

Brane decay in curved space-time

DAN ISRAËL, IAP, UNIV. PIERRE ET MARIE CURIE

From D. I. & E. Rabinovici, JHEP **0701** (2007)
D. I., JHEP **0704** (2007)

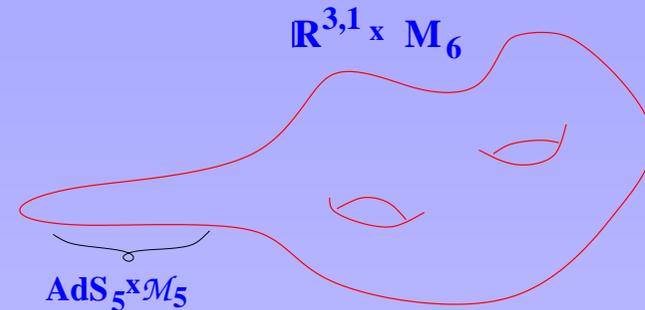
Outline of the Talk

1. Introduction & motivations from cosmology
2. Brane decay in flat space-time
3. Brane decay in AdS
4. Closed and open string emission
5. Open string tachyons in non-critical strings (geometric and universal)
6. Brane decay in non-critical strings
7. Conclusions

① *Brane inflation & open string tachyons*

✓ **Natural setting of string cosmology:** flux compactification of type II string theory, with stabilized moduli

➔ generically warped throats develop



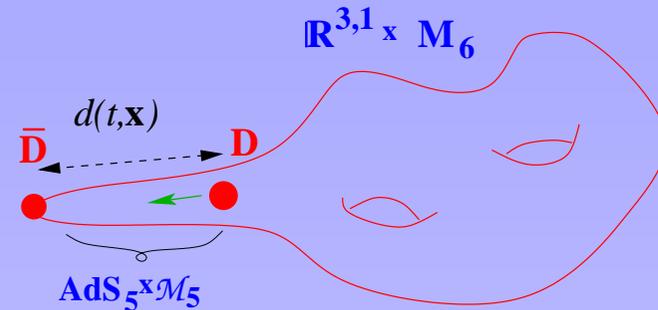
✓ AdS_5 geometry, capped both in the UV (compact 6-manifold) and in the IR (tip of the throat)

[Giddings, Kachru, Polchinski '03]

① Brane inflation & open string tachyons

✓ Natural setting of string cosmology: flux compactification of type II string theory, with stabilized moduli

➔ generically warped throats develop



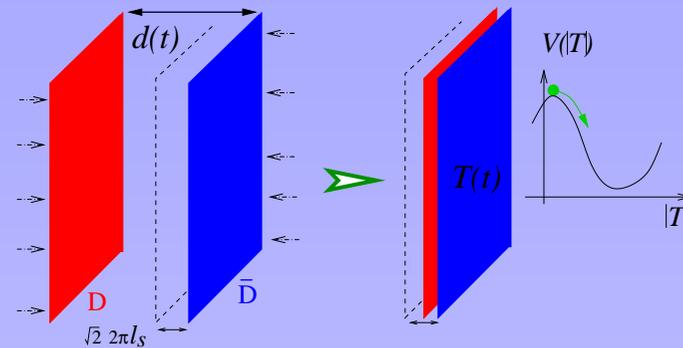
✓ AdS_5 geometry, capped both in the UV (compact 6-manifold) and in the IR (tip of the throat)

[Giddings, Kachru, Polchinski '03]

✓ D-brane/ anti D-brane pair in the throat: Coulombian attraction redshifted by AdS_5 metric ➔ slow-roll inflation (inflaton $d(t, \mathbf{x})$) [Kachru, Kallosh, Linde, Maldacena, McAllister, Trivedi '03]

Tachyon condensation & Brane Reheating

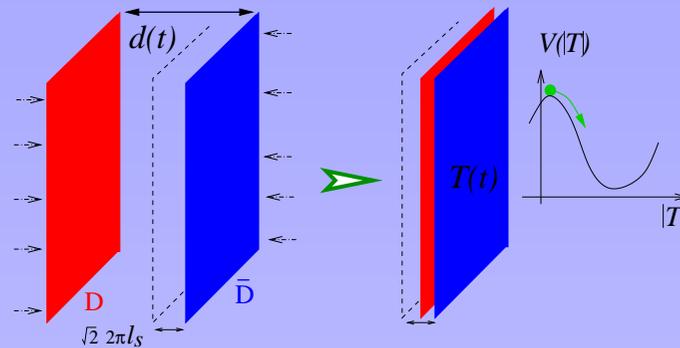
✓ End of inflation: D- \bar{D} annihilation \rightarrow open string tachyon for $d^2 < 8\pi^2\ell_s^2$



★ String theory realization of hybrid inflation

Tachyon condensation & Brane Reheating

- ✓ End of inflation: D- \bar{D} annihilation \rightarrow open string tachyon for $d^2 < 8\pi^2\ell_s^2$



- ★ String theory realization of hybrid inflation

- ✓ Tachyon condensation: involves all the massive string modes ($m > 1/\ell_s$)
 - \rightarrow string corrections important

★ For brane decay in flat spacetime one can use exact tree-level string computations

[Sen '02]

★ For brane decay in flat spacetime one can use exact tree-level string computations

[Sen '02]

↳ "Rolling tachyon" solution of boundary CFT

★ For brane decay in flat spacetime one can use exact tree-level string computations

[Sen '02]

↳ "Rolling tachyon" solution of boundary CFT

↳ produces non-relativistic "tachyon dust" of massive closed strings

★ For brane decay in flat spacetime one can use **exact tree-level string computations**

[Sen '02]

↳ "Rolling tachyon" solution of boundary CFT

↳ produces non-relativistic "tachyon dust" of massive closed strings

★ How much this elementary process of string theory is **universal**?

★ For brane decay in flat spacetime one can use **exact tree-level string computations**

[Sen '02]

↳ "Rolling tachyon" solution of boundary CFT

↳ produces non-relativistic "tachyon dust" of massive closed strings

★ How much this elementary process of string theory is **universal**?

✓ Both low-energy (background geometry) and high energy data (asymptotic density of string states) play a role in the computation of closed string radiation.

✓ both regimes are related through modular invariance

★ For brane decay in flat spacetime one can use **exact tree-level string computations** *[Sen '02]*

↳ "Rolling tachyon" solution of boundary CFT

↳ produces non-relativistic "tachyon dust" of massive closed strings

★ How much this elementary process of string theory is **universal**?

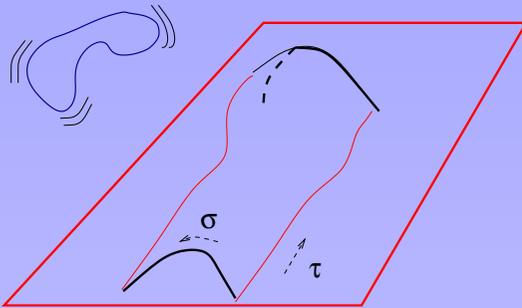
✓ Both low-energy (background geometry) and high energy data (asymptotic density of string states) play a role in the computation of closed string radiation.

