

Complex eigenvalues of quadratised matrices

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based on joint work (J Phys A, 2012) with:

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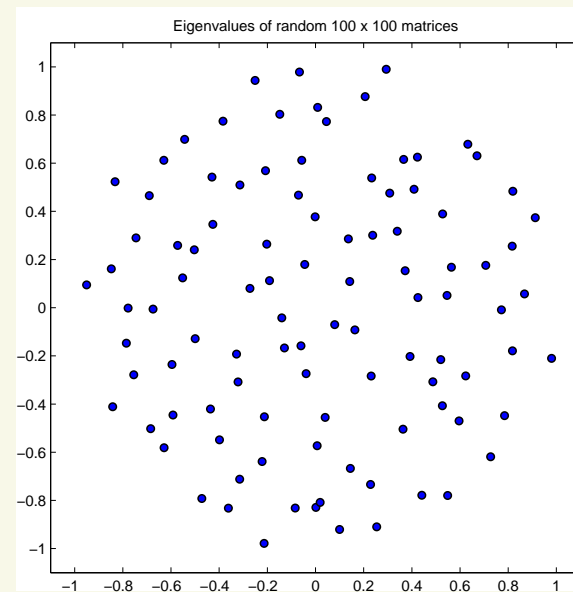
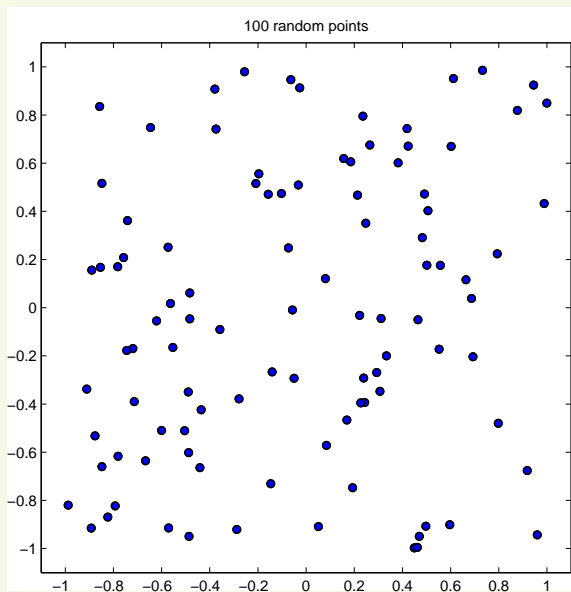
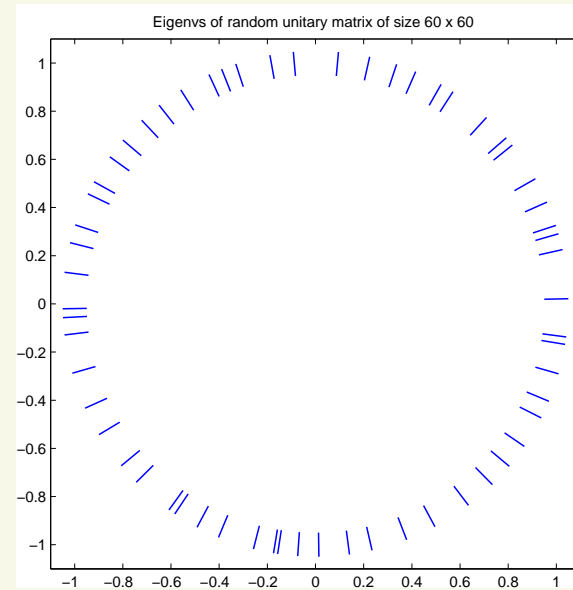
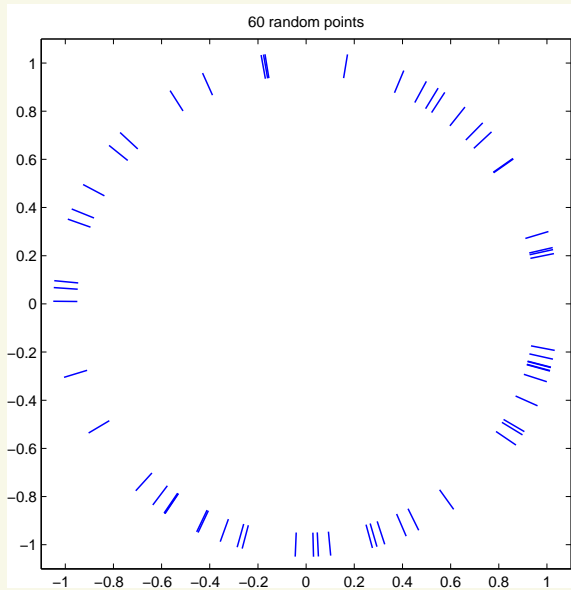
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Outline of talk

