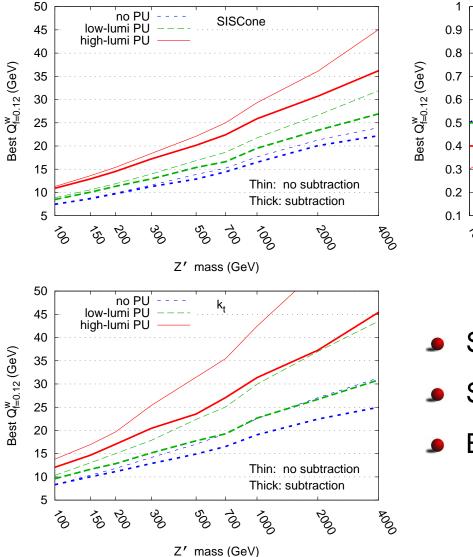
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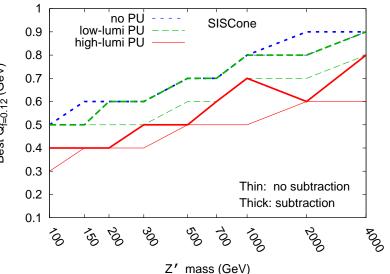
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Message 5: with subtraction, pileup has reasonably little influence





- Subtraction reduces width
- SISCone a bit better than k_t
- Best R not much affected

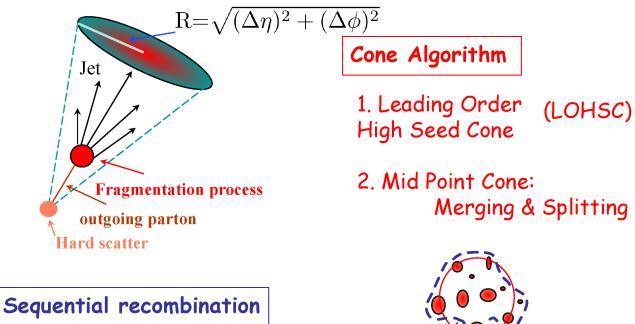


Part 3: heavy-ion background: subtracting more complex background



[S. Salur, J. Putschke, ... (STAR)]





KT
 Cambridge/ Aachen

 K_{T} jet Cone jet

Explore systematics: Use both Clustering & Cone algorithms.

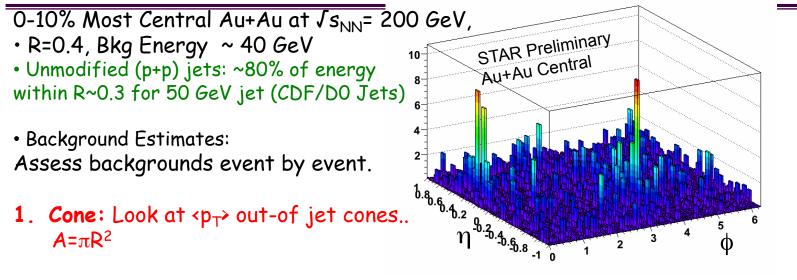
Sevil Salur

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[S. Salur, J. Putschke, ... (STAR)]

Correction for Heavy-Ion Background



2. Sequential Recombination: Estimate the active area of each jet by addition of zero energy particles of known density.

 p_T (Jet Measured) ~ p_T (Parton) + $\rho \times A(Jet) \pm \sigma \sqrt{A(Jet)}$

 ρ = Diffuse noise, σ =noise fluctuations M. Cacciari, G. Salam, G. Soyez 0802.1188 [hep-ph]

Reduction of background fluctuations: p_T cuts, limit R.

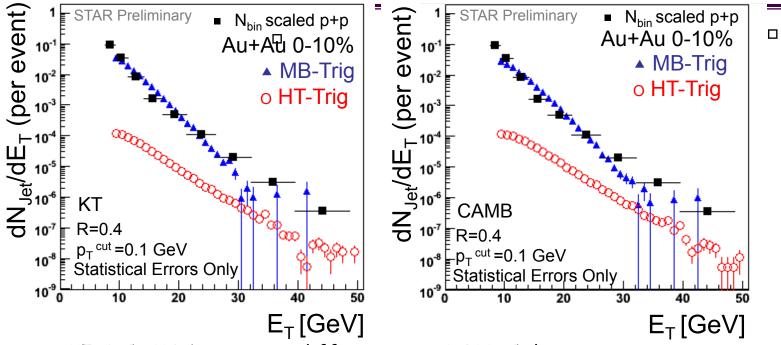
Sevil Salur

Jets at RHIC



[S. Salur, J. Putschke, ... (STAR)]

Jets with Sequential Recombination Algorithm



KT & CAMB biases are different wrt. LOHSC due to:

- -- background subtraction algorithm
- -- no seed
- -- low pt cut

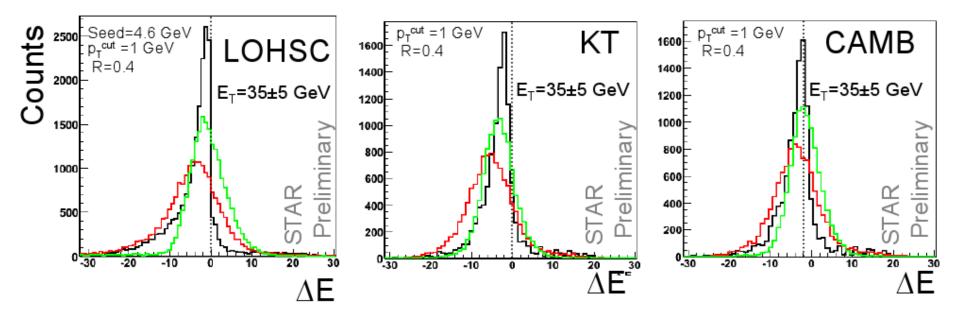
Systematic Uncertainity on Normalization: 50% Good agreement with N_{bin} scaled p+p for unbiased algorithms.

