

Out of equilibrium dynamics of the vortex glass

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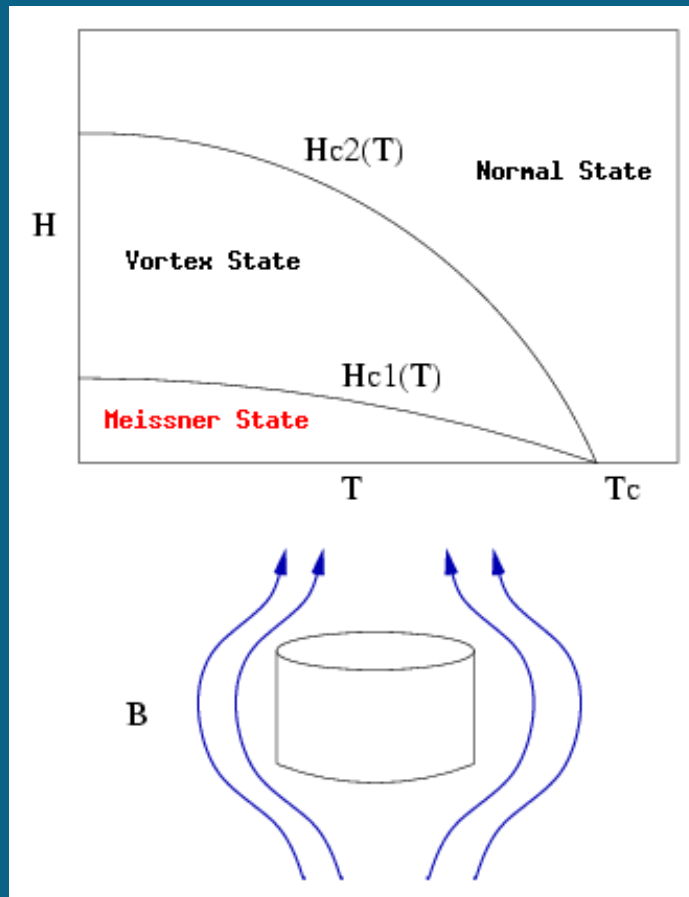


Outline

- Vortices in Superconductors
- Vortex dynamics model and vortex glass
- Single vortex line: Aging and FDT
- Vortex glass: Aging and FDT
- Conclusions

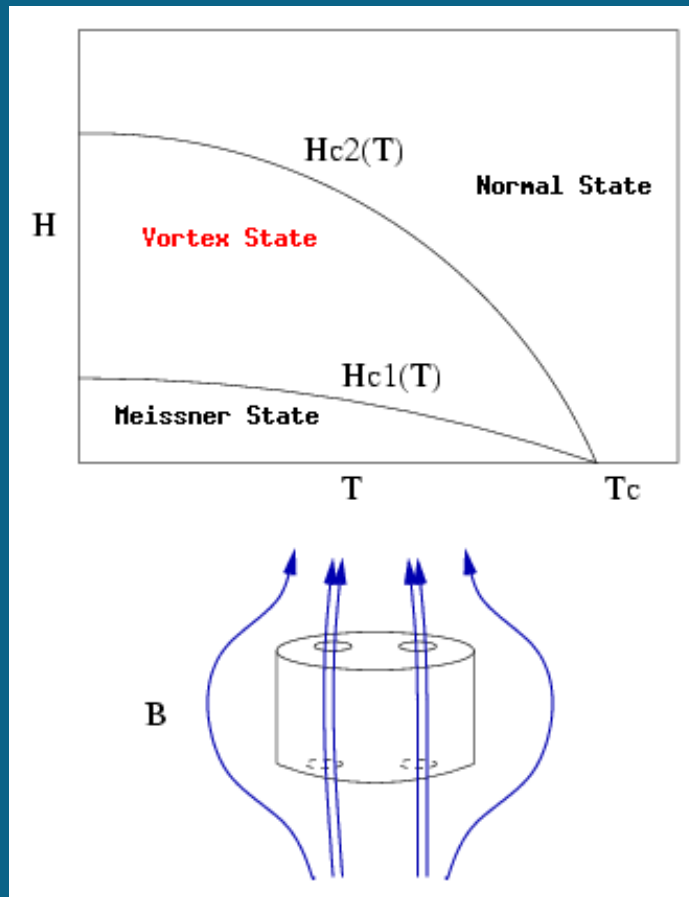


Type II Superconductors: Vortices



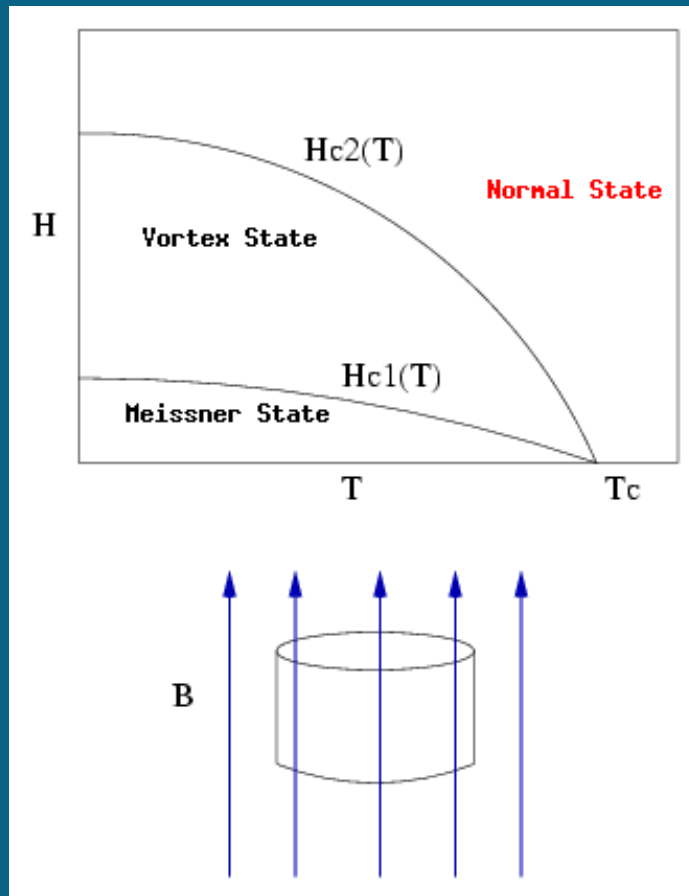


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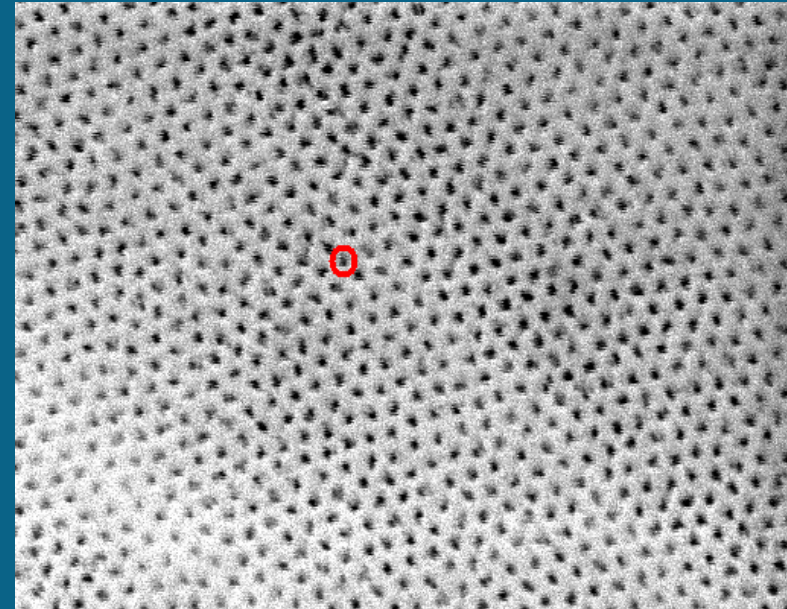
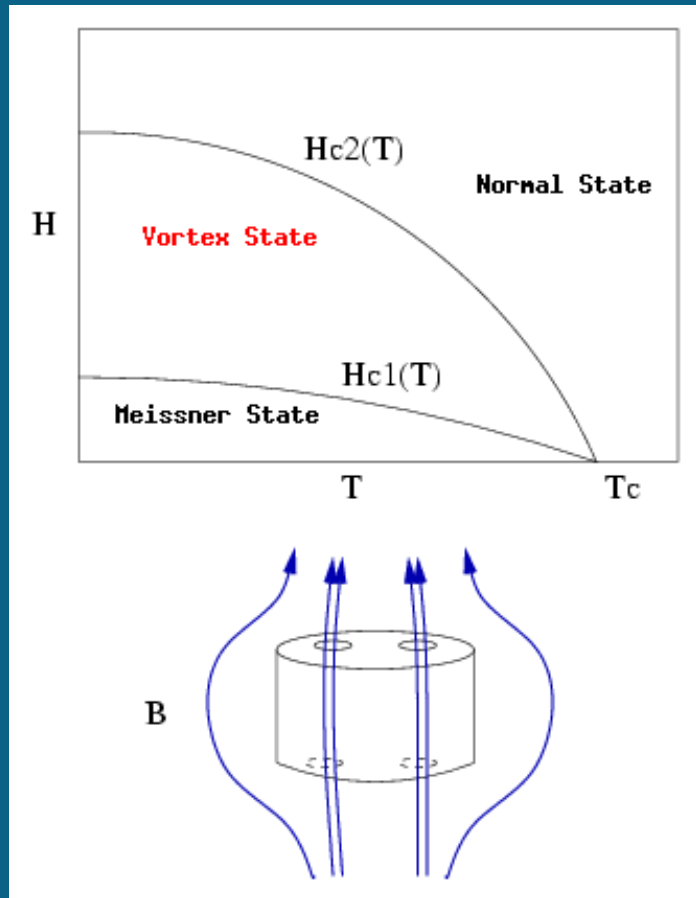