- **EV distributions in the complex plane**
- **Quadratisation of rectangular matrices**
- **Induced complex Ginibre ensemble**
- **Induced real Ginibre ensemble**
- **Conclusions**

Random matrices as a probability machine



Two matrix decompositions

Quite a lot of interest to EV stats in the complex plane in the last 20 yrs.

Two matrix decompositions ('coordinate systems'):

- $$G = \frac{G + G^\dagger}{2} + i \frac{G - G^\dagger}{2i}.$$

If G is complex then $G = H_1 + iH_2$ with $H_{1,2}$ Hermitian; and if G is real then $G = S + A$ with S symmetric and A asymmetric;

- $G = RU$ where $R = \sqrt{G^\dagger G}$ and U is unitary (orthogonal).

Correspondingly, one can describe matrix distributions in terms of the 'real' and 'imaginary' part of G , or 'radial' and 'angular' part. Polar decomposition is natural for rotationally invariant matrix distributions like

$$p(G) \propto \exp[-\text{Tr} \phi(GG^\dagger)]$$

Of course, there are many other matrix decompositions.

Cartesian picture: Complex Matrices

$G = H_1 + i\alpha H_2$, where $H_{1,2}$ are Hermitian, independent Gaussian (GUE)

If $\alpha = 1$ then have Ginibre's ensemble $p(G) \propto \exp[-\text{Tr} GG^\dagger]$. For large N have uniform distr of EV in a disk, and EV corr fncs are known in closed form (Ginibre 1965).

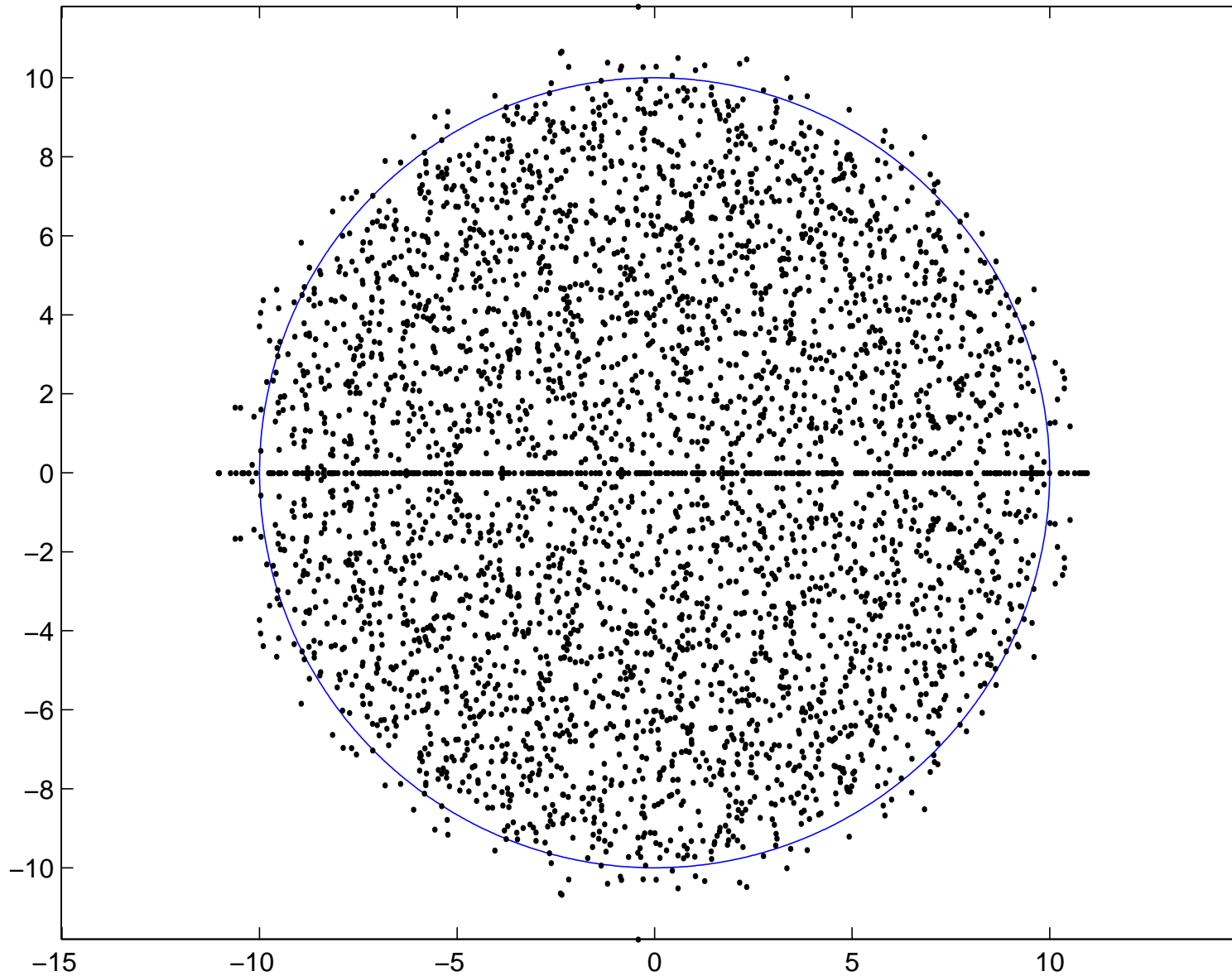
If $|\alpha| < 1$ then for large N have uniform distr of EV in a ellipse (Girko 1985) and the Ginibre EV correlations. If $\alpha \rightarrow 0$ the have weakly non-Hermitian matrices, crossover from Wigner-Dyson to Ginibre EV correlations (Fyodorov, Kh, Sommers 1997).