✓ both regimes are related through modular invariance

↳ look for curved spacetimes (especially warped geometries) where the D-brane decay can be followed exactly (at the perturbative level)

② Brane Annihilation: Flat Space-Time

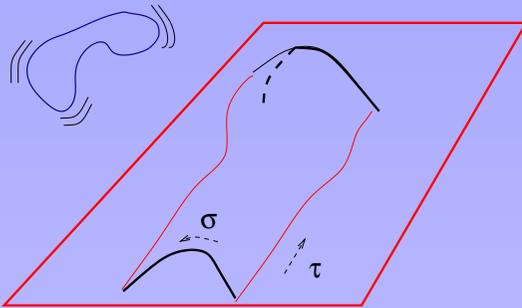
✓ Decay of an unstable D-brane: equivalent to coincident $D-\bar{D}$ pair with no relative velocity (using $(-)^{FL}$ orbifold) → solvable worldsheet string model [Sen '02]



$$\delta S = \lambda \int d\tau \exp\{X^0(\tau)/\ell_s\} \text{ Wick rotation of boundary Liouville theory}$$

② Brane Annihilation: Flat Space-Time

- ✓ Decay of an unstable D-brane: equivalent to coincident $D-\bar{D}$ pair with no relative velocity (using $(-)^{FL}$ orbifold) → solvable worldsheet string model [Sen '02]

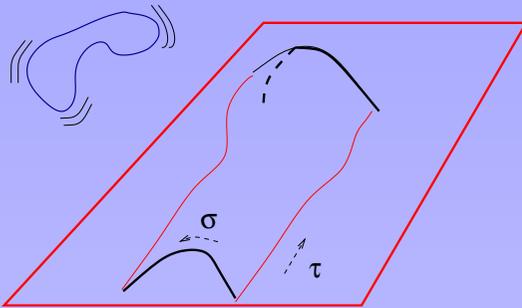


$$\delta S = \lambda \int d\tau \exp\{X^0(\tau)/\ell_s\} \text{ Wick rotation of boundary Liouville theory}$$

- ✓ Couplings to closed strings (*grav. sector*) $\langle V_E \rangle_\lambda = (\pi\lambda)^{-iE} \frac{\pi}{\sinh \pi E}$
- time-dependent source for all closed string modes

② Brane Annihilation: Flat Space-Time

- ✓ Decay of an unstable D-brane: equivalent to coincident $D-\bar{D}$ pair with no relative velocity (using $(-)^{FL}$ orbifold) → solvable worldsheet string model [Sen '02]



$$\delta S = \lambda \int d\tau \exp\{X^0(\tau)/\ell_s\} \text{ Wick rotation of boundary Liouville theory}$$

- ✓ Couplings to closed strings (*grav. sector*) $\langle V_E \rangle_\lambda = (\pi\lambda)^{-iE} \frac{\pi}{\sinh \pi E}$
 → time-dependent source for all closed string modes

- ★ Closed strings production (*coherent state*)

Number of emitted strings (*tree-level*): $\mathcal{N} = \int \frac{dE}{2E} \rho(E) |\langle V_E \rangle_\lambda|^2$ [Lambert, Liu, Maldacena '03]

✓ Density of closed strings oscillators $\rho(N)$

➔ exponentially growing (*cf. Hagedorn transition at high temperature*)

★ In flat space-time, $\rho(N) \sim N^\alpha e^{+4\pi\sqrt{N}}$ with $E = 2\sqrt{N}/\ell_s$

✓ Density of closed strings oscillators $\rho(N)$

➔ exponentially growing (cf. Hagedorn transition at high temperature)

★ In flat space-time, $\rho(N) \sim N^\alpha e^{+4\pi\sqrt{N}}$ with $E = 2\sqrt{N}/\ell_s$

✓ Amplitude $\mathcal{N} \sim \int dE E^{2\alpha-1} e^{2\pi E} \sinh^{-2}(\pi E)$

➔ divergent for D0-branes ($\alpha = 0$) (D3-branes: instable to inhomogeneous decay)

✓ Density of closed strings oscillators $\rho(N)$

➔ exponentially growing (*cf. Hagedorn transition at high temperature*)

★ In flat space-time, $\rho(N) \sim N^\alpha e^{+4\pi\sqrt{N}}$ with $E = 2\sqrt{N}/\ell_s$

✓ Amplitude $\mathcal{N} \sim \int dE E^{2\alpha-1} e^{2\pi E} \sinh^{-2}(\pi E)$

➔ divergent for D0-branes ($\alpha = 0$) (*D3-branes: instable to inhomogeneous decay*)

★ Divergence signals **breakdown of string perturbation theory**

➔ Large gravitational back-reaction from the brane decay!

✓ Density of closed strings oscillators $\rho(N)$

➔ exponentially growing (*cf. Hagedorn transition at high temperature*)

★ In flat space-time, $\rho(N) \sim N^\alpha e^{+4\pi\sqrt{N}}$ with $E = 2\sqrt{N}/\ell_s$

✓ Amplitude $\mathcal{N} \sim \int dE E^{2\alpha-1} e^{2\pi E} \sinh^{-2}(\pi E)$

➔ divergent for D0-branes ($\alpha = 0$) (*D3-branes: instable to inhomogeneous decay*)

★ Divergence signals **breakdown of string perturbation theory**

➔ Large gravitational back-reaction from the brane decay!

★ mass of a D0-brane $m_{D0} \propto 1/\ell_s g_s$

➔ energy conservation not "built-in" the (tree-level) computation

✓ Density of closed strings oscillators $\rho(N)$

➔ exponentially growing (*cf. Hagedorn transition at high temperature*)

★ In flat space-time, $\rho(N) \sim N^\alpha e^{+4\pi\sqrt{N}}$ with $E = 2\sqrt{N}/\ell_s$

✓ Amplitude $\mathcal{N} \sim \int dE E^{2\alpha-1} e^{2\pi E} \sinh^{-2}(\pi E)$

➔ divergent for D0-branes ($\alpha = 0$) (*D3-branes: instable to inhomogeneous decay*)

★ Divergence signals **breakdown of string perturbation theory**

➔ Large gravitational back-reaction from the brane decay!

★ mass of a D0-brane $m_{D0} \propto 1/\ell_s g_s$

➔ energy conservation not "built-in" the (tree-level) computation

✓ One needs a **UV cutoff** at $E \sim m_{D0}$

★ fraction of total energy in strings of mass $m \sim \text{cst.}$ (up to m_{D0})

➔ most energy in strings $m \sim m_{D0}$, non-relativistic ($p \propto 1/\ell_s \sqrt{g_s}$): **tachyon dust**

Sen's Conjecture

1. The closed string description of the brane decay breaks down after $t \sim \ell_s \sqrt{g_s}$
→ all energy is converted into *tachyon dust* of massive closed strings

Sen's Conjecture

1. The closed string description of the brane decay breaks down after $t \sim \ell_s \sqrt{g_s}$
→ all energy is converted into *tachyon dust* of massive closed strings
2. However the **open string description** of the process remains valid
→ may be spoiled by open string pair production (more later)

Sen's Conjecture

1. The closed string description of the brane decay breaks down after $t \sim \ell_s \sqrt{g_s}$
→ all energy is converted into *tachyon dust* of massive closed strings
2. However the **open string description** of the process remains valid
→ may be spoiled by open string pair production (more later)
3. The open string description is *holographically dual* to the closed strings description, hence is *complete*

Sen's Conjecture

1. The closed string description of the brane decay breaks down after $t \sim \ell_s \sqrt{g_s}$
→ all energy is converted into *tachyon dust* of massive closed strings
2. However the **open string description** of the process remains valid
→ may be spoiled by open string pair production (more later)
3. The open string description is *holographically dual* to the closed strings description, hence is *complete*
4. One can use the tachyon low-energy effective action
$$S_T = \int d^d x \cosh(T/\sqrt{2})^{-1} \sqrt{-\det(\eta_{\mu\nu} + \partial_\mu T \partial_\nu T + \dots)}$$
 → late-time "dust"

Sen's Conjecture

1. The closed string description of the brane decay breaks down after $t \sim \ell_s \sqrt{g_s}$
→ all energy is converted into *tachyon dust* of massive closed strings
2. However the **open string description** of the process remains valid
→ may be spoiled by open string pair production (more later)
3. The open string description is *holographically dual* to the closed strings description, hence is *complete*
4. One can use the tachyon low-energy effective action
$$S_T = \int d^d x \cosh(T/\sqrt{2})^{-1} \sqrt{-\det(\eta_{\mu\nu} + \partial_\mu T \partial_\nu T + \dots)}$$
 → late-time "dust"
5. Conjecture has been checked in 2D string theory

Is this an universal process?