Type II Superconductors: Vortices



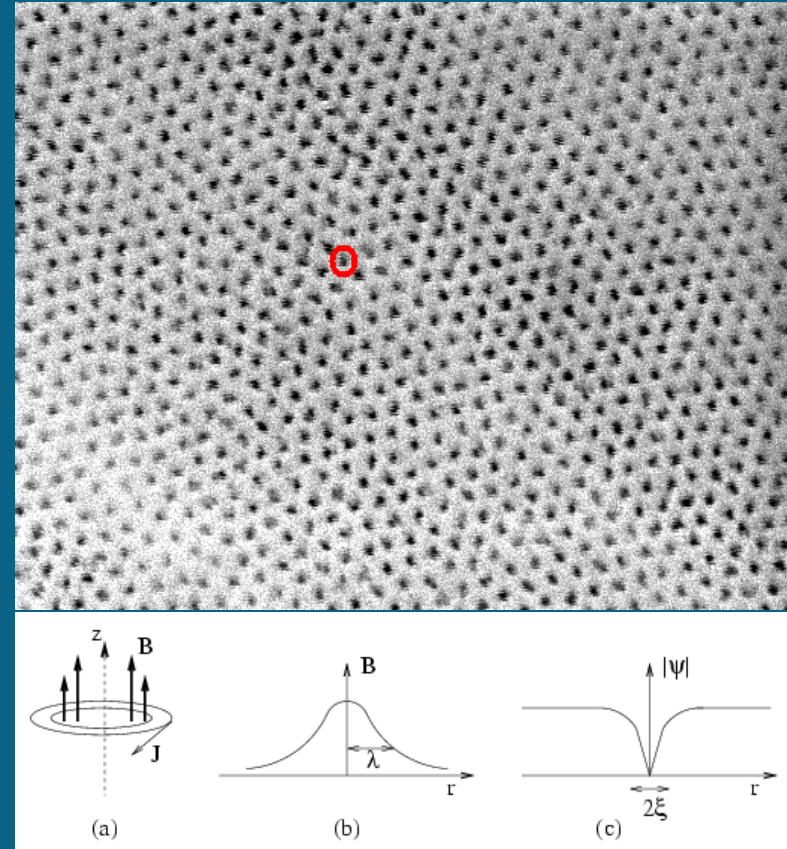
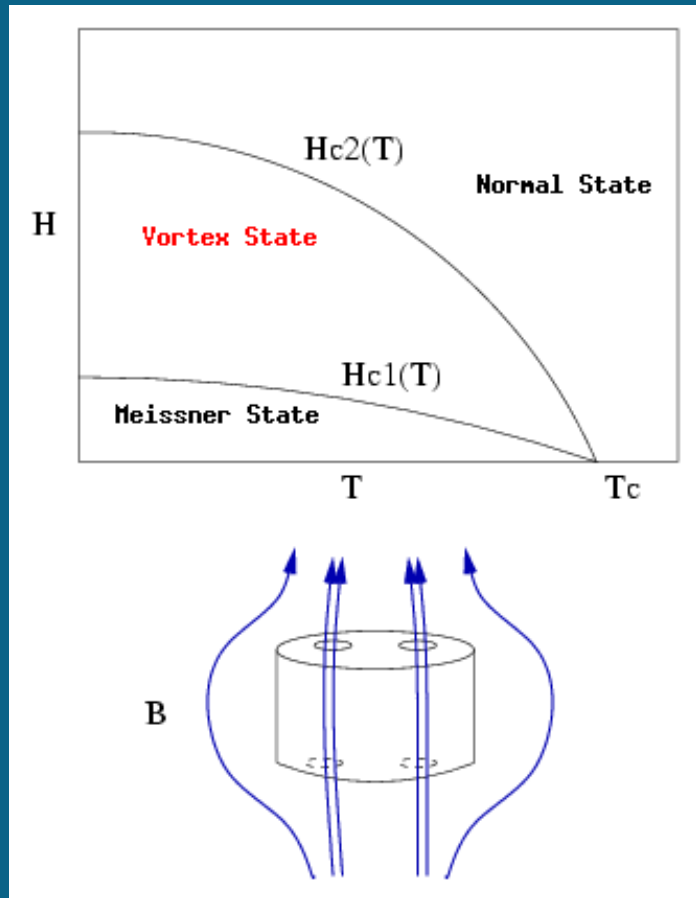


Type II Superconductors: Vortices





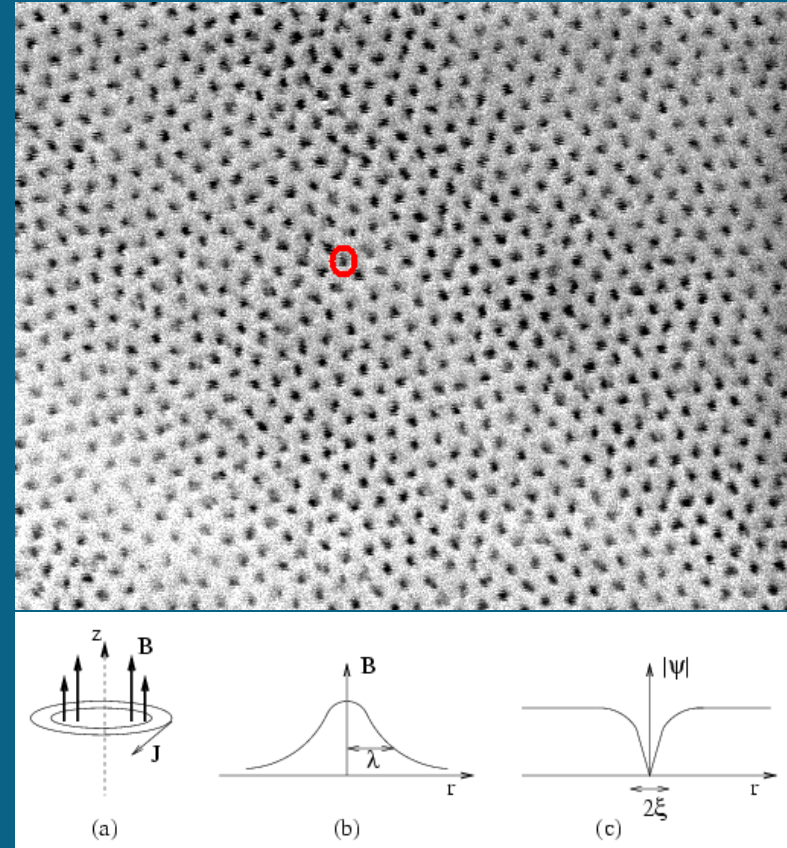
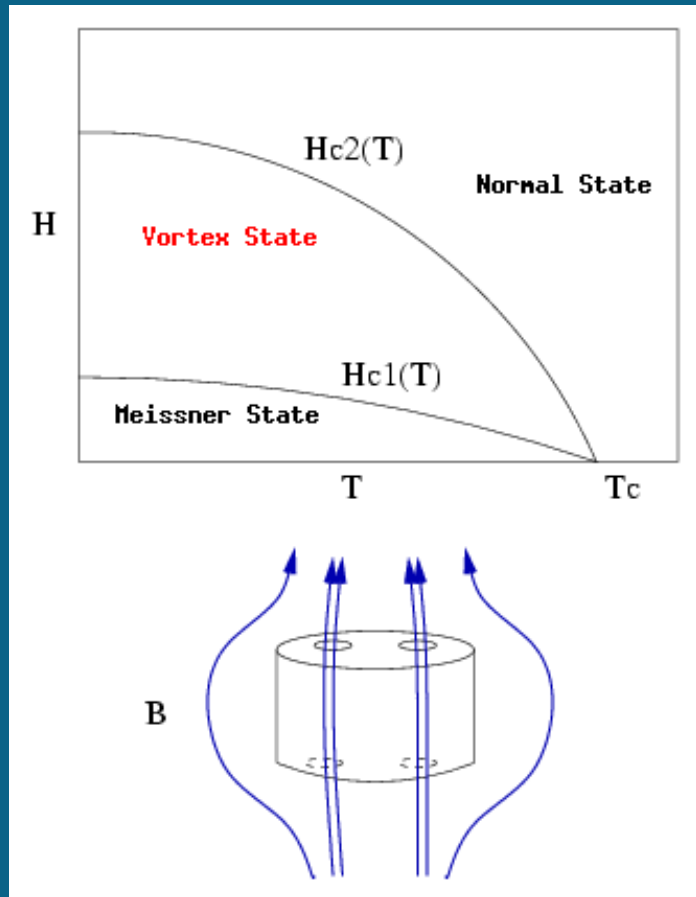
Type II Superconductors: Vortices