Beyond Gaussian distribution:

- G with i.i.d. entries: Circular Law (Girko 1984, Bai 1997, also Götze & Tikhomirov 2010, Tao & Vu 2008, 2010), if allow for correlated pairs symmetric about the main diagonal, $\langle G_{jk} G_{kj} \rangle = 1 - \alpha^2, |\alpha| < 1$, then have Elliptic Law (Girko 1985, Naumov 2012)). EV corr fncs are not known in this case.
- weak non-Hermiticity $H_1 \in \text{GUE}$, H_2 is finite rank, fixed, EV corr fncs known in closed form (Fyodorov & Kh 1999).

Real Matrices

Eigvs of a real Ginibre matrix
N = 100, no. of samples = 40.



Real Matrices

Have the Elliptic Law of distribution of EVs for matrices with pairwise (G_{jk}, G_{kj}) correlations, and the Circular Law for i.i.d.

For Gaussian matrices finer details of EV distribution are available via EV jpdf (Lehmann & Sommers 1991, and Edelman 1993).

The expected no of real EV is propto \sqrt{N} in the limit of large matrix dim N and real EV have uniform distribution (Edelman, Kostlan, & Shub 1994, Forrester & Nagao 2007)

Away from the real line have Ginibre correlations (Akemann & Kanzieper 2007, Forrester & Nagao 2007). New EV correlations on the real line (Forrester & Nagao 2007) and near the real line (Borodin & Sinclair 2009). Alternative derivation by Sommers 2007 and Sommers & Weiczorek 2008.

Weakly non-Hermitian limit $G = S + \alpha A$, with $\alpha \rightarrow 0$ by Efetov 1997 and Forrester & Nagao 2009.

Polar decomposition picture: EV density

Consider random matrices in the form $G = RU$, with R, U independent, $R \geq 0$ and U Haar unitary (orthogonal), equivalently $V\Lambda U$ (SVD)

Introduced by Fyodorov & Sommers 2003, with finite rank R , in the context of resonances in open chaotic sys.

Another example, the Feinberg-Zee ensemble $p_{FZ} \propto \exp[-N \text{Tr} \phi(G^\dagger G)]$, with ϕ polynomial. Can be recast into RU by SVD, $G = V\Lambda U$.

For finite- N the mean EV density of RU is known in terms EV of R (Wei & Fyodorov 2008), the large- N limit performed by Bogomolny 2010.