- ✓ **Cosmological context:** D/\bar{D} in a curved space-time (*e.g. capped AdS_5*)
 - ➔ is the physics of the decay similar? (in string theory, UV-IR relation)

Is this an universal process?

- ✓ **Cosmological context:** D/\bar{D} in a curved space-time (*e.g.* capped AdS_5)
 - ➔ is the physics of the decay similar? (in string theory, UV-IR relation)
- ✓ In particular cancellation between **asymptotic density of closed string states** & **closed string emission amplitude** may not be true anymore
- ★ In a CFT with minimal conformal dimension Δ_m , $\rho(E) \sim \exp\{\sqrt{1 - \Delta_m} 2\pi E\}$
 - UV finite?

Is this an universal process?

- ✓ **Cosmological context:** D/ \bar{D} in a curved space-time (e.g. capped AdS_5)
 - ➔ is the physics of the decay similar? (in string theory, UV-IR relation)

 - ✓ In particular cancellation between **asymptotic density of closed string states** & **closed string emission amplitude** may not be true anymore
 - ★ In a CFT with minimal conformal dimension Δ_m , $\rho(E) \sim \exp\{\sqrt{1 - \Delta_m} 2\pi E\}$
 - UV finite?

 - ✓ Is the process still well-described by the curved background generalization of the **open string tachyon effective action**?
- $$S_T = \int d^{p+1}x \sqrt{-g} \cosh\left(\frac{T}{\sqrt{2}}\right)^{-1} \sqrt{-\det\{(g + B + 2\pi\ell_s^2 F)_{\mu\nu} + \partial_\mu T \partial_\nu T\}} + \int W(T) dT \wedge C_{[p]}$$

Is this an universal process?

- ✓ **Cosmological context:** D/ \bar{D} in a curved space-time (e.g. capped AdS_5)
 - ➔ is the physics of the decay similar? (in string theory, UV-IR relation)

- ✓ In particular cancellation between **asymptotic density of closed string states** & **closed string emission amplitude** may not be true anymore
- ★ In a CFT with minimal conformal dimension Δ_m , $\rho(E) \sim \exp\{\sqrt{1 - \Delta_m} 2\pi E\}$
 - **UV finite?**

- ✓ Is the process still well-described by the curved background generalization of the **open string tachyon effective action**?

$$S_T = \int d^{p+1}x \sqrt{-g} \cosh\left(\frac{T}{\sqrt{2}}\right)^{-1} \sqrt{-\det\{(g + B + 2\pi\ell_s^2 F)_{\mu\nu} + \partial_\mu T \partial_\nu T\}} + \int W(T) dT \wedge C_{[p]}$$
- ★ In particular, if all the brane energy is not radiated into massive closed strings, the whole picture may be challenged

③ *Decay in Curved Space (I): Anti-de Sitter*

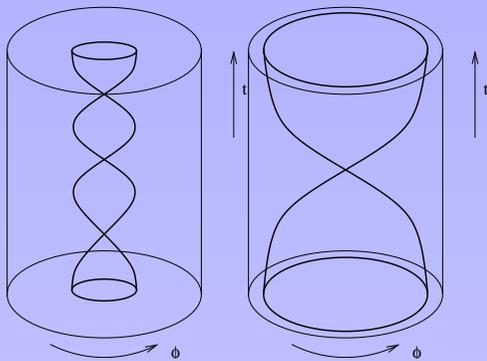
✓ Despite recent progress AdS_5 string theory not solvable

③ Decay in Curved Space (I): Anti-de Sitter

- ✓ Despite recent progress AdS₅ string theory not solvable
- ✓ Solvable "toy model": three-dimensional AdS → conformal field theory on the string worldsheet: Wess-Zumino Witten model for the group manifold $SL(2, \mathbb{R})$
 $ds^2 = \ell_s^2 k [d\rho^2 + \sinh^2 \rho d\phi^2 - \cosh^2 \rho d\tau^2]$, with a B-field $B = \ell_s^2 k \cosh 2\rho d\tau \wedge d\phi$

③ Decay in Curved Space (I): Anti-de Sitter

- ✓ Despite recent progress AdS₅ string theory not solvable
 - ✓ Solvable "toy model": three-dimensional AdS → conformal field theory on the string worldsheet: Wess-Zumino Witten model for the group manifold $SL(2, \mathbb{R})$
- $$ds^2 = \ell_s^2 k [d\rho^2 + \sinh^2 \rho d\phi^2 - \cosh^2 \rho d\tau^2], \text{ with a B-field } B = \ell_s^2 k \cosh 2\rho d\tau \wedge d\phi$$



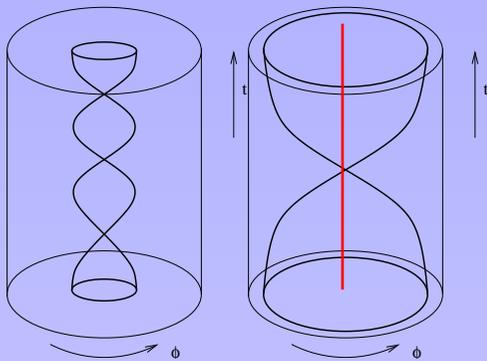
Two types of string modes:

short strings trapped in AdS (exponentially decreasing wave-functions)

long strings, macroscopic solutions winding w -times around ϕ

③ Decay in Curved Space (I): Anti-de Sitter

- ✓ Despite recent progress AdS₅ string theory not solvable
- ✓ Solvable "toy model": three-dimensional AdS → conformal field theory on the string worldsheet: Wess-Zumino Witten model for the group manifold SL(2,ℝ)
 $ds^2 = \ell_s^2 k [d\rho^2 + \sinh^2 \rho d\phi^2 - \cosh^2 \rho d\tau^2]$, with a B-field $B = \ell_s^2 k \cosh 2\rho d\tau \wedge d\phi$



Two types of string modes:
short strings trapped in AdS (exponentially decreasing wave-functions)
long strings, macroscopic solutions winding w -times around ϕ

- ✓ Unstable D0-brane of type IIB superstrings in AdS₃ × M₇: localized at the origin $\rho = 0$ (infrared) → decay of the brane solvable (equivalent to D- \bar{D} annihilation)

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$
- ➔ writing $SL(2, \mathbb{R}) \sim SL(2, \mathbb{R})/U(1) \times \mathbb{R}^{0,1}$ the dynamics involves only the time direction (up to an orbifold)

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$
- ➔ writing $SL(2, \mathbb{R}) \sim SL(2, \mathbb{R})/U(1) \times \mathbb{R}^{0,1}$ the dynamics involves only the time direction (up to an orbifold)
- ➔ decay given by same deformation as in flat space $\lambda \int_{\partial\Sigma} dt \mathbb{I} \times \exp\{\sqrt{k/2} \tau(t)\}$

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$
- ➔ writing $SL(2, \mathbb{R}) \sim SL(2, \mathbb{R})/U(1) \times \mathbb{R}^{0,1}$ the dynamics involves only the time direction (up to an orbifold)
- ➔ decay given by same deformation as in flat space $\lambda \int_{\partial\Sigma} dt \mathbb{I} \times \exp\{\sqrt{k/2} \tau(t)\}$
- ★ One gets the **D-brane couplings to closed string modes**, e.g. for long strings with radial momentum p_ρ and winding w :

$$|\langle V_{p_\rho, w, E} \rangle_\lambda| \propto \sqrt{\frac{\sinh 2\pi p_\rho \sinh \frac{2\pi p_\rho}{k}}{\cosh 2\pi \rho + \cos \pi(E - kw)}} \frac{1}{|\sinh \frac{\pi E}{\sqrt{2k}}|} \text{ with } E = \frac{k w}{2} + \frac{2}{w} \left[\frac{p_\rho^2 + \frac{1}{4}}{k} + N + \dots \right]$$