- $B = n_v \Phi_0$



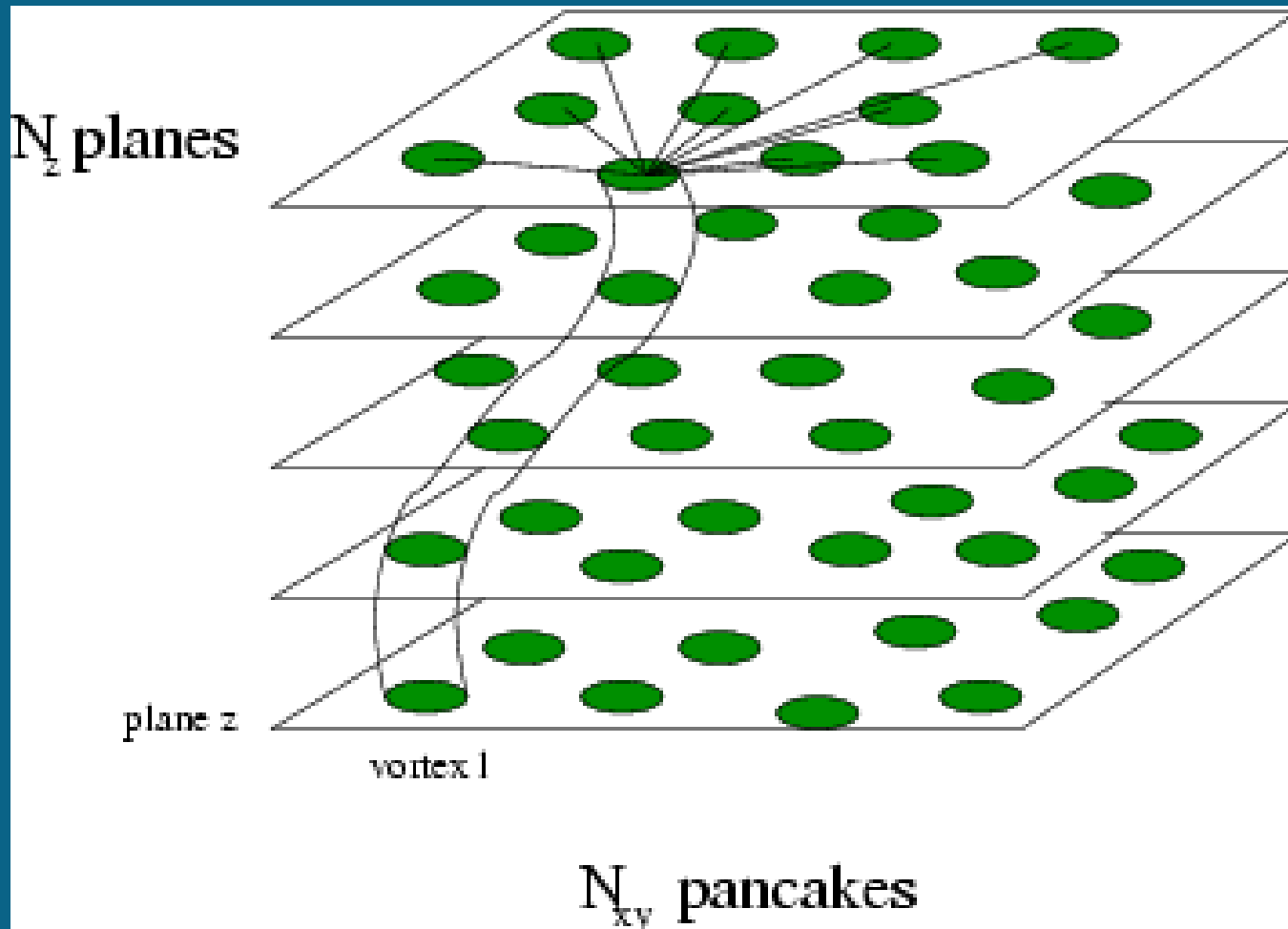
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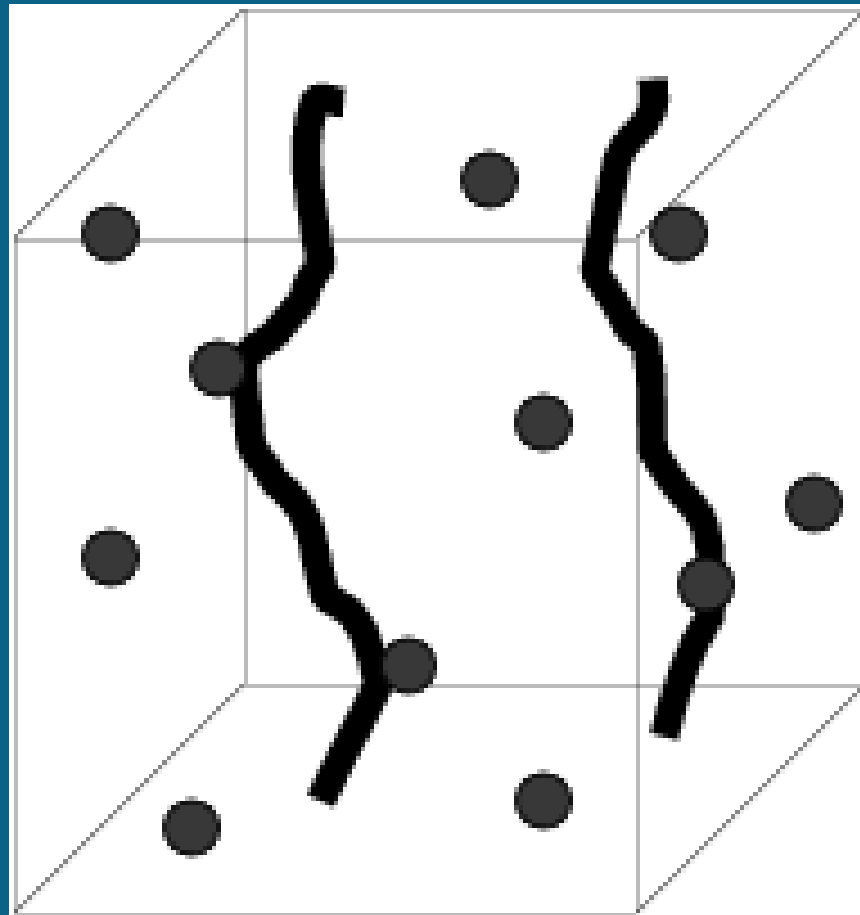


Vortex Interactions





Random Pinning





Vortex dynamics model

$$\overbrace{\eta \frac{\partial \mathbf{r}_{lz}(t)}{\partial t}}^{\text{friction}} = \overbrace{\mathbf{f}_{lz}^{int}(t) + \mathbf{f}_{lz}^{elast}(t)}^{\text{interactions}} + \overbrace{\mathbf{f}_{lz}^P(t)}^{\text{randompinning}} + \overbrace{\mathbf{f}_{lz}^T(t)}^{\text{thermalnoise}}, \quad (1)$$



Vortex dynamics model

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- \mathbf{r}_{lz} : 2D position of l -th vortex in z -th plane.
- Friction force: $\eta = \frac{\Phi_0 H_c 2d}{c^2 \rho_n}$



Vortex dynamics model

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- \mathbf{r}_{lz} : 2D position of l -th vortex in z -th plane.
- Friction force: $\eta = \frac{\Phi_0 H_c 2d}{c^2 \rho_n}$
- In-plane interaction force: $\mathbf{f}_{lz}^{int} = -\nabla_{\mathbf{r}_{lz}} U^p[\{\mathbf{r}_{lz}(t)\}]$ with

$$U^p[\{\mathbf{r}_{lz}(t)\}] = \sum_{l' \neq l} U \left(\frac{|\mathbf{r}_{lz'} - \mathbf{r}_{lz}|}{\lambda} \right), \quad (2)$$

with

$$U(x) = \frac{\Phi_0}{8\pi^2 \lambda^2} K_0(x), \quad (3)$$



- Elastic (inter-plane) force: $\mathbf{f}_{lz}^{elast} = -\nabla_{\mathbf{r}_{lz}} U^e[\{\mathbf{r}_{lz}(t)\}]$ with

$$U^e[\{\mathbf{r}_{lz}(t)\}] = V \left(\frac{|\mathbf{r}_{(z+1)l} - \mathbf{r}_{lz}|}{2r_g} \right) + V \left(\frac{|\mathbf{r}_{(z-1)l} - \mathbf{r}_{lz}|}{2r_g} \right), \quad (4)$$

$$V(|\mathbf{x}|) = \begin{cases} 2c_j(|\mathbf{x}| - 1) & \text{for } |\mathbf{x}| \geq 1 \\ c_j(|\mathbf{x}|^2 - 1) & \text{for } |\mathbf{x}| < 1, \end{cases} \quad (5)$$

with $\mathbf{x} = |\mathbf{r}_{z'l} - \mathbf{r}_{lz}|/2r_g$, $z' = z \pm 1$,
 $r_g = d\xi_{ab}/\xi_c$, d interplane distance, and

$$c_j = \frac{\epsilon_0 a}{\pi d} \left(1 + \ln \frac{\lambda}{a} \right)$$



- Random pinning force:

$$\mathbf{f}_{lz}^P = -\nabla_{\mathbf{r}_{lz}} \int d^2\mathbf{r}' u(\mathbf{r}', z) p(|\mathbf{r}_{lz} - \mathbf{r}'|), \quad (6)$$

$$\langle u(\mathbf{r}, z) u(\mathbf{r}', z') \rangle = \gamma \delta(\mathbf{r} - \mathbf{r}') \delta_{zz'},$$

and

$$p(r) = \frac{2\xi^2}{r^2 + 2\xi^2}.$$



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- Thermal (Langevin) force:

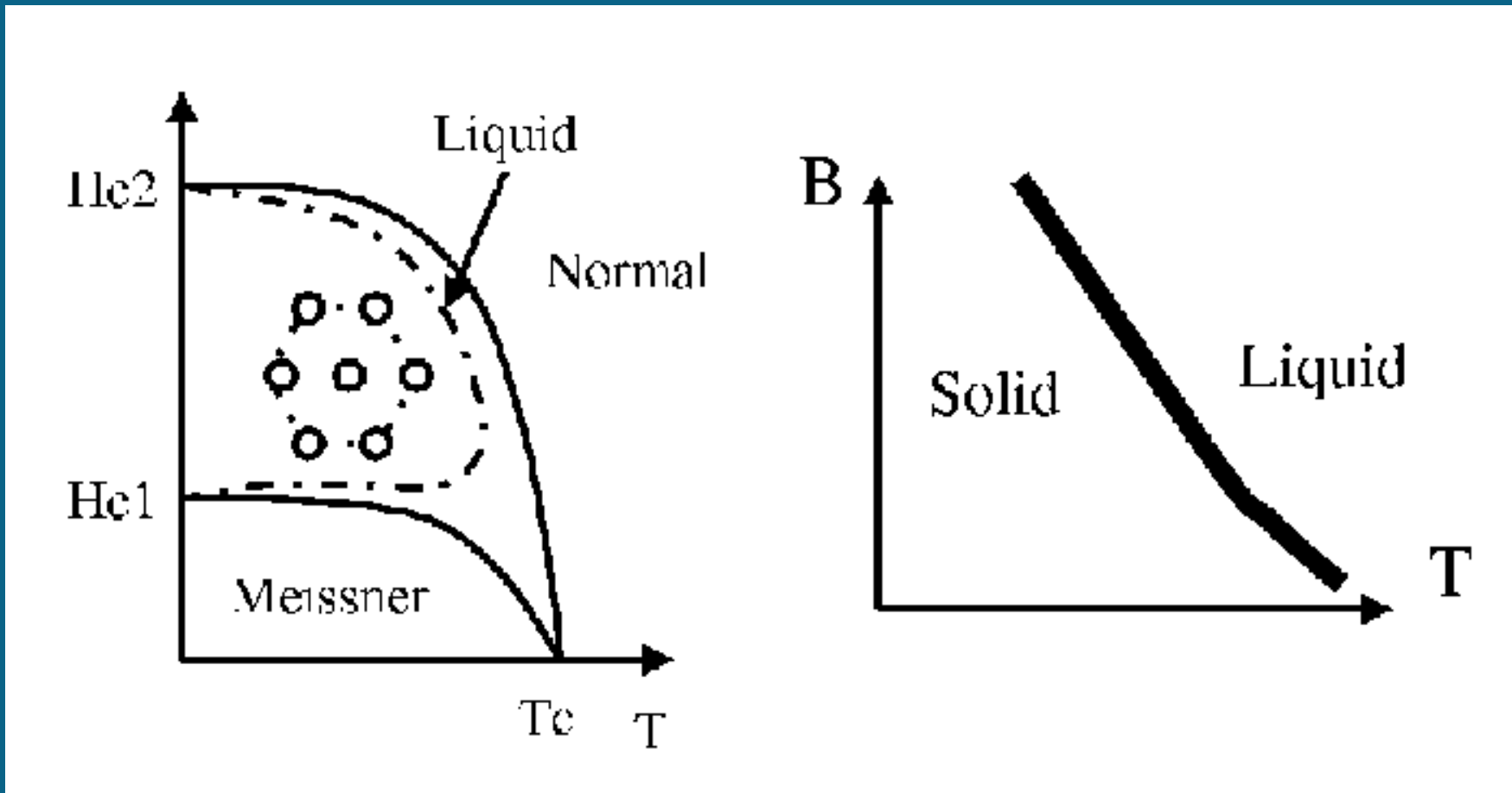
$$\langle f_{lz,\mu}^T(t) \rangle = 0,$$