Single Ring Theorem for $G = V\Lambda U$ by Guionnet, Krishnapur & Zeitouni 2011, proving Feinberg-Zee 1997 – an analogue of the Circular Law.

Finer details of EV distribution, e.g. corr fnc or distribution of real EVs for real matrices are difficult (if possible) for general R . Known in several exactly solvable cases beyond Gaussian G : spherical ensemble (Forrester & Krishnapur 2008, Forrester & Mays 2010) truncations of Haar unitaries/orthogonals (Życzkowski & Sommers 2000, Kh, Sommers & Życzkowski 2010)

Quadratisation of rectangular matrices

Consider X with M rows and N columns, $M > N$, matrices are 'standing'.

Y and Z are the upper $N \times N$ block and lower $(M - N) \times N$ of X . Want a unitary transformation $W \in U(M)$ s.t.

$$W^\dagger X = W^\dagger \begin{bmatrix} Y \\ Z \end{bmatrix} = \begin{bmatrix} G \\ 0 \end{bmatrix}.$$

W exists, def. up to right multiplication by $\text{diag}[U, V]$, $U \in U(N), U(M - N)$

Thus $W \in U(M)/U(N) \times U(M - N) \dots$

Parameter count:

- X has $2MN$ real parameters.
- W has $M^2 - N^2 - (M - N)^2 = 2MN - N^2$ real parameters
- G has N^2 real parameters.

Call G quadratization of X

Quadratisation of rectangular matrices

Can parametrise cosets by $N \times (M - N)$ matrices C :

$$\tilde{W} = \begin{bmatrix} (1_N - CC^\dagger)^{1/2} & C \\ -C^\dagger & (1_{M-N} - C^\dagger C)^{1/2} \end{bmatrix},$$

Prop. 1 Let $M > N$. For any $X \in \text{Mat}(M, N)$ of full rank there exist unique \tilde{W} as above and $G \in \text{Mat}(N, N)$ such that

$$\tilde{W}^\dagger X = \tilde{W}^\dagger \begin{bmatrix} Y \\ Z \end{bmatrix} = \begin{bmatrix} G \\ 0 \end{bmatrix}.$$

The square matrix G is given by $G = (1_N + \frac{1}{Y^\dagger} Z^\dagger Z \frac{1}{Y})^{1/2} Y$.

Prop. 2 If $X \in \text{Mat}(M, N)$ is Gaussian with density $p(X) \propto e^{-\frac{\beta}{2} \text{Tr} X^\dagger X}$ ($\beta = 1$ or $\beta = 2$ depending on whether X is real or complex) then its quadratisation $G \in \text{Mat}(N, N)$ has density

$$p_{\text{IndG}}(G) \propto (\det G^\dagger G)^{\frac{\beta}{2} (M-N)} \exp\left(-\frac{\beta}{2} \text{Tr} G^\dagger G\right)$$

Proof of Prop 2

Here is one based on SVD (works for non-Gaussian weights as well):

Ignoring a set of zero prob., $X = Q \Sigma^{1/2} P^\dagger$, where $\Sigma^{1/2}$ is diag mat of ordered SVs of X , and $Q \in U(M)/U(M-N)$ and $P \in U(N)/U(1)^N$. By computation, $d\nu(X) = d\mu(Q)d\mu(P)d\sigma(\Sigma)$, where $d\mu$ is Haar and

$$d\sigma(\Sigma) \propto (\det \Sigma)^{\frac{\beta}{2}(M-N+1-\frac{2}{\beta})} e^{-\frac{\beta}{2} \text{Tr} \Sigma} \prod_{j < k} |s_k - s_j|^\beta \prod_{j=1}^N ds_j .$$

Introduce independent Haar unitary $U \in U(N)$ and rewrite SVD as

$$X = QUU^\dagger \Sigma^{1/2} P^\dagger = QUG; \quad G = U^\dagger \Sigma^{1/2} P^\dagger .$$

The pdf of G follows by rolling the argument back from $U^\dagger \Sigma^{1/2} P^\dagger$ to G .