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$
- ➔ writing $SL(2, \mathbb{R}) \sim SL(2, \mathbb{R})/U(1) \times \mathbb{R}^{0,1}$ the dynamics involves only the time direction (up to an orbifold)
- ➔ decay given by same deformation as in flat space $\lambda \int_{\partial\Sigma} dt \mathbb{I} \times \exp\{\sqrt{k/2} \tau(t)\}$
- ★ One gets the **D-brane couplings to closed string modes**, e.g. for long strings with radial momentum p_ρ and winding w :

$$|\langle V_{p_\rho, w, E} \rangle_\lambda| \propto \sqrt{\frac{\sinh 2\pi p_\rho \sinh \frac{2\pi p_\rho}{k}}{\cosh 2\pi \rho + \cos \pi(E - kw)}} \frac{1}{|\sinh \frac{\pi E}{\sqrt{2k}}|} \text{ with } E = \frac{k w}{2} + \frac{2}{w} \left[\frac{p_\rho^2 + \frac{1}{4}}{k} + N + \dots \right]$$

- ➔ also coupling to discrete states (i.e. localized strings)

Closed Strings Emission by the brane decay

- ✓ Open string sector on the D0-brane: tachyon + tower of string modes built on the *identity representation* of $SL(2, \mathbb{R})$
- ➡ writing $SL(2, \mathbb{R}) \sim SL(2, \mathbb{R})/U(1) \times \mathbb{R}^{0,1}$ the dynamics involves only the time direction (up to an orbifold)
- ➡ decay given by same deformation as in flat space $\lambda \int_{\partial\Sigma} dt \mathbb{I} \times \exp\{\sqrt{k/2} \tau(t)\}$
- ★ One gets the **D-brane couplings to closed string modes**, e.g. for long strings with radial momentum p_ρ and winding w :

$$|\langle V_{p_\rho, w, E} \rangle_\lambda| \propto \sqrt{\frac{\sinh 2\pi p_\rho \sinh \frac{2\pi p_\rho}{k}}{\cosh 2\pi \rho + \cos \pi(E - kw)}} \frac{1}{|\sinh \frac{\pi E}{\sqrt{2k}}|} \text{ with } E = \frac{k w}{2} + \frac{2}{w} \left[\frac{p_\rho^2 + \frac{1}{4}}{k} + N + \dots \right]$$

- ➡ also coupling to discrete states (i.e. localized strings)
- ★ Total number of emitted closed strings given by the **imaginary part** of the annulus one-loop amplitude, using *optical theorem + open/closed channel duality*

$$\mathcal{N} = \text{Im} \left[\int \frac{ds}{2s} \text{Tr}_{\text{open}} e^{-\pi s \mathcal{H}} \right]$$

✓ As in flat space, an important input is the asymptotic density of string states

$$\star E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\left\{2\pi \sqrt{\left(1 - \frac{1}{2k}\right)wE}\right\} \text{ (while } |\langle V_E \rangle|^2 \sim \exp\left\{-\sqrt{\frac{2}{k}}\pi E\right\}\text{)}$$

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

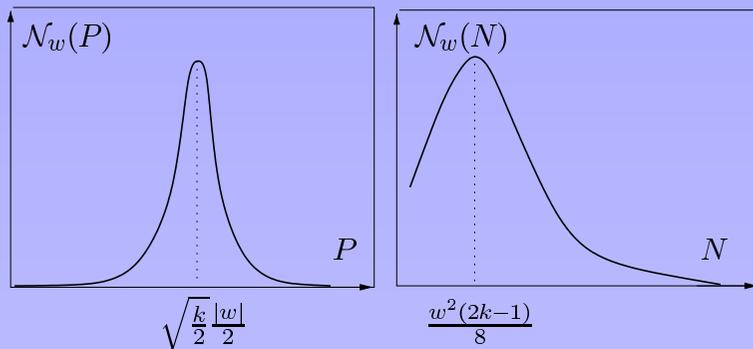
★ Like a 2D field theory (cf. AdS₃/CFT₂)

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\left\{2\pi \sqrt{\left(1 - \frac{1}{2k}\right)wE}\right\}$ (while $|\langle V_E \rangle|^2 \sim \exp\left\{-\sqrt{\frac{2}{k}}\pi E\right\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!

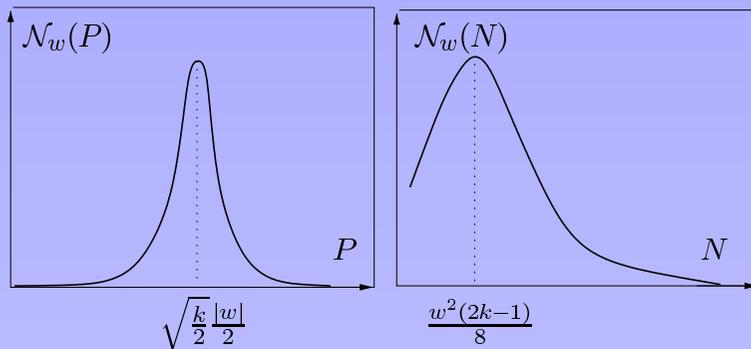


✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!



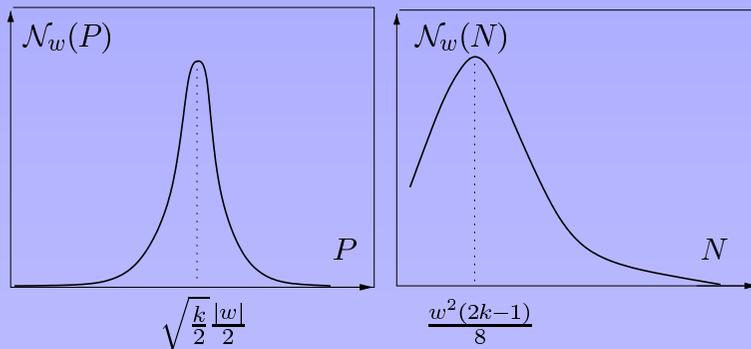
- Displacement of $\langle p_\rho \rangle$ due to non-perturbative corrections in ℓ_s^2 (worldsheet instantons)
 - not seen in semi-classical limit

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!



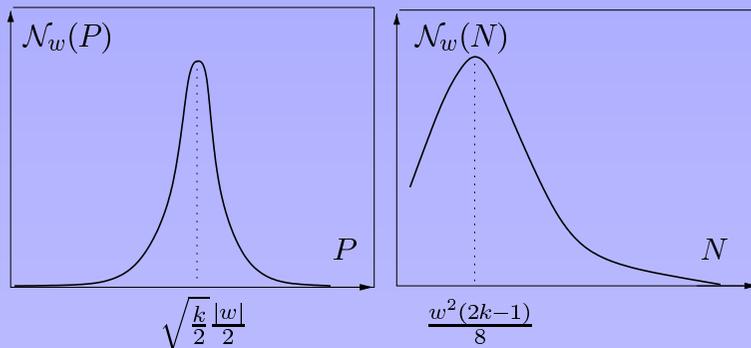
- Displacement of $\langle p_\rho \rangle$ due to non-perturbative corrections in ℓ_s^2 (worldsheet instantons)
 - not seen in semi-classical limit
- For large w , $\langle E \rangle \sim kw$

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!