Sevil Salur

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[S. Salur, J. Putschke, ... (STAR)]



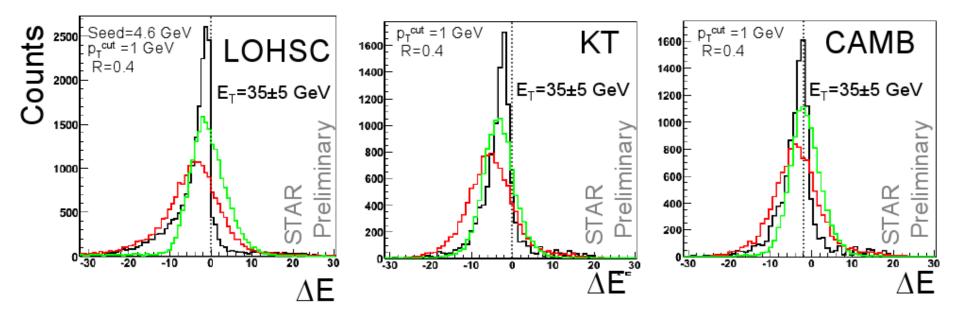
Red curves: $\Delta E = E^{PyEmbed} - E^{PyDet}$ with PyDet \equiv Pythia pp + detector effects

 $PyEmbed \equiv same with real AuAu event added$

 \Rightarrow measure of the subtraction efficiency



[S. Salur, J. Putschke, ... (STAR)]



Red curves: $\Delta E = E^{\text{PyEmbed}} - E^{\text{PyDet}}$

with $PyDet \equiv Pythia \ pp$ + detector effects

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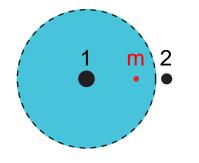
 \Rightarrow measure of the subtraction efficiency

Work under progress: removing the last few GeV shift in ΔE

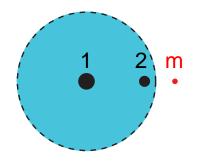


Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
- Modify the hard scattering (back-reaction)
 - can be pointlike or diffuse
 - gain: p_2 gained when adding p_m



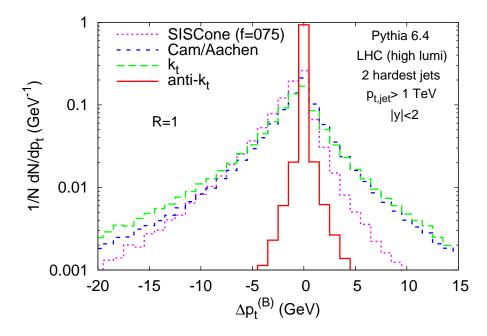
• loss: p_2 lost when adding p_m





Additional soft background has 2 effects:

- Throw soft particles in the hard jet: dealt with by subtraction
- Modify the hard scattering (back-reaction)
 - can be pointlike or diffuse
 - tractable analytically (similar to areas)
 - $k_t \gtrsim \text{Cambridge} > \text{SISCone} \gg \text{anti-}k_t$





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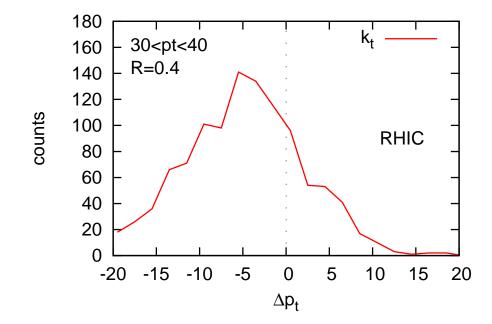
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 - $k_t \gtrsim \text{Cambridge} > \text{SISCone} \gg \text{anti-}k_t$

+ For heavy-ion collisions: fluctuating underlying event background \rightarrow median estimation of ρ might oversimplified

Results (3)



Work under progress: test subtraction using Pythia pp hard event + HYDJET AA background

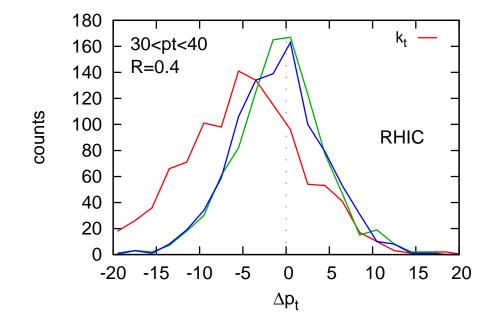


Similar shifts than STAR

Results (3)



Work under progress: test subtraction using Pythia pp hard event + HYDJET AA background



Target: reach a precision of 1-2 GeV by

- carefully tuning the algorithm to reduce UE sensitivity and back-reaction
- carefully tuning the subtraction to deal with the fluctuating background



Use IRC-safe algorithms

Jet-finding in pp at the LHC

- SISCone and Cam+filt. do a slightly better job
- strong R dependence: important to choose R_{best}
- R_{best} increases with the scale
- same for quark and gluon jets, larger R_{best} for gluons
- with subtraction, pileup has reasonably little influence
- \Rightarrow flexibility in jet physics at the LHC
- Jet-finding in AA at RHIC and the LHC
 - First measurement at RHIC
 - Work under progress: improve subtraction down to 1-2 GeV