Thus, with $\tilde{Q} := QU \in U(M)/U(M-N)$ and $W \in U(M)/U(M-N) \times U(N)$,

$$X = \tilde{Q}G = W \begin{bmatrix} G \\ 0 \end{bmatrix} ,$$

as required. □

Polar decomposition revisited

Relation $G = \left(1_N + \frac{1}{Y^\dagger} Z^\dagger Z \frac{1}{Y}\right)^{1/2} Y$ provides a recipe for sampling the distribution

$$p_{\text{IndG}}(G) \propto (\det G^\dagger G)^{\frac{\beta}{2}(M-N)} \exp\left(-\frac{\beta}{2} \text{Tr} G^\dagger G\right)$$

Interestingly, by rearranging $X = QUU^\dagger \Sigma^{1/2} P^\dagger = QUG$ one obtains another recipe (G is just Haar unitary times sq. root of Wishart).

Prop 3 Suppose that U is $N \times N$ Haar unitary (real orthogonal for $\beta = 1$) and X is $M \times N$ Gaussian, independent of U . Then the $N \times N$ matrix

$$G = U(X^\dagger X)^{1/2}$$

has the distribution $p_{\text{IndG}}(G)$.

Proof. Since $(X^\dagger X)^{1/2} = P\Sigma^{1/2}P^\dagger$, then

$$G = U^\dagger Q^\dagger X = U^\dagger P^\dagger (X^\dagger X)^{1/2} = \tilde{U}^\dagger (X^\dagger X)^{1/2}.$$

\tilde{U}^\dagger is Haar unitary and independent of X . □

A few comments on quadratisation of random matrices

Quadratisation is well defined: as density of quadratised matrix doesn't depend on which rows W nullifies. Thus, can introduce *quasi spectrum* as the spectrum of the quadratised matrix.

Distributions other than Gaussian: Consider $X \in \text{Mat}(M, N)$ with density $p_{\text{FZ}}(X) \propto \exp[-\text{Tr} \phi(X^\dagger X)]$. On applying the procedure of quadratisation, one obtains the induced Feinberg-Zee ensemble

$$p_{\text{IndFZ}}(G) \propto (\det G^\dagger G)^{\frac{\beta}{2}(M-N)} \exp[-\text{Tr} \phi(G^\dagger G)] .$$

'Eigenvalue' map: Embed $\text{Mat}(M, N)$ into $\text{Mat}(M, M)$ by augmenting X with $M - N$ zero column-vectors and write the quadratisation rule in terms of square matrices albeit with zero blocks:

$$W^\dagger \begin{bmatrix} Y & 0 \\ Z & 0 \end{bmatrix} = \begin{bmatrix} G & 0 \\ 0 & 0 \end{bmatrix}; \quad \text{or} \quad W^\dagger \tilde{X} = \tilde{G}$$

EV map: the zero EV of \tilde{X} stays put, and its multiplicity is conserved, and, otherwise, the EV of Y are mapped onto those of G .

Quadratising complex Gaussian matrices: EV density

Complex matrices are straightforward, can apply the method of OP.

The EV-corr fnc are $R_n^{(N)}(\lambda_1, \dots, \lambda_n) = \det(K_N(\lambda_k, \lambda_l))$,

$$K_N(\lambda_k, \lambda_l) = \frac{1}{\pi} e^{-\frac{1}{2}|\lambda_k|^2 - \frac{1}{2}|\lambda_l|^2} \sum_{j=M-N}^{M-1} \frac{(\lambda_k \bar{\lambda}_l)^j}{j!}.$$

For M, N large and $M - N = \alpha N$ the EV density $\rho_N(\lambda) = K_N(\lambda, \lambda)$ is uniform in the ring $r_{in} < |\lambda| < r_{out}$, $r_{in} = \sqrt{M - N}$, $r_{out} = \sqrt{M}$, i.e.

$$\lim_{N \rightarrow \infty} \rho_N(\sqrt{N}z) = \frac{1}{\pi} \text{ for all } |z| \in (\sqrt{\alpha}, \sqrt{1 + \alpha}).$$

Close to the circular edges of EV support,

$$\lim_{N \rightarrow \infty} \rho_N((r_{out} + \xi)e^{i\phi}) = \lim_{N \rightarrow \infty} \rho_N((r_{in} - \xi)e^{i\phi}) = \frac{1}{2\pi} \operatorname{erfc}(\sqrt{2}\xi),$$

where $\operatorname{erfc}(x) := \frac{2}{\sqrt{\pi}} \int_x^\infty e^{-t^2} dt$, hence the EV density falls from $1/\pi$ to zero very fast, at a Gaussian rate.

