

Topologically localized insulators

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in collaboration with
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New Recipes



Topologically localized insulators (TLI)

Stuff you need:

- Anderson localization



- Topology



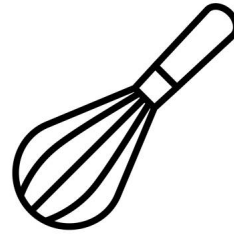
New Recipes



Topologically localized insulators (TLI)

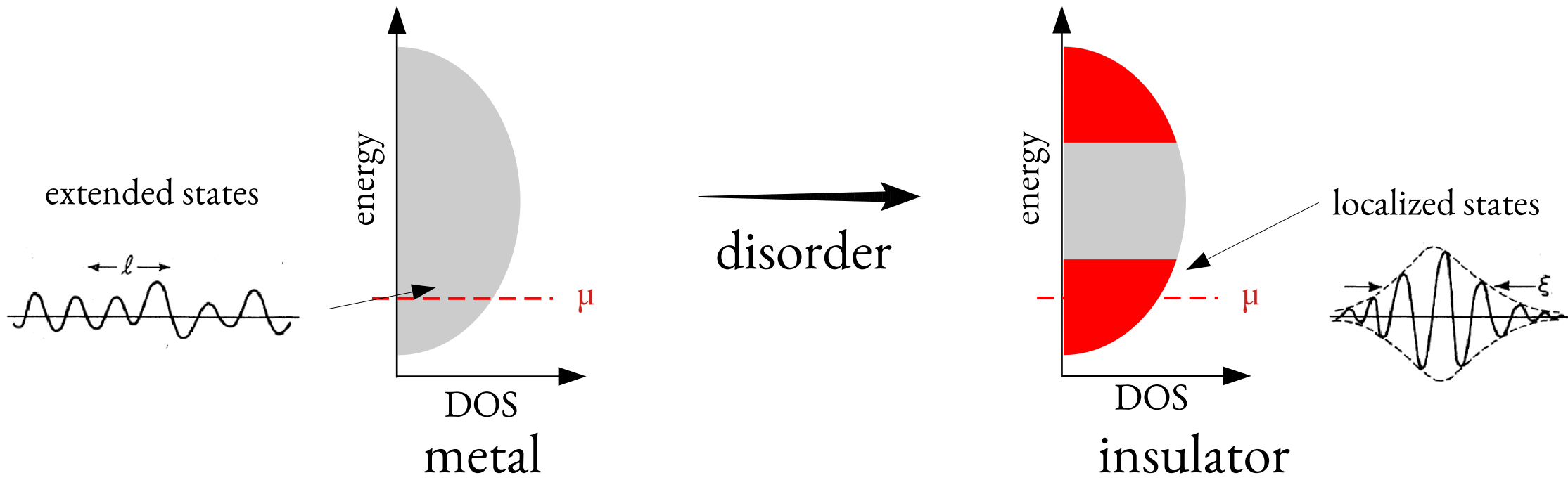
Stuff you need:

- Anderson localization



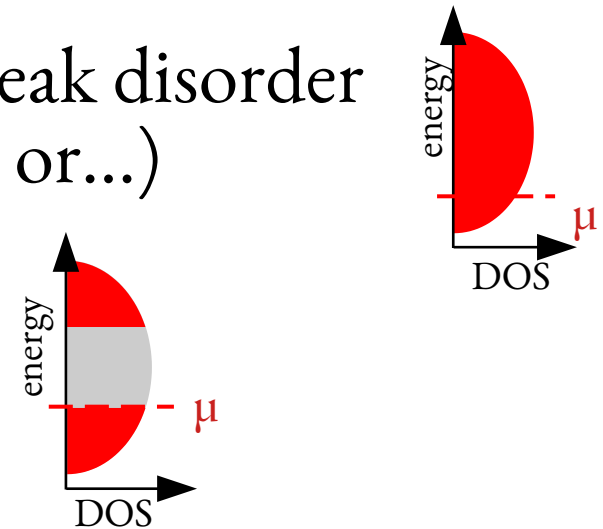
TLI

Stuff you need to do: ...