- Displacement of $\langle p_\rho \rangle$ due to non-perturbative corrections in ℓ_s^2 (worldsheet instantons)
 - not seen in semi-classical limit
- For large w , $\langle E \rangle \sim kw$

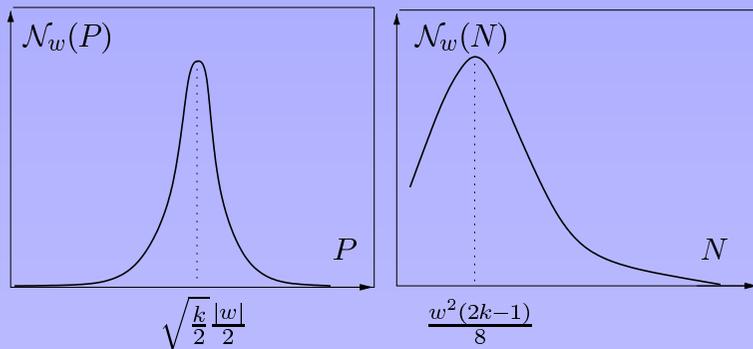
✓ Summation over spectral flow: $\mathcal{N}_{\text{long}} \sim \sum_{w=1}^{\infty} 1/w \rightarrow$ divergence at high energies

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!



- Displacement of $\langle p_\rho \rangle$ due to non-perturbative corrections in ℓ_s^2 (worldsheet instantons)
 - not seen in semi-classical limit
- For large w , $\langle E \rangle \sim kw$

✓ Summation over spectral flow: $\mathcal{N}_{\text{long}} \sim \sum_{w=1}^{\infty} 1/w \rightarrow$ divergence at high energies

★ Needs non-perturbative UV cutoff:

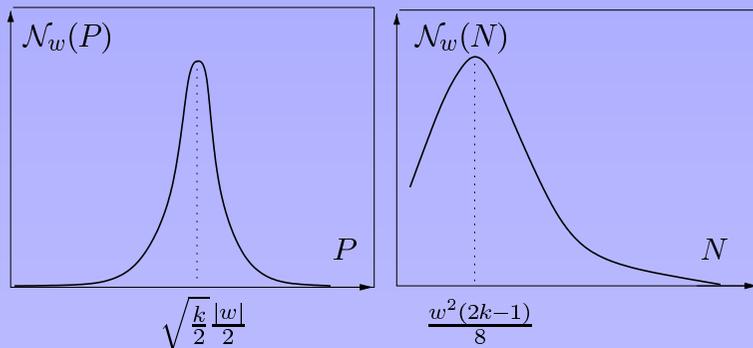
- $w \lesssim 1/g_s^2$ (NS-NS charge conservation)
- $w \lesssim 1/g_s$ (energy conservation)

✓ As in flat space, an important input is the asymptotic density of string states

★ $E \sim \frac{2N}{w} \rightarrow \rho(E) \sim E^\alpha \exp\{2\pi \sqrt{(1 - \frac{1}{2k})wE}\}$ (while $|\langle V_E \rangle|^2 \sim \exp\{-\sqrt{\frac{2}{k}}\pi E\}$)

★ Like a 2D field theory (cf. AdS₃/CFT₂)

★ for given winding w , long strings emission is (exponentially) UV-finite!



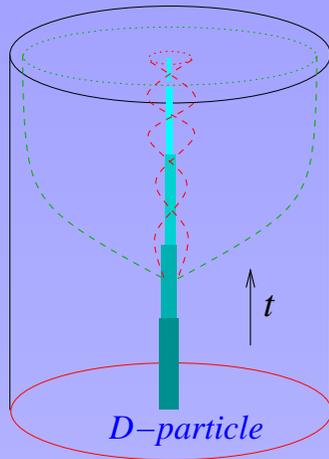
- Displacement of $\langle p_\rho \rangle$ due to non-perturbative corrections in ℓ_s^2 (worldsheet instantons)
 - not seen in semi-classical limit
- For large w , $\langle E \rangle \sim kw$

✓ Summation over spectral flow: $\mathcal{N}_{\text{long}} \sim \sum_{w=1}^{\infty} 1/w \rightarrow$ divergence at high energies

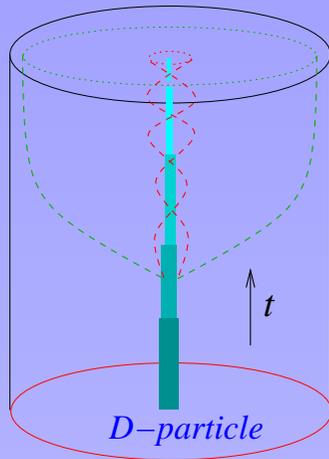
★ Needs non-perturbative UV cutoff:

- $w \lesssim 1/g_s^2$ (NS-NS charge conservation)
- $w \lesssim 1/g_s$ (energy conservation)

★ On the contrary, emission of short strings (localized strings) stays finite

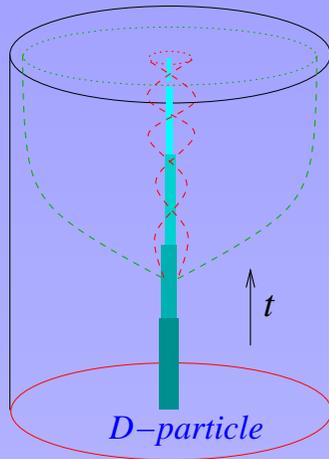


✓ **Conclusion:** most of the energy converted into highly excited long strings of winding $w \sim 1/g_s$, expanding at speed $d\rho/dt \sim 1/\ell_s\sqrt{k}$



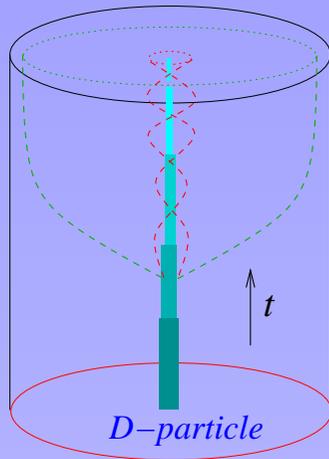
✓ **Conclusion:** most of the energy converted into highly excited long strings of winding $w \sim 1/g_s$, expanding at speed $d\rho/dt \sim 1/\ell_s\sqrt{k}$

★ Closed string emission fails to converge because of non-pert. effects in $\alpha' = \ell_s^2$



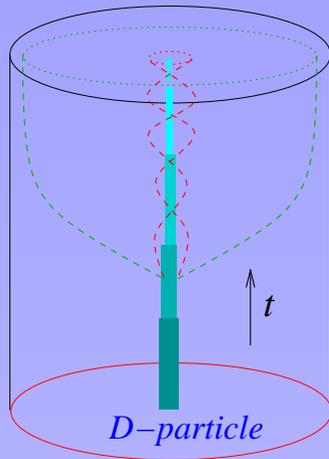
✓ **Conclusion:** most of the energy converted into highly excited long strings of winding $w \sim 1/g_s$, expanding at speed $d\rho/dt \sim 1/\ell_s\sqrt{k}$

- ★ Closed string emission fails to converge because of non-pert. effects in $\alpha' = \ell_s^2$
- ★ Production of short strings negligible in the perturbative regime $g_s \ll 1$ (since it does not depend on the coupling constant)



✓ **Conclusion:** most of the energy converted into highly excited long strings of winding $w \sim 1/g_s$, expanding at speed $d\rho/dt \sim 1/\ell_s\sqrt{k}$

- ★ Closed string emission fails to converge because of non-pert. effects in $\alpha' = \ell_s^2$
- ★ Production of short strings negligible in the perturbative regime $g_s \ll 1$ (since it does not depend on the coupling constant)
- ✓ **AdS₃/CFT₂ correspondence** string theory on AdS₃ dual to a symmetric product 2D CFT → dual description of tachyon decay?