Quadratising complex Gaussian matrices: EV correlations

Not surprisingly in the regime of strong non-rectangularity ($M, N \rightarrow \infty$ and $M - N = \alpha N$) one recovers the Ginibre correlations both in the bulk and at the circular edges of the eigenvalue distribution.

In the bulk: Set $\lambda_k = \sqrt{N}u + z_k$. Then for any $|u| \in (\sqrt{\alpha}, \sqrt{\alpha + 1})$:

$$\lim_{N \rightarrow \infty} R_n^{(N)}(\lambda_1, \dots, \lambda_n) = \det \left[\frac{1}{\pi} \exp \left(-\frac{|z_j|^2}{2} - \frac{|z_k|^2}{2} + z_j \bar{z}_k \right) \right]_{j,k=1}^n$$

In particular, $R_2(\lambda_1, \lambda_2) = 1 - \exp(-|z_1 - z_2|^2)$.

At the circular edges: Let $|u| = 1$ and set $\lambda_k = \sqrt{M}u + z_k$. Then:

$$\lim_{N \rightarrow \infty} R_n^{(N)}(\lambda_1, \dots, \lambda_n) = \det \left[\frac{1}{2\pi} e^{-\frac{|z_j|^2}{2} - \frac{|z_k|^2}{2} + z_j \bar{z}_k} \operatorname{erfc} \left(\frac{z_j \bar{u} + \bar{z}_k u}{\sqrt{2}} \right) \right].$$

The same limiting expression is found around the inner edge $\sqrt{M - N}$ of the eigenvalue density by setting $\lambda_k = \sqrt{M - N}u - z_k$.

Almost square matrices: emergence of the hole in the spectrum

Explore a different regime when the rectangularity index $L = M - N \ll N$. This corresponds to the quadratisation of almost square matrices.

At the origin the EV density vanishes algebraically, $\rho_N(\lambda) \sim \frac{1}{\pi} \frac{|\lambda|^{2L}}{L!}$ as $\lambda \rightarrow 0$, uniformly in N . Away from the origin, the density reaches its asymptotic value $1/\pi$ very quickly. This plateau extends to a full circle of radius \sqrt{N} .

Also, $\lim_{N \rightarrow \infty} R_n^{(N)}(\lambda_1, \dots, \lambda_n) = \det (K_{origin}(\lambda_j, \lambda_k))_{j,k=1}^n$, with

$$K_{origin}(\lambda_j, \lambda_k) = \frac{1}{\pi} \frac{1}{\Gamma(L)} \int_0^{\lambda_j \bar{\lambda}_k} t^{L-1} e^{-t} dt .$$

The hole prob $A(s)$ is the prob that no EV lies inside the disk $|\lambda| < s$:

$$A(s) = \prod_{j=1}^N \frac{\Gamma(j+L, s^2)}{\Gamma(j+L)} \text{ with } \Gamma(a, x) := \int_x^\infty e^{-t} t^{a-1} dt .$$

For almost square matrices $A(s) = 1 - \frac{s^{2(L+1)}}{(L+1)!} + O\left(\frac{s^{2(L+2)}}{(L+2)!}\right)$.