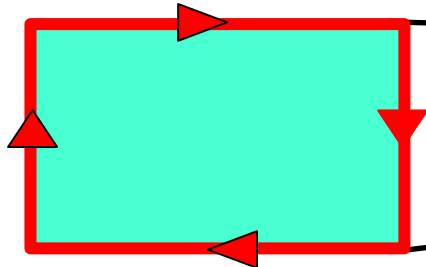
- in dimensions $d=1$ and 2 , localized for arbitrary weak disorder (exceptions in $d=2$ in presence of topology or...)

- $d=3$: disorder induced metal-insulator transition





chiral mode



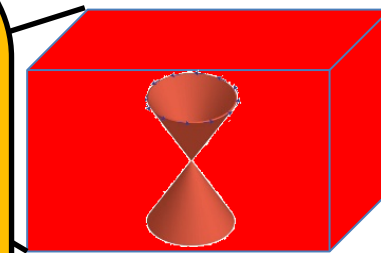
class	\mathcal{T}	\mathcal{P}	\mathcal{C}	$d=0$	$d=1$	$d=2$	$d=3$
A	0	0	0	\mathbb{Z}	0	\mathbb{Z}	0
AIII	0	0	1	0	\mathbb{Z}	0	\mathbb{Z}
AI	+	0	0	\mathbb{Z}	0	0	0
BDI	+	-	1	\mathbb{Z}_2	\mathbb{Z}	0	0

Tenfold-way

2d Quantum Hall insulator
or Chern insulator

Topological quantum material
 \updownarrow
 (robustly) quantized observable

helical surface mode

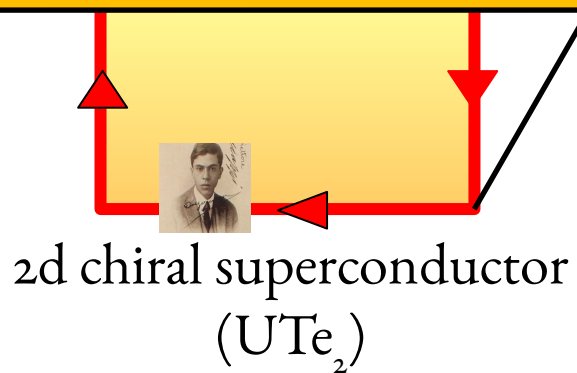


topological insulator
(Bi_2Se_3)