$$\langle f_{lz,\mu}^T(t) f_{z'l',\mu}^T(t') \rangle = 2\eta k_B T \delta(t - t') \delta_{zz'} \delta_{ll'} \delta_{\mu\mu'},$$

$$\mu, \mu' = x, y$$

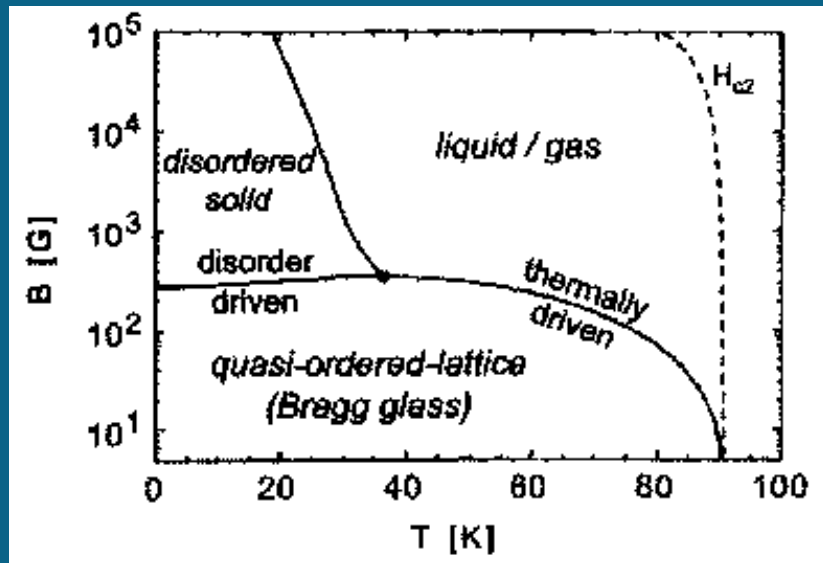


Vortex Phase diagram without disorder

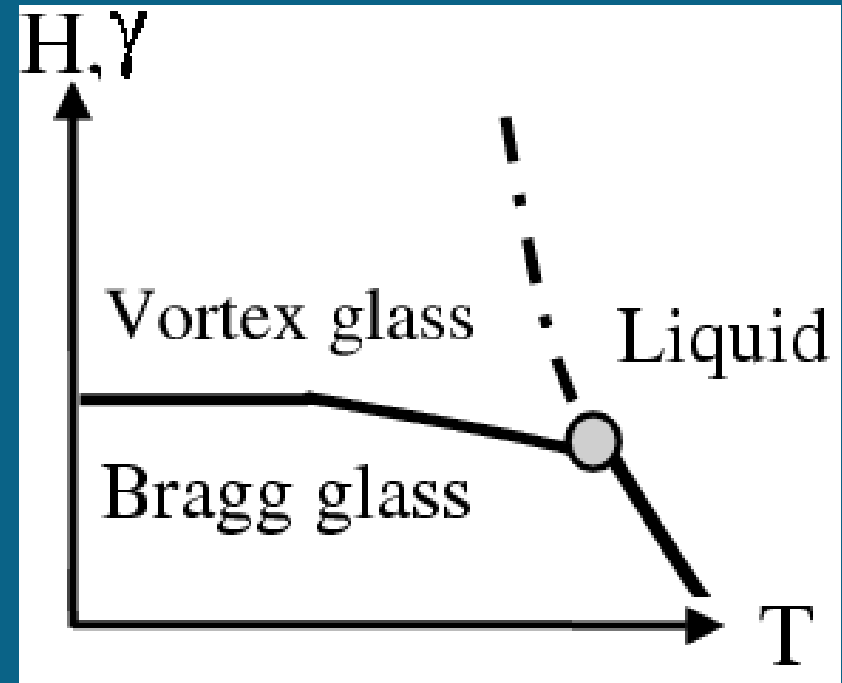




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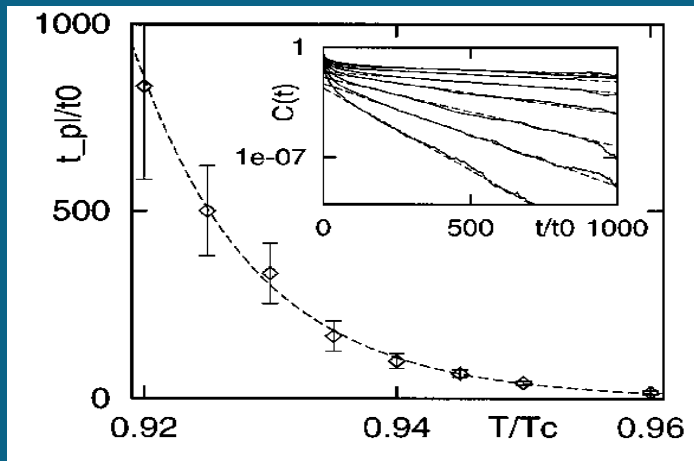
From experimental data in BiSrCaCuO



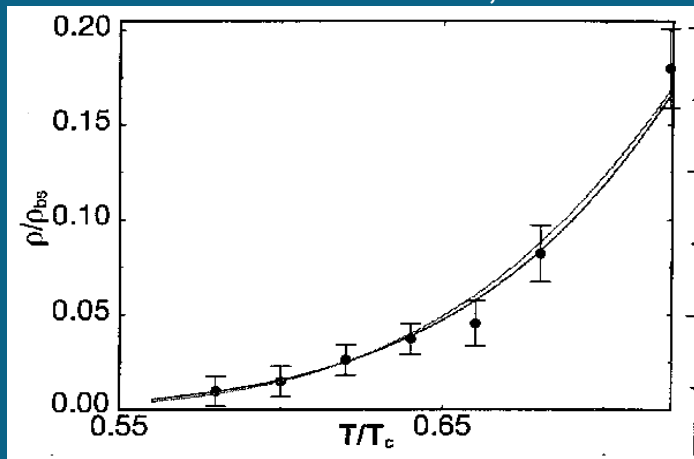
Schematic Phase Diagram



Vortex Glass transition



van Otterlo et al, PRL 1998



Reichhardt et al, PRL 2000

- $C(t) = e^{-\langle [r_{lz}(t) - r_{lz}(0)]^2 \rangle / a^2} \sim e^{-t/t_{pl}}$

- Scaling theory : $\rho \sim |T - T_g|^{\nu(z-1)}$
 $\nu = 1.3 - 2, z = 3.1 - 6.5$

- Volger-Fulcher law : $\rho \sim e^{-1/(T-T_g)}$



Out of equilibrium systems

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★ *Very slow relaxation* $\tau_{equilibration} \gg \tau_{observation}$



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 - ★ *External forces:* $f(x) \neq -\partial V/\partial x, f(x, t)$



Out of equilibrium systems

- Out of equilibrium:
 - ★ *Very slow relaxation* $\tau_{equilibration} \gg \tau_{observation}$
 - ★ *External forces:* $f(x) \neq -\partial V/\partial x, f(x, t)$
- Examples: *spin glasses, window glasses polymers, colloids, gels, dense liquides, granular matter, sheared fluids, etc,..... Vortex Matter?*



Out of equilibrium dynamics of glasses

Energy flow, partial equilibration, and effective temperatures in systems with slow dynamics,
Cugliandolo, Kurchan, Peliti, PRE, 1997.



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- Fluctuations: $C(t + t_w, t_w) = \langle O(t + t_w)O(t_w) \rangle$, $\langle O(t) \rangle = 0$



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Energy flow, partial equilibration, and effective temperatures in systems with slow dynamics,
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- Fluctuations: $C(t + t_w, t_w) = \langle O(t + t_w)O(t_w) \rangle, \langle O(t) \rangle = 0$
- Response: $\chi(t + t_w, t_w) = \lim_{\epsilon \rightarrow 0} \frac{\langle O(t) \rangle_{\epsilon} - \langle O(t) \rangle_0}{\epsilon}, (H_{\epsilon} = H - \epsilon(t)O)$



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- Equilibrium: *Fluctuation - dissipation theorem* ($t_w \gg \tau_{eq}$)

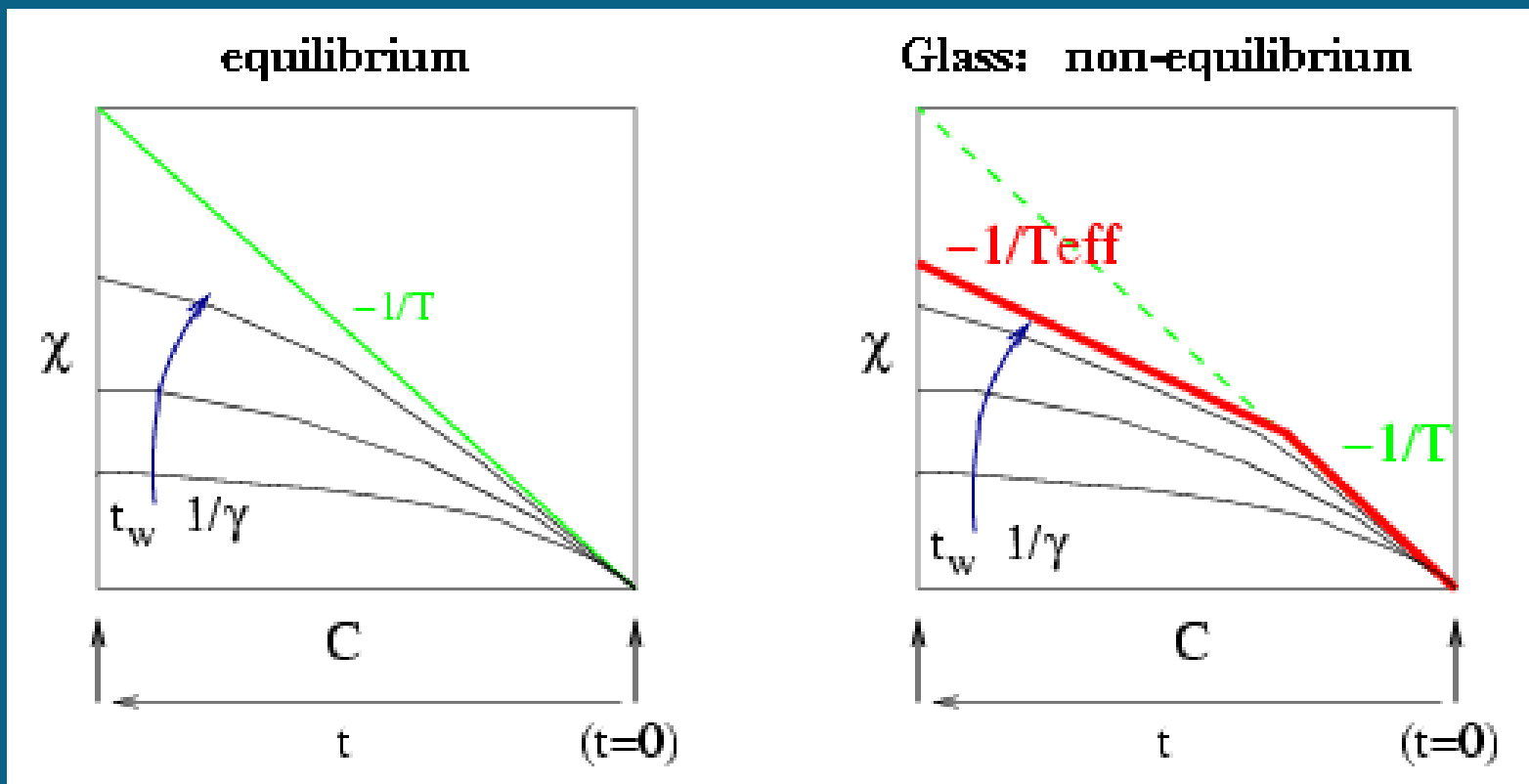
$$\chi(t) = \frac{1}{k_B T} [C(0) - C(t)]$$