✓ **Conclusion:** most of the energy converted into highly excited long strings of winding $w \sim 1/g_s$, expanding at speed $d\rho/dt \sim 1/\ell_s\sqrt{k}$

- ★ Closed string emission fails to converge because of non-pert. effects in $\alpha' = \ell_s^2$
- ★ Production of short strings negligible in the perturbative regime $g_s \ll 1$ (since it does not depend on the coupling constant)
- ✓ **AdS₃/CFT₂ correspondence** string theory on AdS₃ dual to a symmetric product 2D CFT → dual description of tachyon decay?
- ★ For AdS₅/CFT₄, corresponds to **sphaleron** decay *[Drukker, Gross, Itzhaki]*
- ★ Difficult here since CFT is **singular** (*unstable to fragmentation ↔ long strings emission*)

Remarks on Open String Pair Production

✓ Open string point of view: time-dependent Hamiltonian \rightarrow pair production

Mini-superspace limit : $[\partial_t^2 + \lambda e^t + \mathbf{p}^2 + N - 1] \psi(t) = 0$ *[Gutperle, Strominger '03]*

Remarks on Open String Pair Production

✓ Open string point of view: time-dependent Hamiltonian → pair production

Mini-superspace limit : $[\partial_t^2 + \lambda e^t + \mathbf{p}^2 + N - 1] \psi(t) = 0$ *[Gutperle, Strominger '03]*

★ String theory naturally "chooses" (from Liouville theory) the $|out\rangle$ vacuum:

$$\psi \propto H_{-2iE}^{(2)}(2\sqrt{\lambda}e^{t/2}) \stackrel{t \rightarrow -\infty}{\sim} e^{-iEt} + R(E)e^{iEt} \quad (R(E): \text{reflection coefficient})$$

Remarks on Open String Pair Production

✓ Open string point of view: time-dependent Hamiltonian → pair production

Mini-superspace limit : $[\partial_t^2 + \lambda e^t + \mathbf{p}^2 + N - 1] \psi(t) = 0$ *[Gutperle, Strominger '03]*

★ String theory naturally "chooses" (from Liouville theory) the *|out* vacuum:

$$\psi \propto H_{-2iE}^{(2)}(2\sqrt{\lambda}e^{t/2}) \stackrel{t \rightarrow -\infty}{\sim} e^{-iEt} + R(E)e^{iEt} \quad (R(E): \text{reflection coefficient})$$

→ Bogolioubov coefficient $\gamma = \frac{\beta_E}{\alpha_E} \leftrightarrow$ open string two-point function $\langle e^{iEt(\tau)} e^{-iEt(\tau')} \rangle$

★ Tension with Sen's conjecture in flat space?

Remarks on Open String Pair Production

✓ Open string point of view: time-dependent Hamiltonian → pair production

Mini-superspace limit : $[\partial_t^2 + \lambda e^t + \mathbf{p}^2 + N - 1] \psi(t) = 0$ *[Gutperle, Strominger '03]*

★ String theory naturally "chooses" (from Liouville theory) the $|out\rangle$ vacuum:

$$\psi \propto H_{-2iE}^{(2)}(2\sqrt{\lambda}e^{t/2}) \stackrel{t \rightarrow -\infty}{\sim} e^{-iEt} + R(E)e^{iEt} \quad (R(E): \text{reflection coefficient})$$

→ Bogolioubov coefficient $\gamma = \frac{\beta_E}{\alpha_E} \leftrightarrow$ open string two-point function $\langle e^{iEt(\tau)} e^{-iEt(\tau')} \rangle$

★ Tension with Sen's conjecture in flat space?

Rate of pair production $W = -\text{Re} \ln \langle out|in \rangle \sim \int dE \rho(E) e^{-2\pi E}$

→ power-law convergent only (divergent for $D_{p>22}$ in bosonic strings)

✓ High energy behavior of open string pair production in AdS_3

★ For open strings with angular momentum r , one gets (orbifold construction)

$$|R(E)| = \left| \frac{\sinh \pi(E+r/\sqrt{k}) \sinh \pi(E-r/\sqrt{k})}{\sinh^2 2\pi E} \right|$$

→ same large E asymptotics as in flat space

✓ High energy behavior of open string pair production in AdS_3

★ For open strings with angular momentum r , one gets (orbifold construction)

$$|R(E)| = \left| \frac{\sinh \pi(E+r/\sqrt{k}) \sinh \pi(E-r/\sqrt{k})}{\sinh^2 2\pi E} \right| \rightarrow \text{same large } E \text{ asymptotics as in flat space}$$

★ Density of states smaller ($\Delta_{\min} > 0$): $\rho(E) \sim \exp\{2\pi \sqrt{1 - 1/2k} \ell_s E\}$

→ open string production rate exponentially convergent for very massive open strings on the D0-brane in AdS_3

★ One gets that open string perturbative string (field) theory remains a valid description (*despite the disappearance of the brane!*)

④ *Open string tachyons in non-Critical Strings*

- ✓ **Non-critical superstrings:** superstrings in spacetime dimension $d < 10$
 - ➔ extra ($\mathcal{N} = 2$) Liouville (super-)field ϕ

④ *Open string tachyons in non-Critical Strings*

- ✓ **Non-critical superstrings:** superstrings in spacetime dimension $d < 10$
 - ➔ extra ($\mathcal{N} = 2$) Liouville (super-)field ϕ
- ★ Corresponds to string theory near genuine CY singularities

④ *Open string tachyons in non-Critical Strings*

- ✓ **Non-critical superstrings:** superstrings in spacetime dimension $d < 10$
 - ➔ extra ($\mathcal{N} = 2$) Liouville (super-)field ϕ
- ★ Corresponds to string theory near genuine CY singularities
- ★ **Example:** $\mathbb{C}^2/\mathbb{Z}_k$ orbifold singularity ➔ by T-duality, equivalent to k NS5-branes distributed on a transverse circle

④ Open string tachyons in non-Critical Strings

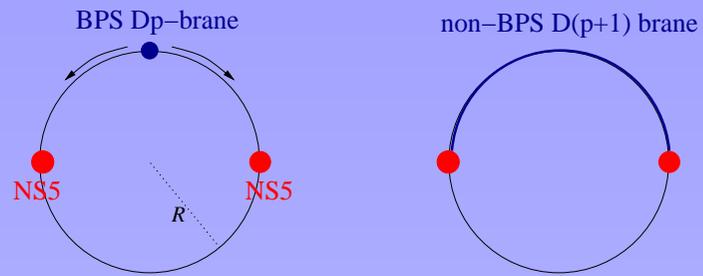
- ✓ **Non-critical superstrings:** superstrings in spacetime dimension $d < 10$
 - ➔ extra ($\mathcal{N} = 2$) Liouville (super-)field ϕ
- ★ Corresponds to string theory near genuine CY singularities
- ★ **Example:** $\mathbb{C}^2/\mathbb{Z}_k$ orbifold singularity ➔ by T-duality, equivalent to k NS5-branes distributed on a transverse circle

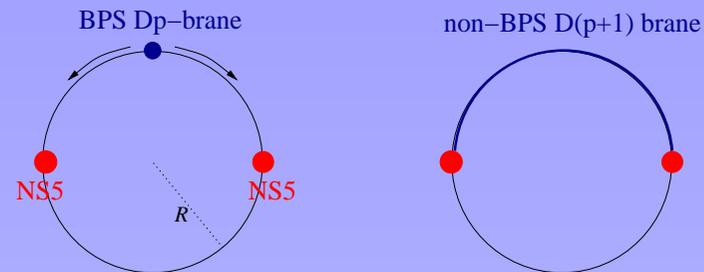
- ✓ **Universal and geometric tachyons**
- ★ BPS Dp -brane halfway between two NS5-branes at a maximum of its potential (attracted by both fivebranes) ➔ **geometric open string tachyon** on its worldvolume whose dynamics is very similar to the usual "universal" open string tachyon [Kutasov]