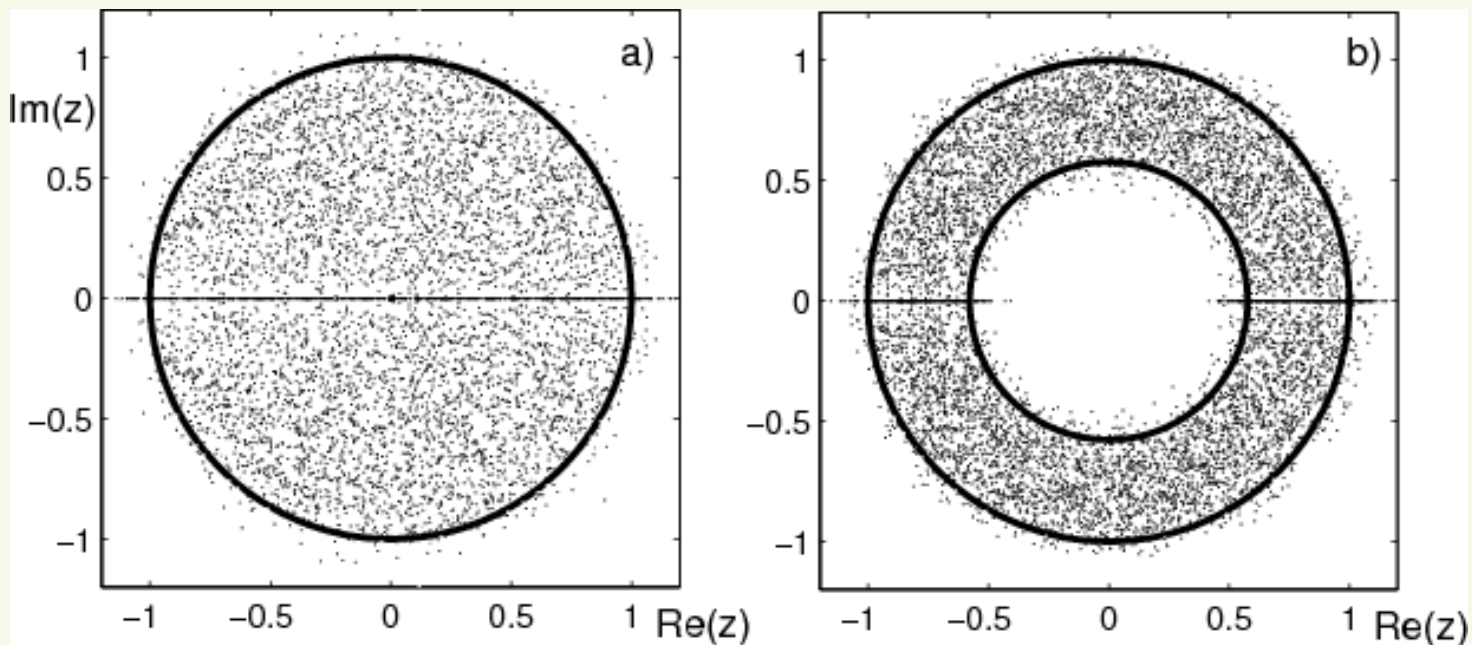
Quadratisation of real matrices

Consider real Gaussian matrices $X \in \text{Mat}(M, N)$, $X \mapsto G \in \text{Mat}(N, N)$

Gaussian measure on X induces a measure on G with density

$$p_{\text{IndG}}(G) \propto (\det G^T G)^{\frac{1}{2}(M-N)} \exp\left(-\frac{1}{2} \text{Tr} G^T G\right)$$

Spectra of G from the induced distribution for $N = 128$ and a) $M = N = 128$ (no hole) and b) $M = N + 32$. Each picture consists of 128 independent realisations and the spectra are rescaled by a factor $1/\sqrt{M}$



Real matrices

Building on the recent advances for the real Ginibre ensemble, the EV corr fnc are given in Pfaffian form, with the kernel expressed in terms of skew-OPs $\{q_j\}_{j=1,\dots}$, monic polynomials defined by

$$(q_{2j}, q_{2k}) = (q_{2j+1}, q_{2k+1}) = 0$$

$$(q_{2j}, q_{2k+1}) = -(q_{2j+1}, q_{2k}) = r_j \delta_{jk}$$

with the skew-symmetric inner product $(-, -)$:

$$(f, g) = (f, g)_{\mathbb{R}} + (f, g)_{\mathbb{C}}$$

$$(f, g)_{\mathbb{R}} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{1}{2}(x^2+y^2)} |xy|^L \operatorname{sgn}(y-x) f(x)g(y) dx dy$$

$$(f, g)_{\mathbb{C}} = 2i \int_{\mathbb{R}_+^2} e^{y^2-x^2} \operatorname{erfc}(\sqrt{2}y) (x^2 + y^2)^L \times \\ [f(x+iy)g(x-iy) - g(x+iy)f(x-iy)] dx dy.$$

Finding skew-OPs is a formidable task. However in the RM set up:

$$2 \sum_{j=0}^{\frac{N}{2}-1} \frac{1}{r_j} [q_{2j}(w)q_{2j+1}(w') - q_{2j+1}(w)q_{2j}(w')] = \\ \frac{1}{r_N} (w - w') \langle \det(G - w1_{N-2}) \det(G - w'1_{N-2}) \rangle_{N-2}$$

2nd moment of the characteristic polynomial

$\langle \det(G - w1_{N-2}) \det(G - w'1_{N-2}) \rangle_{N-2}$ can be evaluated in closed form.