Out of equilibrium dynamics of glasses

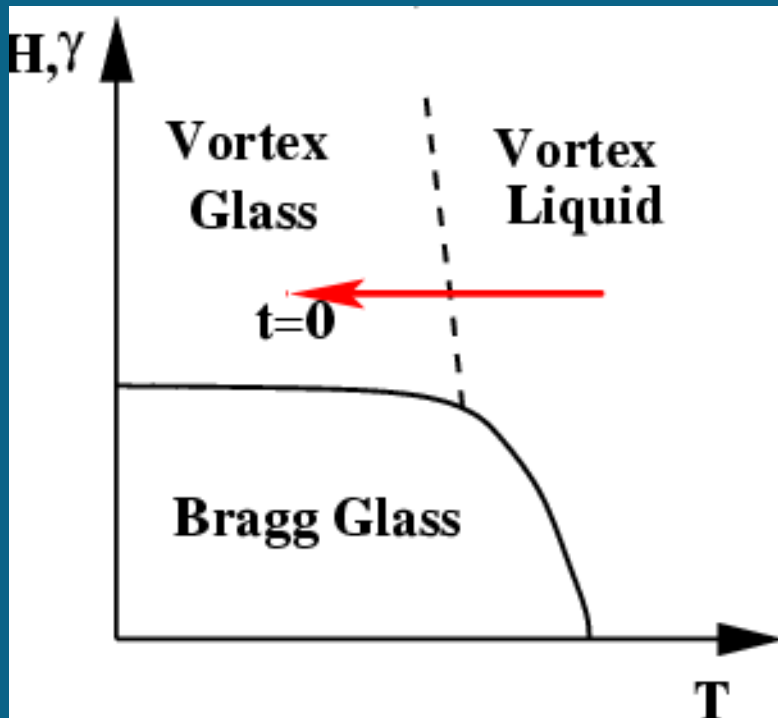
Energy flow, partial equilibration, and effective temperatures in systems with slow dynamics,
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- Fluctuation - dissipation relations

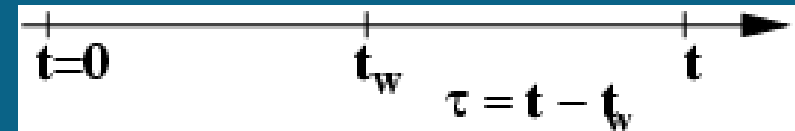




Simulations



- 56 pancakes per plane
- 50 planes
- Quench from $T_{start} = 0.3 > T_g$ to T





Correlation functions measured

- Dynamic wandering

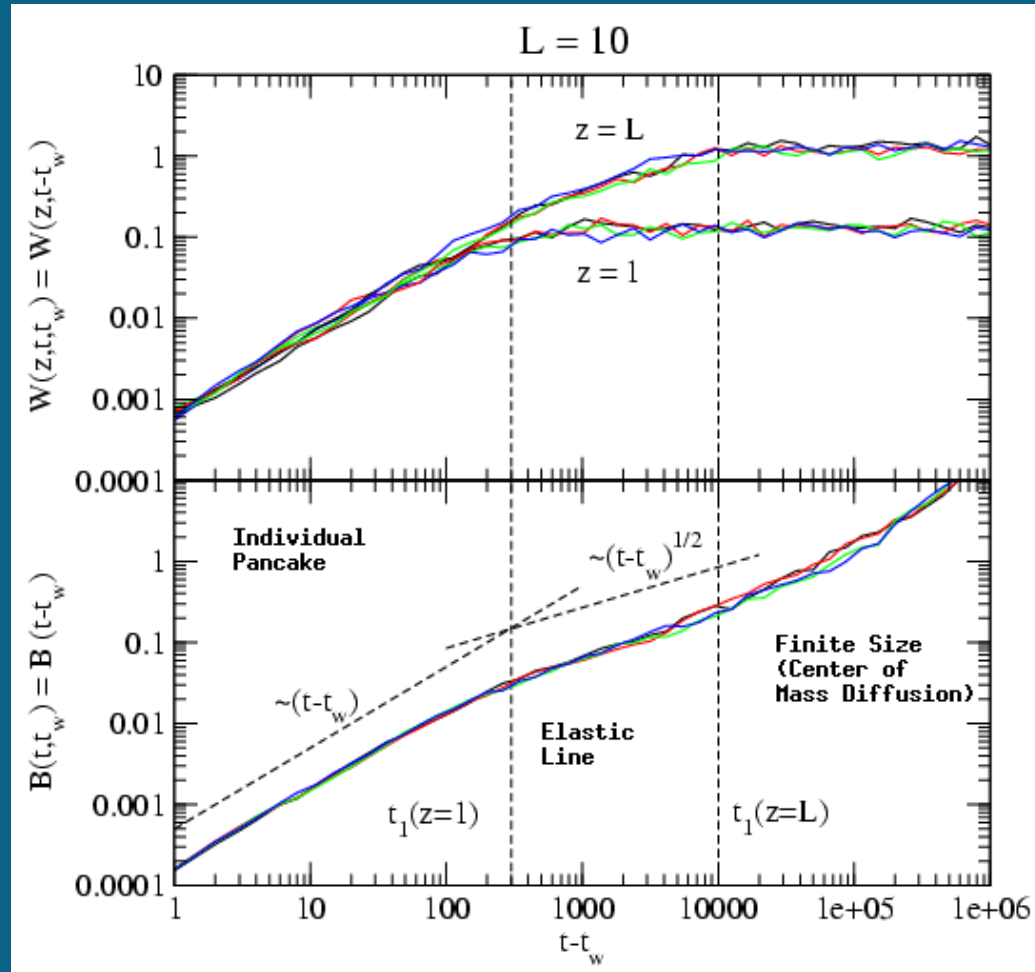
$$W(z, t, t_w) = \left\langle \frac{1}{N_{xy}} \sum_l \left| [\mathbf{r}_{lz}(t) - \mathbf{r}_{l0}(t)] - [\mathbf{r}_{lz}(t_w) - \mathbf{r}_{l0}(t_w)] \right|^2 \right\rangle$$

- Mean square displacement (MSD)

$$B(t, t_w) = \left\langle \frac{1}{N} \sum_{lz} [x_{lz}(t) - x_{lz}(t_w)]^2 \right\rangle$$

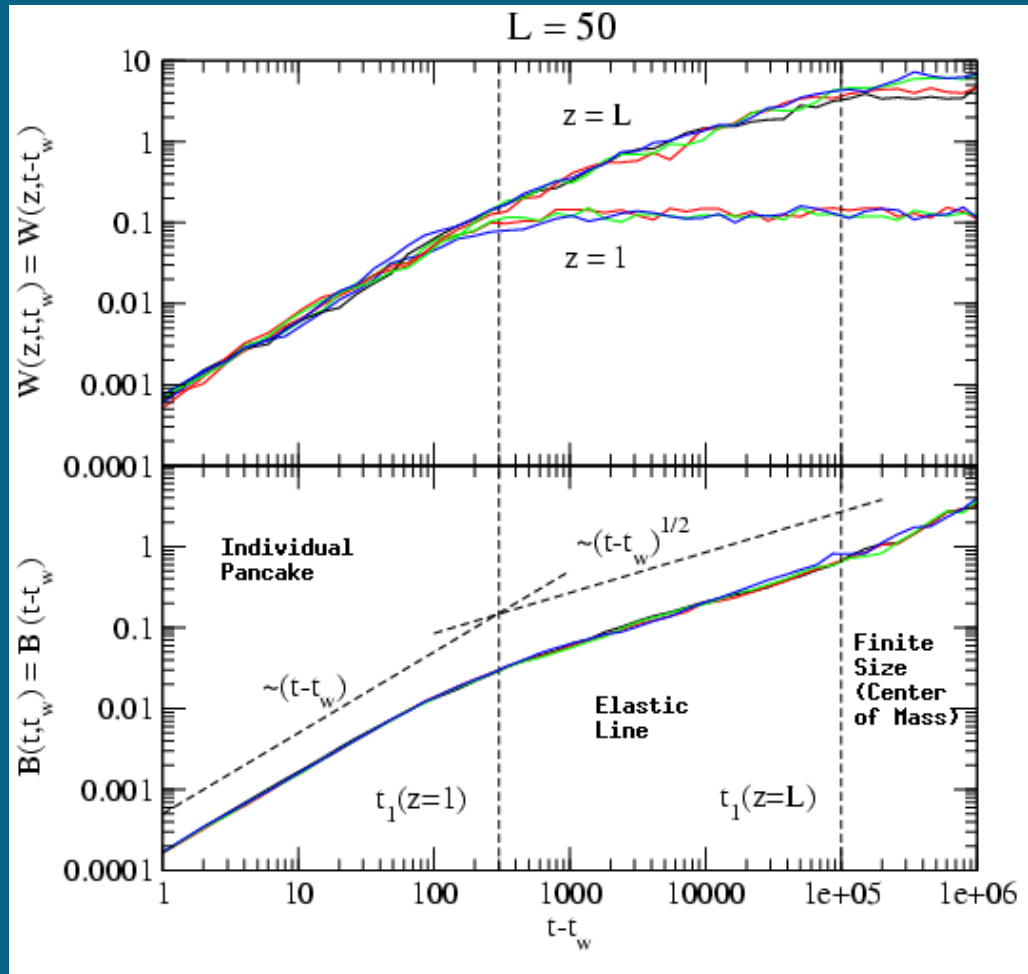


Single Vortex Line without disorder (Vortex Liquid)





Single Vortex Line without disorder (Vortex Liquid)