④ Open string tachyons in non-Critical Strings

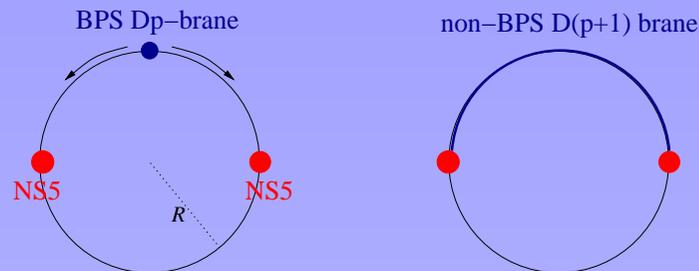
- ✓ **Non-critical superstrings:** superstrings in spacetime dimension $d < 10$
 - ➔ extra ($\mathcal{N} = 2$) Liouville (super-)field ϕ
- ★ Corresponds to string theory near genuine CY singularities
- ★ **Example:** $\mathbb{C}^2/\mathbb{Z}_k$ orbifold singularity ➔ by T-duality, equivalent to k NS5-branes distributed on a transverse circle

- ✓ **Universal and geometric tachyons**
- ★ BPS Dp -brane halfway between two NS5-branes at a maximum of its potential (attracted by both fivebranes) ➔ **geometric open string tachyon** on its worldvolume whose dynamics is very similar to the usual "universal" open string tachyon [Kutasov]
- ★ The universal tachyon appears on non-BPS $D(p + 1)$ -branes stretched between the NS5-branes

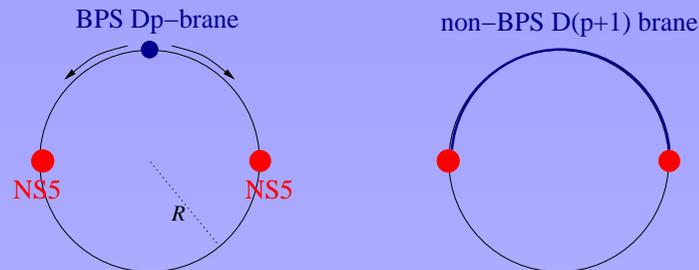




★ Both are solutions of the $D(p+1)$ -brane open string field theory (D_p is a kink)

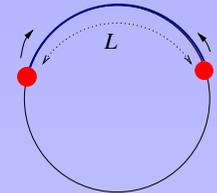


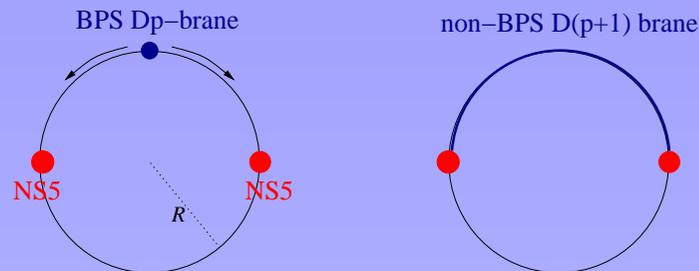
- ★ Both are solutions of the $D(p+1)$ -brane open string field theory (Dp is a kink)
- ★ Sen conjectured that these two solutions **merge** for $R < R_c$ *[Sen'07]*
- ➡ only true for one or two fivebranes, which leads to strong curvatures and require some non-renormalization of the tachyonic DBI effective action ➡ is there an exact string description?



- ★ Both are solutions of the $D(p+1)$ -brane open string field theory (Dp is a kink)
- ★ Sen conjectured that these two solutions **merge** for $R < R_c$ *[Sen'07]*
- ➡ only true for one or two fivebranes, which leads to strong curvatures and require some non-renormalization of the tachyonic DBI effective action ➡ is there an exact string description?

✓ We prove the conjecture in the **double scaling limit**: $g_s \rightarrow 0$, $\frac{L}{g_s \ell_s}$ fixed

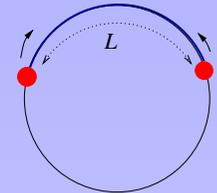




- ★ Both are solutions of the $D(p+1)$ -brane open string field theory (Dp is a kink)
- ★ Sen conjectured that these two solutions **merge** for $R < R_c$ *[Sen'07]*
- ➡ only true for one or two fivebranes, which leads to strong curvatures and require some non-renormalization of the tachyonic DBI effective action ➡ is there an exact string description?

✓ We prove the conjecture in the **double scaling limit**: $g_s \rightarrow 0$, $\frac{L}{g_s \ell_s}$ fixed

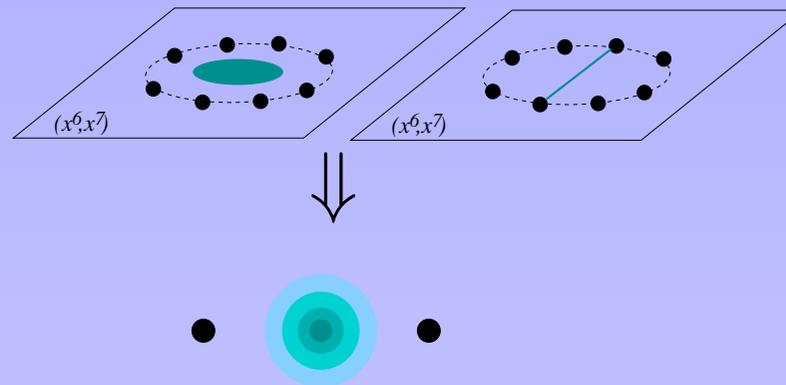
- ★ Gives near-horizon geometry of $k=2$ close parallel fivebranes in non-compact transverse space ➡ six-dimensional non-critical string



- ★ Start from the large k weakly curved regime (near-horizon geometry of k NS5-branes in non-compact transverse space), where one can trust the geometrical interpretation
- ★ For any $k \geq 2$ an exact worldsheet CFT description is available

- ★ Start from the large k weakly curved regime (near-horizon geometry of k NS5-branes in non-compact transverse space), where one can trust the geometrical interpretation
- ★ For any $k \geq 2$ an exact worldsheet CFT description is available
 - ➔ take k as the (discrete) control parameter

- ★ Start from the large k weakly curved regime (near-horizon geometry of k NS5-branes in non-compact transverse space), where one can trust the geometrical interpretation
- ★ For any $k \geq 2$ an exact worldsheet CFT description is available
 - ➔ take k as the (discrete) control parameter
- ★ One finds a **transition** at $k = 2$ where the Dp -brane with a geometric tachyon and the non-BPS suspended $D(p + 1)$ -brane boundary states become identical



- ★ Proof of geometric/universal tachyon equivalence including all α' corrections.

Rolling tachyon in non-critical strings

✓ **Mass gap** in the closed string spectrum $\ell_s m > \sqrt{8-d}/2$ in the closed string sector (δ -normalizable states)

➔ lower density of states $\rho(E) \sim \exp\{2\pi\sqrt{1 - \frac{8-d}{16}E}\}$ (higher Hagedorn temp.)

Rolling tachyon in non-critical strings

✓ **Mass gap** in the closed string spectrum $\ell_s m > \sqrt{8-d}/2$ in the closed string sector (δ -normalizable states)

➔ lower density of states $\rho(E) \sim \exp\{2\pi\sqrt{1 - \frac{8-d}{16}E}\}$ (higher Hagedorn temp.)

✓ From these considerations, it has been suggested that **closed string emission in non-critical string is UV-FINITE**

[Karczmarek, Liu, Hong, Maldacena, Strominger]

Rolling tachyon in non-critical strings

✓ **Mass gap** in the closed string spectrum $\ell_s m > \sqrt{8-d}/2$ in the closed string sector (δ -normalizable states)

➔ lower density of states $\rho(E) \sim \exp\{2\pi\sqrt{1 - \frac{8-d}{16}E}\}$ (higher Hagedorn temp.)