Expand the product of determinants in Schur pols, decompose $G = O\sqrt{G^T G}$ and then integrate over the orthogonal group. The remaining integral is over the radial part of G :

$$\langle \det(G - wI) \det(G - w'I) \rangle_{N-2} = \sum_{j=0}^{N-2} \frac{\langle \epsilon_j(GG^T) \rangle_{N-2}}{\binom{n}{j}} (ww')^{N-2-j}$$

Here $\epsilon_j(GG^T)$ denotes the j -th elementary symm pol in EVs of GG^T and its average is given by the Selberg-Aomoto integral, leading to

$$2 \sum_{j=0}^{\frac{N}{2}-1} \frac{1}{r_j} [q_{2j}(w)q_{2j+1}(w') - q_{2j+1}(w)q_{2j}(w')] = \frac{1}{r_N} (w - w') \frac{(L+N-2)!}{\sqrt{2\pi}} \sum_{j=0}^{N-2} \frac{(ww')^j}{(L+j)!},$$

from which $q_0(w) = 1$, $q_1(w) = w$ and

$$q_{2j}(w) = w^{2j}, q_{2j+1}(w) = w^{2j+1} - (2j + L)w^{2j-1},$$

$$(q_{2j}, q_{2j+1}) = 2\sqrt{2\pi}\Gamma(L + 2j + 1).$$

Large- N limit for real matrices

Density of complex EVs:

$$\rho_N^C(x + iy) = \sqrt{\frac{2}{\pi}} y \operatorname{erfc}(\sqrt{2}y) e^{y^2 - x^2} \sum_{j=0}^{N-2} \frac{(x^2 + y^2)^{j+L}}{(j+L)!}$$

vanishes on the real line. Same distribution in the bulk and at the edges as for complex matrices, though convergence is not uniform near \mathbb{R} .

Expected number of real EVs:

In the leading order the average no of real EV is $\sqrt{\frac{2}{\pi}}(\sqrt{N+L} - \sqrt{L})$

(Mean) Density of real EVs:

$$\lim_{N \rightarrow \infty} \rho_N^{(R)}(\sqrt{N}x) = \begin{cases} \frac{1}{\sqrt{2\pi}} & \text{if } |x| \in (\sqrt{\alpha}, \sqrt{\alpha+1}) \\ 0 & \text{if } |x| \notin [\sqrt{\alpha}, \sqrt{\alpha+1}] \end{cases}$$

Tails of real density EVs:

$$\lim_{N \rightarrow \infty} \rho_N^R(\sqrt{L}-\xi) = \lim_{N \rightarrow \infty} \rho_N^R(\sqrt{L+N}+\xi) = \frac{1}{\sqrt{2\pi}} \left[\operatorname{erfc}(\sqrt{2}\xi) + \frac{1}{2\sqrt{2}} e^{-\xi^2} \operatorname{erfc}(-\xi) \right].$$

Same as in the real Ginibre, also the real Ginibre EV correlations in the bulk and at the edges!

Conclusions

A quadratisation procedure for rectangular matrices is introduced. When applied to Gaussian matrices, it leads to modified Ginibre ensemble with a pre-exponential term (the radial part squared is Wishart/Laguerre). This ensemble can be solved exactly for EV densities and corr fnc:

- for complex matrices the EV density is uniform in a ring, almost square matrices describe cross-over from disk to ring;
- for real matrices EV the density of complex EV is uniform in a ring, and the density of real EV is uniform on two intervals symmetric about the origin. The expected no of real EVs scales with \sqrt{N} ;
- the EV corr functions are same as in the pure Ginibre ensembles

Work in progress (Fischmann's PHD thesis): applying quadratisation procedure to non-Gaussian distribution, e.g. rectangular truncations of Haar unitaries or orthogonals, EV densities are no longer uniform however recover the Ginibre corr fnc after unfolding the spectrum.

THANK YOU!