Finite Size effect

$$\tau = t - t_w$$

$$W(z, \tau) \sim \begin{cases} \tau & \tau < t_{1,(z)} \\ z^{2\zeta_T} & \tau > t_{1,(z)} \end{cases}$$

$\zeta = \zeta_T = 1/2$ (thermal line wandering)

For $t > t_{1,(z=N_z)}$ finite size effect



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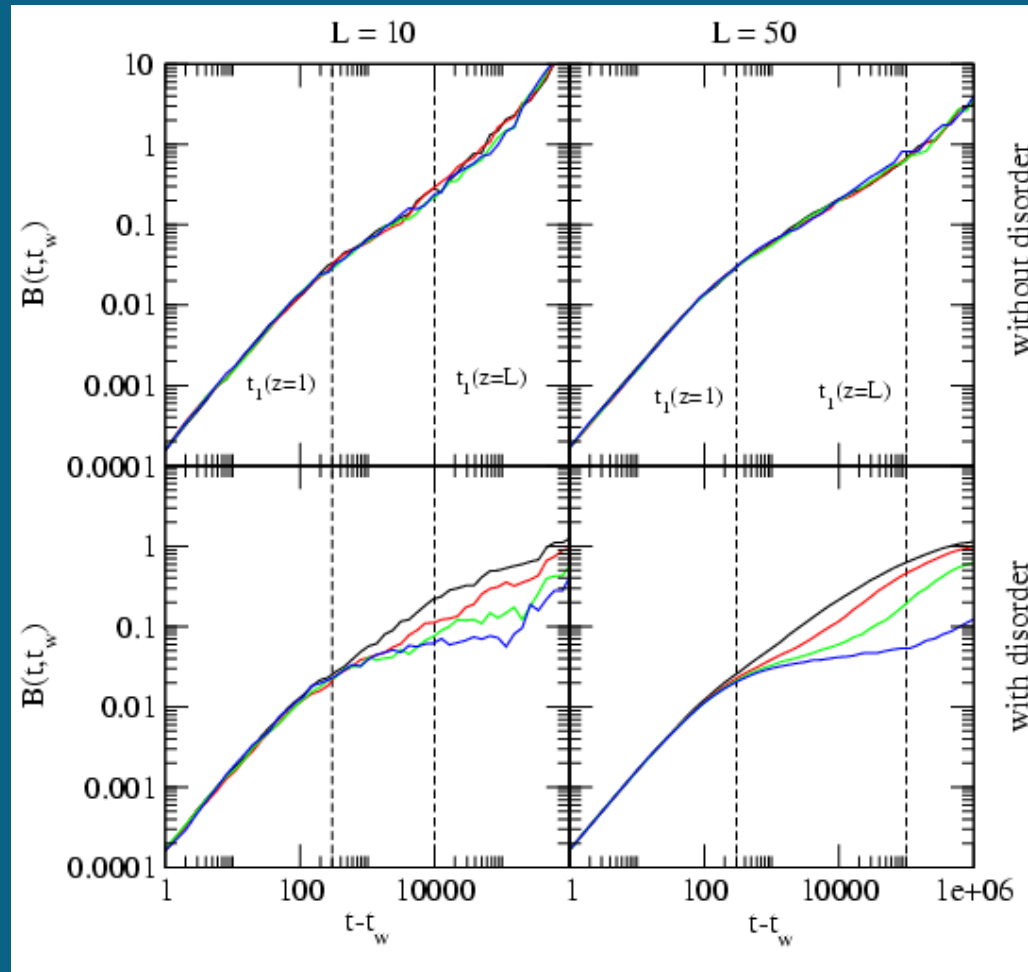
$\zeta = \zeta_T = 1/2$ (thermal line wandering)

For $t > t_{1,(z=N_z)}$ finite size effect

$$B(\tau) \sim \begin{cases} \tau & \tau < t_{1,(z=1)} \\ \tau^{1/2} & t_{1,(z=1)} < \tau < t_{1,(z=N_z)} \\ \tau & t_{1,(z=N_z)} < \tau \end{cases} \begin{array}{l} \text{single pancake} \\ \text{elastic line} \\ \text{center of mass} \end{array}$$



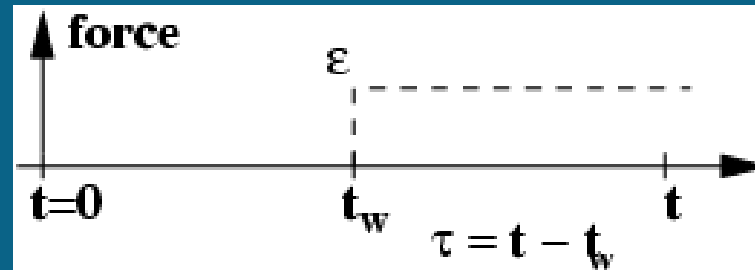
Single Vortex Line with Disorder: Aging





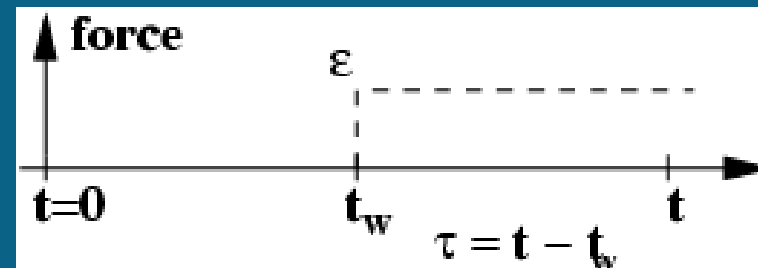
Response Function

- Integrated Response



Response Function

- Integrated Response



Random force for $t \geq t_w$: $\mathbf{f}_{lz} = \epsilon s_{lz} \hat{\mathbf{x}}$,

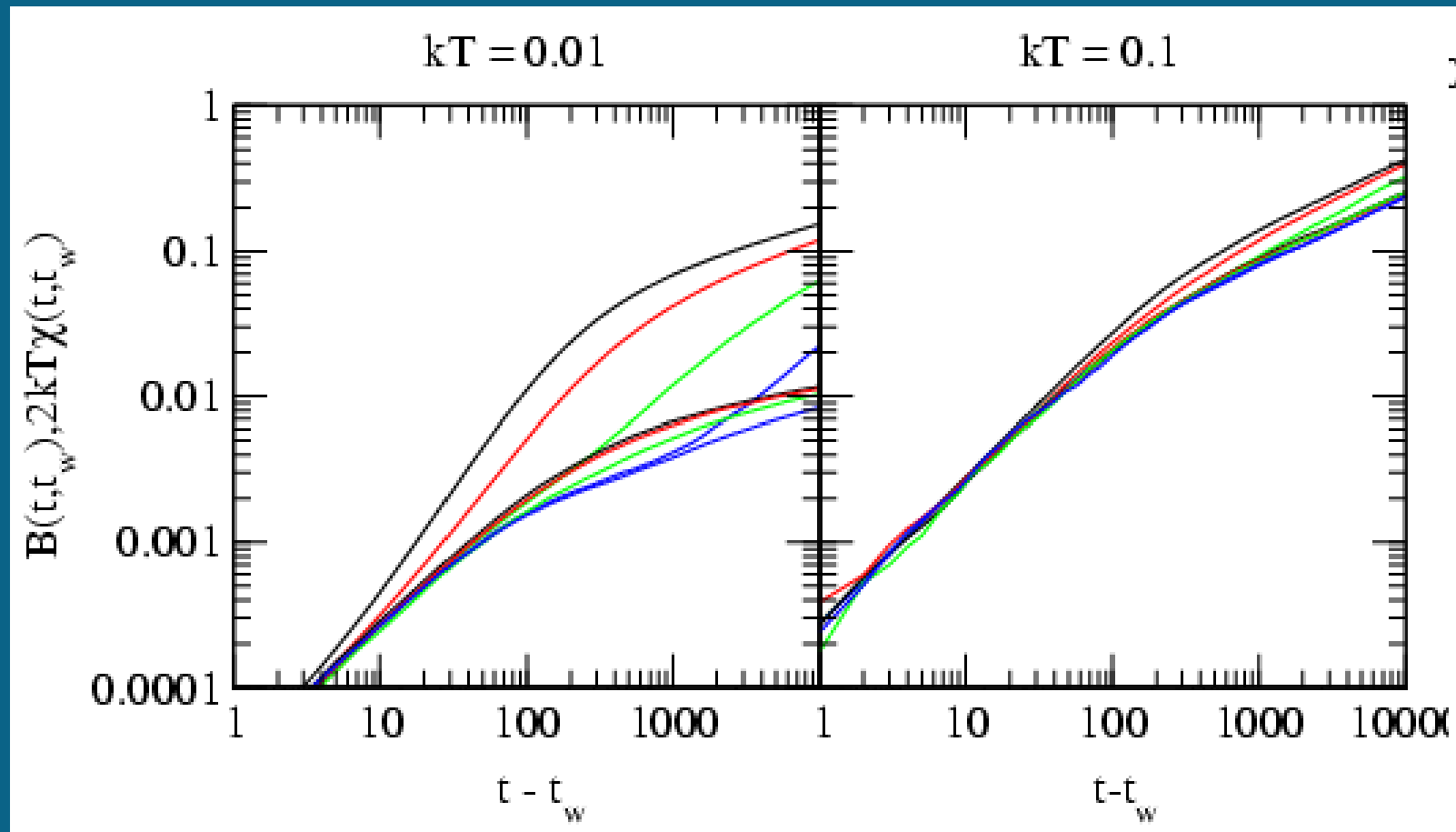
ϵ : perturbation strength, $s_{lz} = \pm 1$, random

$$\chi(t, t_w) = \frac{\langle \frac{1}{N} \sum_{lz} s_{lz} x_{lz}(t) \rangle_{\epsilon} - \langle \frac{1}{N} \sum_{lz} x_{lz}(t) \rangle_{\epsilon=0}}{\epsilon}$$

x_{lz} , x -component of vortex position,



Single Vortex Line with Disorder: Aging





Multiplicative Scaling

Yoshino, PRL 1998, Directed Polimer in Random Media (1+1)

Define

$$B(t, t_w) = \tilde{B}(\tilde{t})t^{-\alpha}$$
$$2kT\chi(t, t_w) = \tilde{\chi}(\tilde{t})t^{-\alpha}$$