✓ From these considerations, it has been suggested that **closed string emission in non-critical string is UV-FINITE** [Karczmarek, Liu, Hong, Maldacena, Strominger]

- would raise a puzzle: what is the leftover of the brane mass? ($\ell_s m_D \sim 1/g_s^{\text{local}}$)
- would challenge Sen's conjecture ("universality" of DBI tachyonic action)

Decay of localized branes

- ✓ **Brane localized** in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
- ★ D_p with a geometric tachyon/suspended non-BPS $D(p + 1)$ seen before

Decay of localized branes

- ✓ Brane localized in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
- ★ Dp with a geometric tachyon/suspended non-BPS $D(p + 1)$ seen before
- ★ Discrete spectrum built on the *identity representation* of $\mathcal{N} = 2$ Liouville
 - ➔ identity \mathbb{I} is a normalizable state

Decay of localized branes

- ✓ Brane localized in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
- ★ D p with a geometric tachyon/suspended non-BPS D($p + 1$) seen before
- ★ Discrete spectrum built on the *identity representation* of $\mathcal{N} = 2$ Liouville
 - ➔ identity \mathbb{I} is a normalizable state
- ★ D-brane has an open string tachyon built on the identity
 - ➔ decay corresponds to $\delta S = \lambda \sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{X_0(x)/\sqrt{2}\ell_s}$

Decay of localized branes

- ✓ Brane localized in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
- ★ Dp with a geometric tachyon/suspended non-BPS D(p + 1) seen before
- ★ Discrete spectrum built on the *identity representation* of $\mathcal{N} = 2$ Liouville
 - ➔ identity \mathbb{I} is a normalizable state
- ★ D-brane has an **open string tachyon built on the identity**
 - ➔ decay corresponds to $\delta S = \lambda \sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{X_0(x)/\sqrt{2}\ell_s}$
- ✓ One-point function for the rolling tachyon: $\langle V \rangle_\lambda \propto \frac{\sinh 2\pi p_\phi \sinh \pi p_\phi}{\cosh \pi p_\phi + \cos \pi s} \frac{(\pi \lambda)^{2iE}}{\sinh \pi E}$

Decay of localized branes

- ✓ Brane localized in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
 - ★ Dp with a geometric tachyon/suspended non-BPS D(p + 1) seen before
 - ★ Discrete spectrum built on the *identity representation* of $\mathcal{N} = 2$ Liouville
 - ➔ identity \mathbb{I} is a normalizable state
 - ★ D-brane has an **open string tachyon built on the identity**
 - ➔ decay corresponds to $\delta S = \lambda \sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{X_0(x)/\sqrt{2}\ell_s}$
 - ✓ One-point function for the rolling tachyon: $\langle V \rangle_\lambda \propto \frac{\sinh 2\pi p_\phi \sinh \pi p_\phi}{\cosh \pi p_\phi + \cos \pi s} \frac{(\pi\lambda)^{2iE}}{\sinh \pi E}$
 - ★ Gives closed strings production
- $$\mathcal{N} \sim \int dE dp_\phi d\mathbf{p} \sum_N \rho(N) \left| \langle V_{p_\phi E \mathbf{p} s} \rangle_\lambda \right|^2 \delta(E^2 - p_\phi^2 - 2N - \mathbf{p}^2 + d/8)$$

Decay of localized branes

- ✓ Brane localized in the strong coupling end in $\mathcal{N} = 2$ Liouville (*cf.* ZZ brane)
- ★ Dp with a geometric tachyon/suspended non-BPS D(p + 1) seen before
- ★ Discrete spectrum built on the *identity representation* of $\mathcal{N} = 2$ Liouville
 - ➔ identity \mathbb{I} is a normalizable state
- ★ D-brane has an **open string tachyon built on the identity**
 - ➔ decay corresponds to $\delta S = \lambda \sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{X_0(x)/\sqrt{2}\ell_s}$
- ✓ One-point function for the rolling tachyon: $\langle V \rangle_\lambda \propto \frac{\sinh 2\pi p_\phi \sinh \pi p_\phi}{\cosh \pi p_\phi + \cos \pi s} \frac{(\pi\lambda)^{2iE}}{\sinh \pi E}$
- ★ Gives closed strings production
 - $\mathcal{N} \sim \int dE dp_\phi d\mathbf{p} \sum_N \rho(N) \left| \langle V_{p_\phi E \mathbf{p} s} \rangle_\lambda \right|^2 \delta(E^2 - p_\phi^2 - 2N - \mathbf{p}^2 + d/8)$
 - ➔ $\rho(N)$ smaller than in flat space, but $\int dp_\phi$ gives UV divergent production of closed strings as in flat space (due to non-pert. α' corrections)

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1 - d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1-d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$
- ✓ Non-BPS D-brane (*or* D/\bar{D} pair): open string tachyon of mass $\ell_s m = i\sqrt{d}/4$

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1-d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$
- ✓ Non-BPS D-brane (*or* D/\bar{D} pair): open string tachyon of mass $\ell_s m = i\sqrt{d}/4$
- ★ Homogeneous decay: $\delta S = \lambda \sigma^1 \oint dx G_{-1/2} e^{-\sqrt{1-d/8}\phi(x) + \frac{\sqrt{d}}{4\ell_s} X_0(x)}$
 - ➔ not a known conformal field theory (non-trivial along the dilaton gradient)

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1-d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$
- ✓ Non-BPS D-brane (*or* D/\bar{D} pair): open string tachyon of mass $\ell_s m = i\sqrt{d}/4$
- ★ Homogeneous decay: $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} e^{-\sqrt{1-d/8}\phi(x) + \frac{\sqrt{d}}{4\ell_s} X_0(x)}$
 - ➔ not a known conformal field theory (non-trivial along the dilaton gradient)
- ✓ One could instead deform the worldsheet with $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{\frac{X_0(x)}{\sqrt{2}\ell_s}}$

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1-d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$
- ✓ Non-BPS D-brane (*or* D/\bar{D} pair): open string tachyon of mass $\ell_s m = i\sqrt{d}/4$
- ★ Homogeneous decay: $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} e^{-\sqrt{1-d/8}\phi(x) + \frac{\sqrt{d}}{4\ell_s} X_0(x)}$
 - ➔ not a known conformal field theory (non-trivial along the dilaton gradient)
- ✓ One could instead deform the worldsheet with $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{\frac{X_0(x)}{\sqrt{2}\ell_s}}$
- ★ However the identity \mathbb{I} is **not normalizable** on the extended brane in Liouville theory (measure $\propto d\phi e^{\sqrt{4-d/2}\phi}$)

Decay of extended branes

- ✓ Brane extended along the dilaton gradient in $\mathcal{N} = 2$ Liouville (*cf.* FZZT brane)
- ★ Continuous spectrum (δ -norm) above a gap
 - ➔ vertex operators: $V_p(x) = \exp\{-\sqrt{1-d/8} + ip_\phi\}\phi(x) + p_\mu X^\mu(x) + \dots\}$
- ✓ Non-BPS D-brane (*or* D/\bar{D} pair): open string tachyon of mass $\ell_s m = i\sqrt{d}/4$
- ★ Homogeneous decay: $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} e^{-\sqrt{1-d/8}\phi(x) + \frac{\sqrt{d}}{4\ell_s} X_0(x)}$
 - ➔ not a known conformal field theory (non-trivial along the dilaton gradient)
- ✓ One could instead deform the worldsheet with $\delta S = \lambda\sigma^1 \oint dx G_{-1/2} \mathbb{I} \times e^{\frac{X_0(x)}{\sqrt{2}\ell_s}}$
- ★ However the identity \mathbb{I} is **not normalizable** on the extended brane in Liouville theory (measure $\propto d\phi e^{\sqrt{4-d/2}\phi}$)
 - ➔ does not represent the decay of the open string tachyon but changes the boundary conditions at $\phi \rightarrow -\infty$ (*however leads to a UV-finite result*)

Conclusions

- D-brane decay \rightarrow involves all the tower of string modes
- Non-perturbative α' effects & asympt. density of states are crucial ingredients
- String theory clever enough to convert all brane mass into closed strings although the detailed outcome may be different
 \rightarrow Sen's conjecture seems universal !
- However, perturbative string theory leaves many issues open (backreaction)
- Future work: tachyon condensation and D/\bar{D} scattering, relevant for cosmology