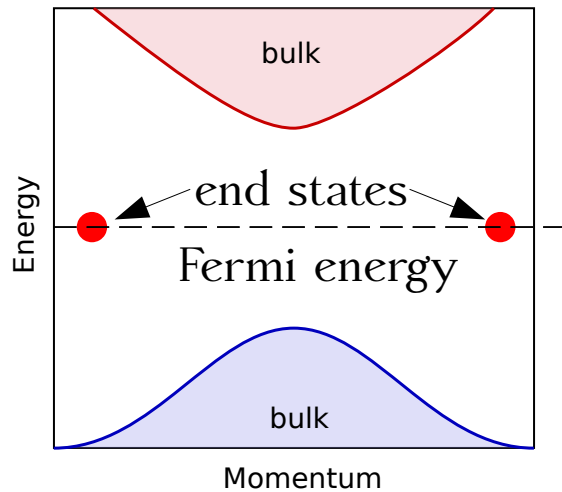
Majorana



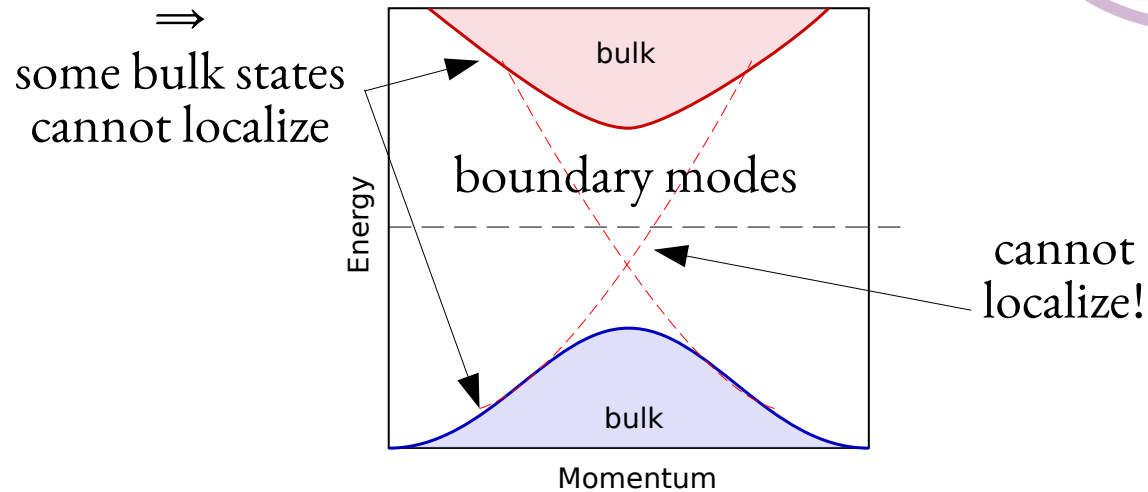
1d Kitaev chain
(InAs proximitized with SC)



Type A



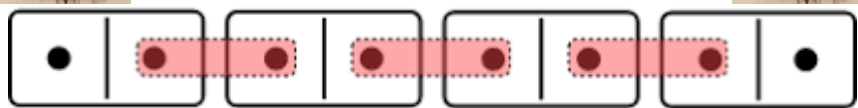
Type B



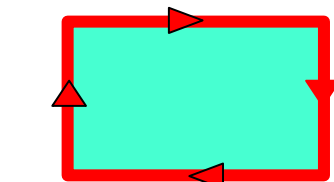
...the title of the talk says «insulators»!

(insulators & superconductors)

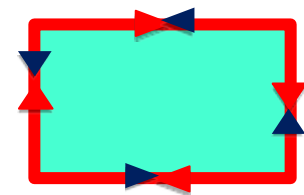
(only for superconductors)



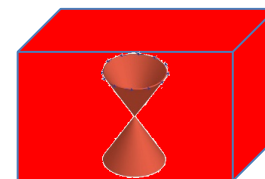
unpaired Majorana modes



Quantum Hall insulator

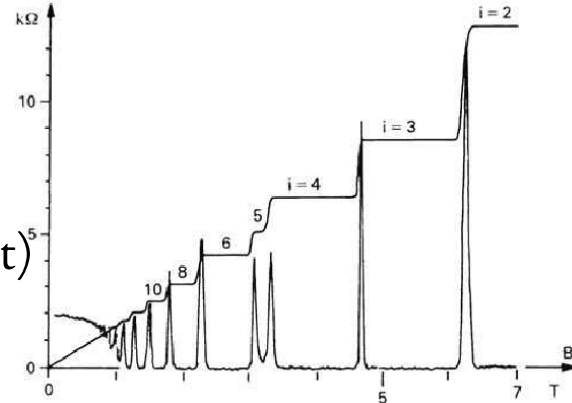
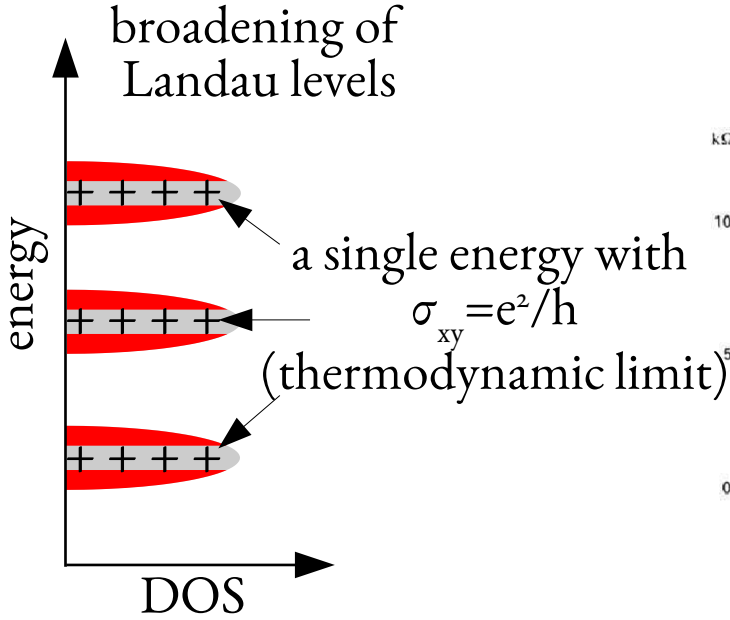