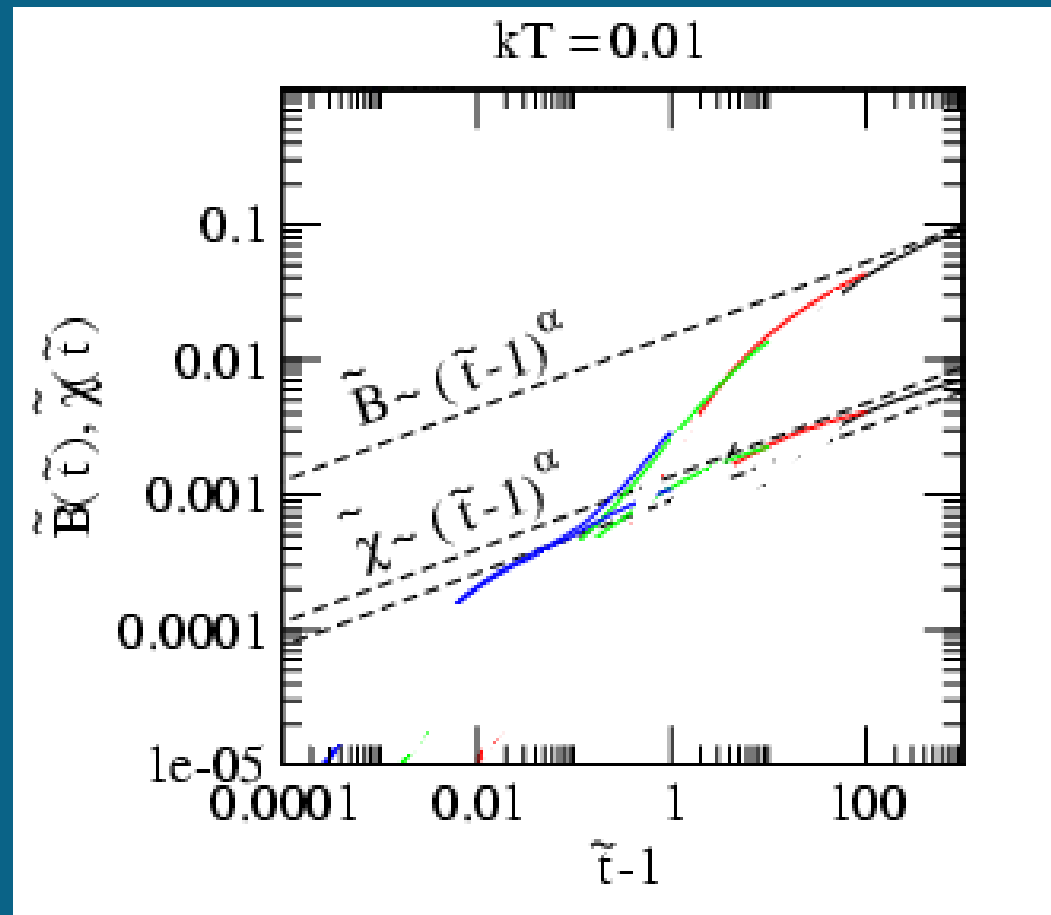
with $\tilde{t} = t/t_w$ and

$$\tilde{B}(\tilde{t}) = \begin{cases} c_1(T)(\tilde{t} - 1)^{-\alpha(T)} & \tilde{t} \ll 1 \\ c_2(T)(\tilde{t} - 1)^{-\alpha(T)} & \tilde{t} \gg 1 \end{cases}$$

$$\tilde{\chi}(\tilde{t}) = \begin{cases} c_1(T)(\tilde{t} - 1)^{-\alpha(T)} & \tilde{t} \ll 1 \\ y(T)c_2(T)(\tilde{t} - 1)^{-\alpha(T)} & \tilde{t} \gg 1 \end{cases}$$



Single Vortex Line with Disorder: Scaling





Violation of FDT

$$\chi(t, t_w) = \frac{y(t, t_w)}{2T} B(t, t_w)$$

with $y(t, t_w) = \frac{T}{T_{\text{ef}}(t, t_w)}$



Violation of FDT

$$\chi(t, t_w) = \frac{y(t, t_w)}{2T} B(t, t_w)$$

with $y(t, t_w) = \frac{T}{T_{\text{ef}}(t, t_w)}$

- for $t < t_w$,

$$\tilde{\chi} = \tilde{B}$$

- for $t > t_w$

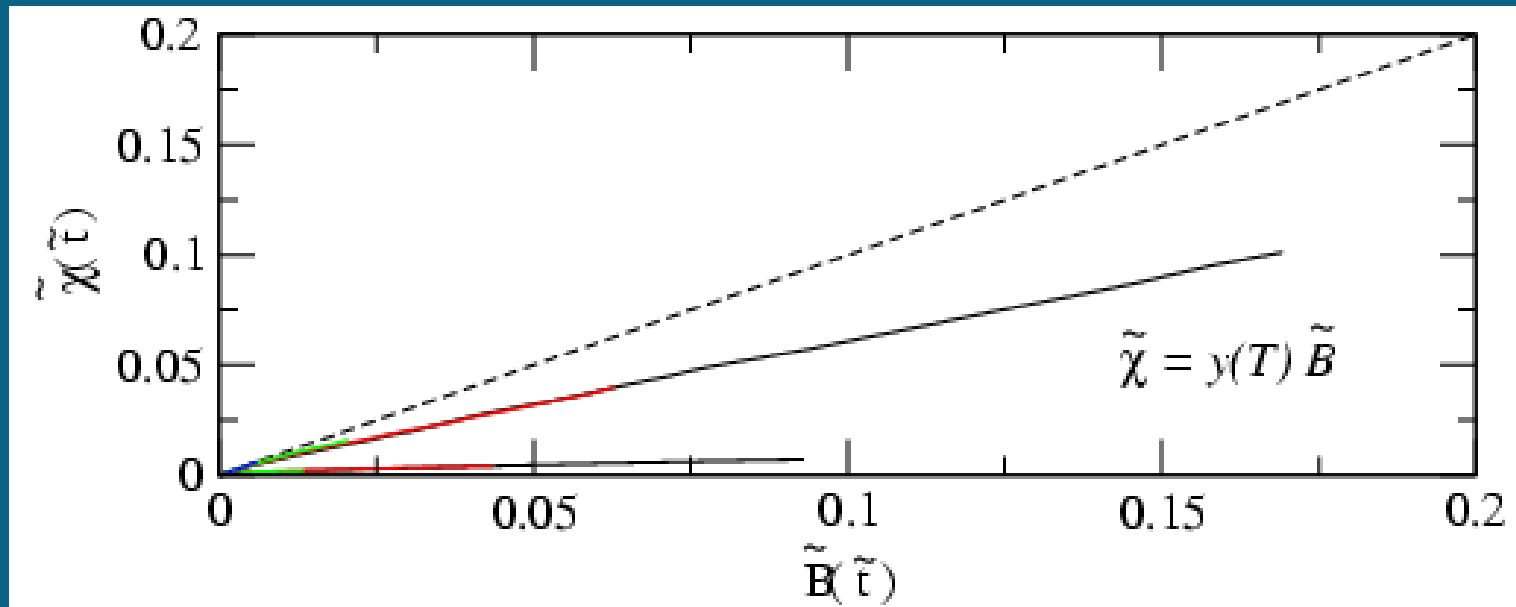
$$\tilde{\chi} = y(T) \tilde{B}$$

- Effective temperature:

$$y(T) = \frac{T}{T_{\text{eff}}}$$

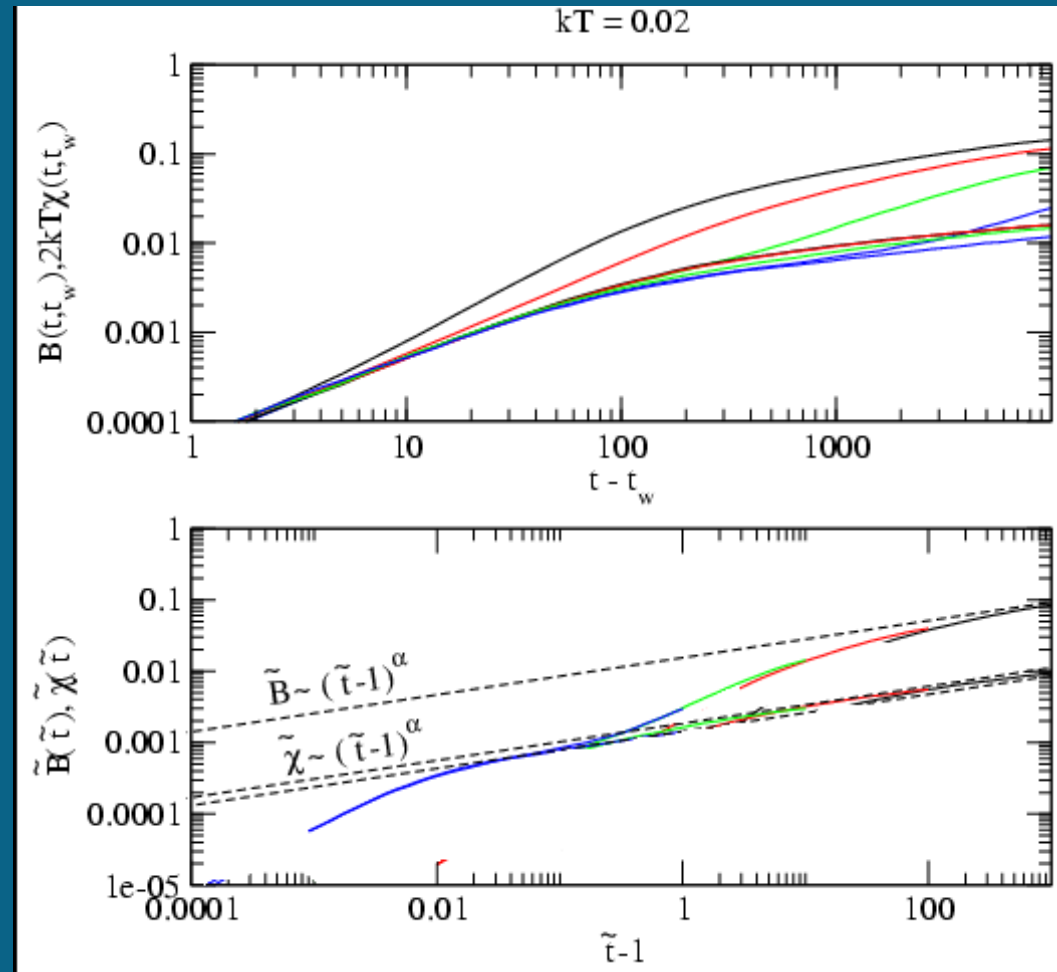


Single Vortex Line with Disorder: Violation of FDT



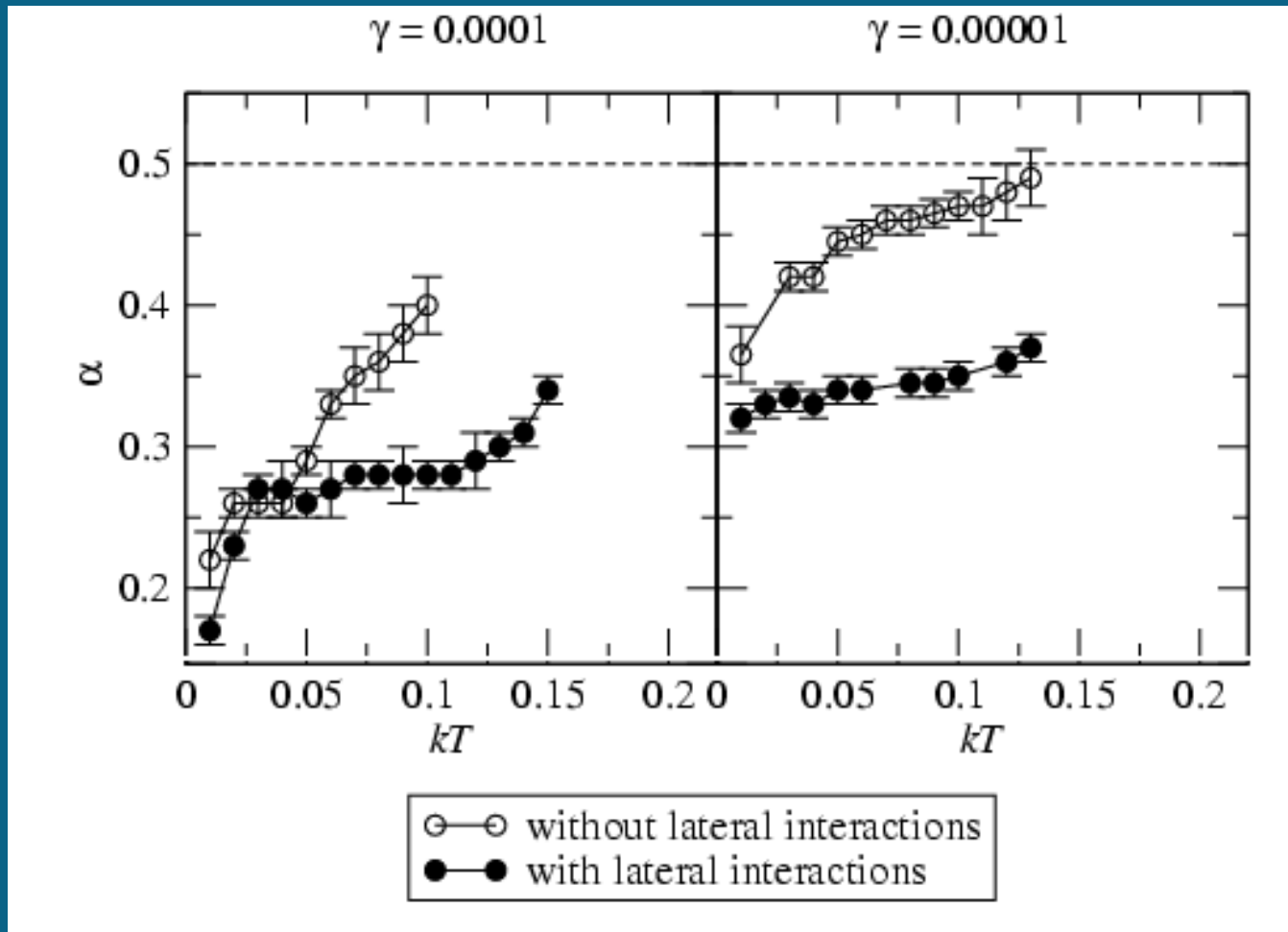


Vortex Glass: Aging and Scaling



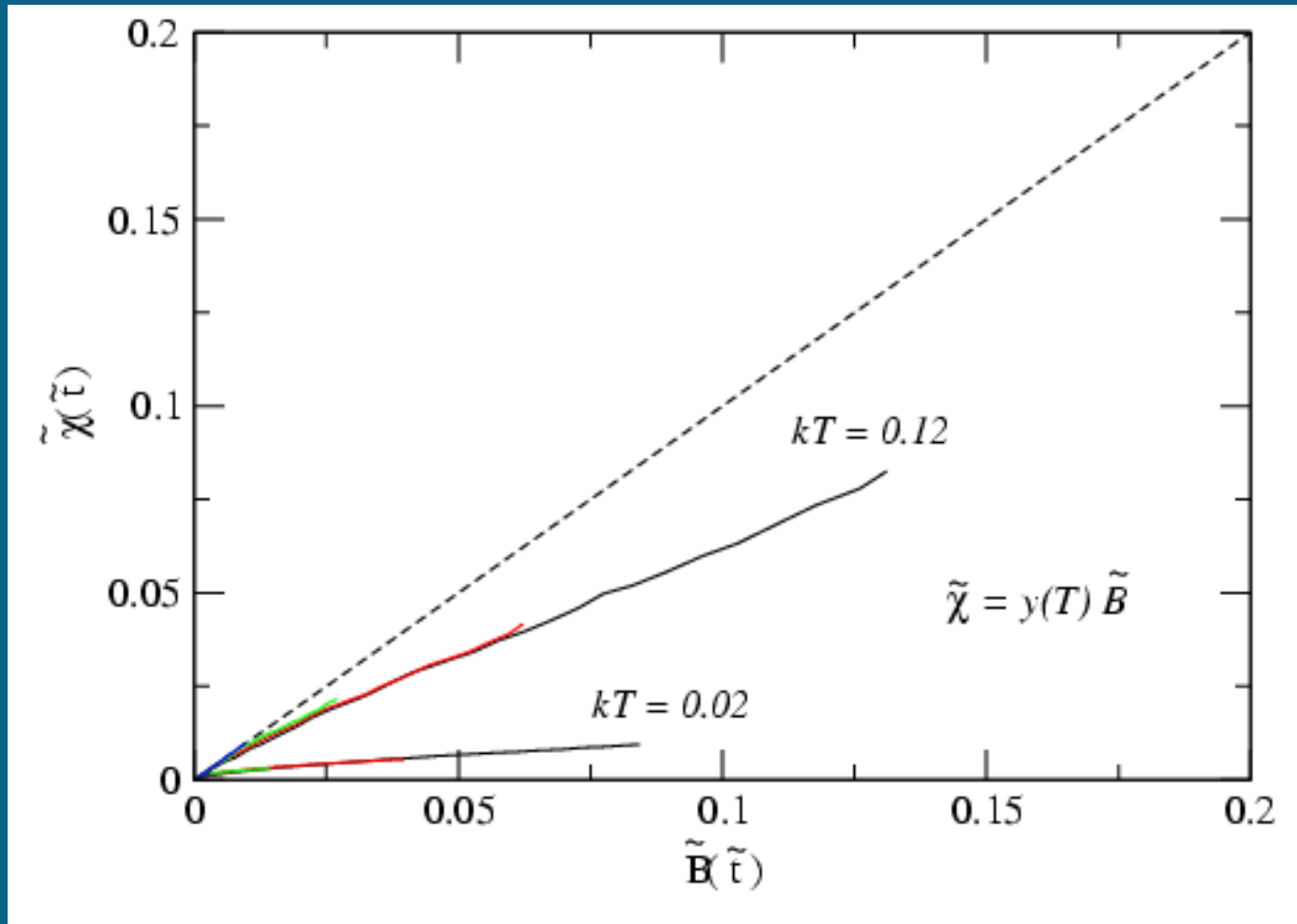


Scaling Exponents $\alpha(T)$



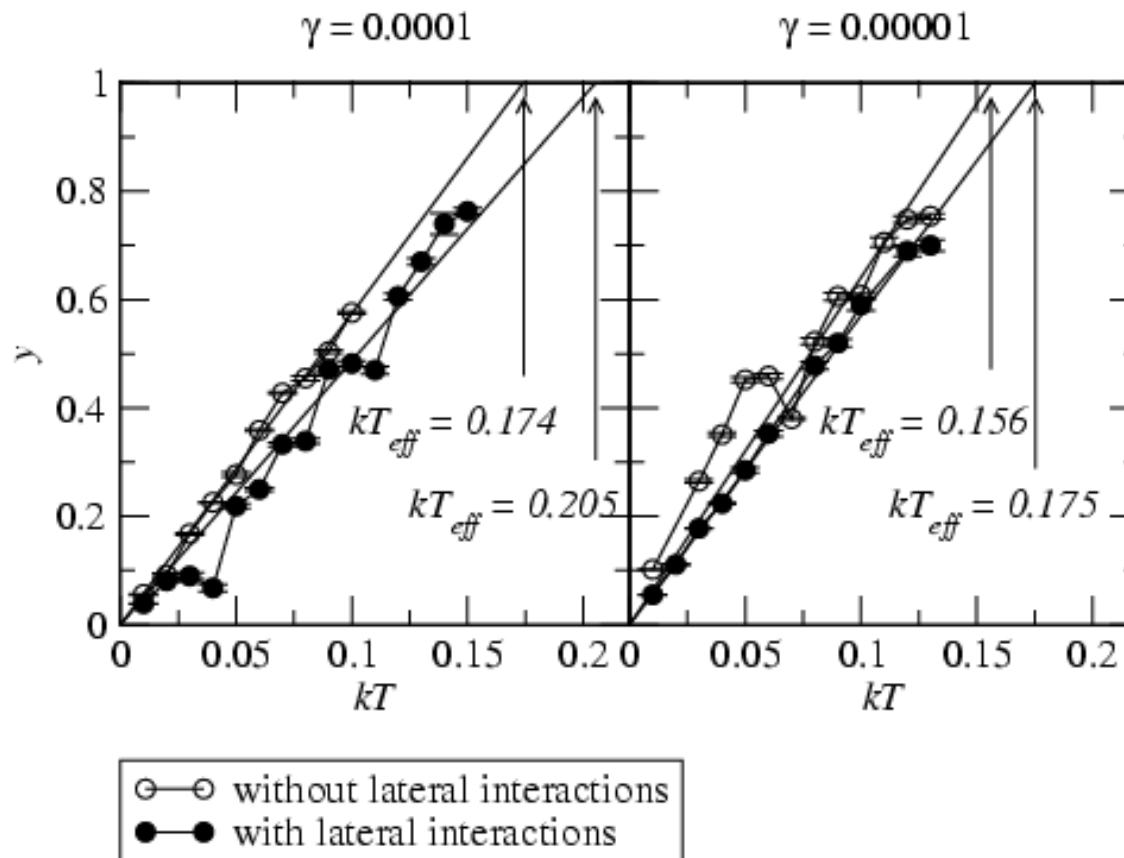


Vortex Glass: Violation of FDT





Effective Temperature



$$y(T) = \frac{T}{T_{eff}(T)}$$



Conclusions

- Slow dynamics dominated by elastic line relaxation
- Aging : multiplicative scaling (as polimers in random media)
- Violation of FDT after a quench to the vortex glass