Quantum Spin Hall insulator

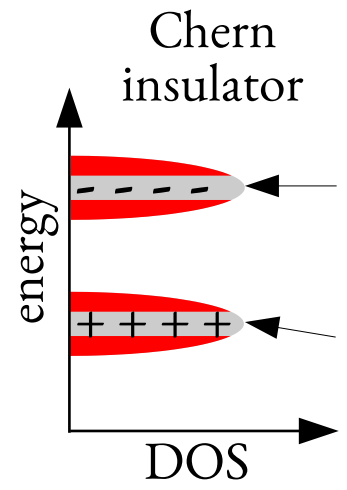


3d topological insulator

Quantum Hall/Chern insulator

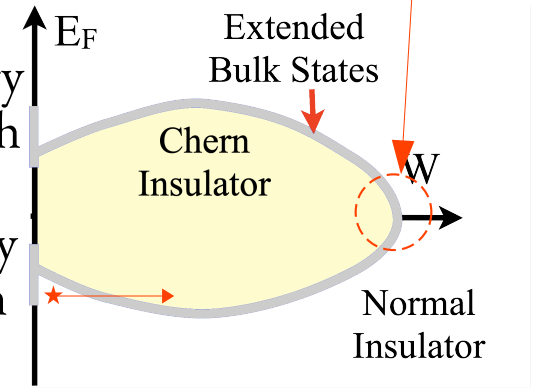


«pair annihilation»



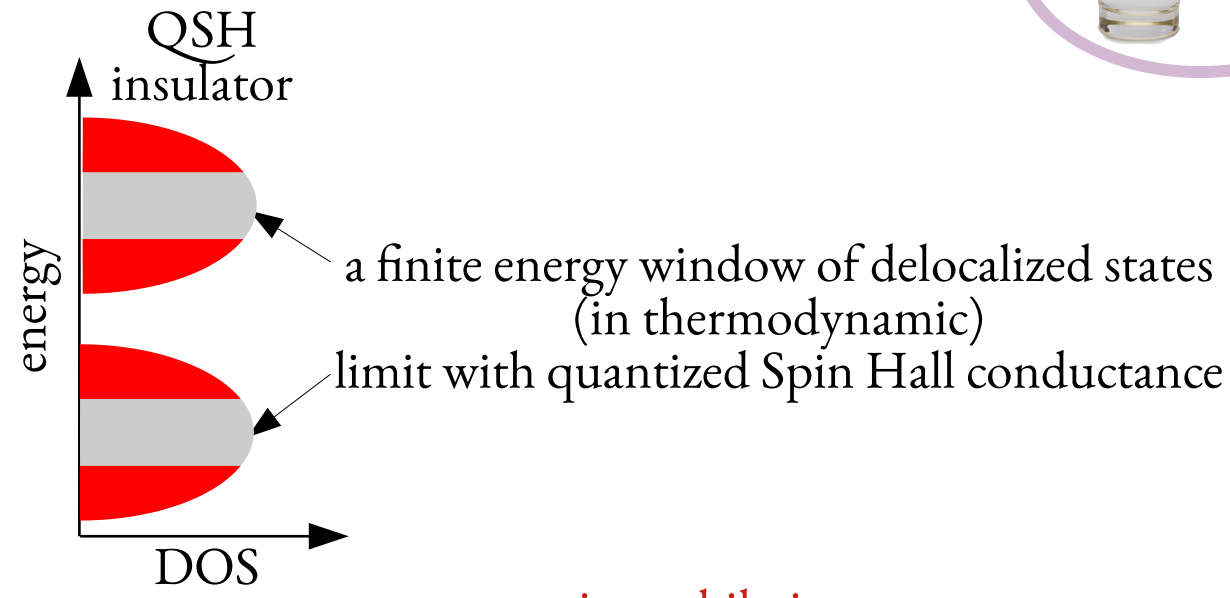
a single energy with $\sigma_{xy} = -e^2/h$

a single energy with $\sigma_{xy} = e^2/h$

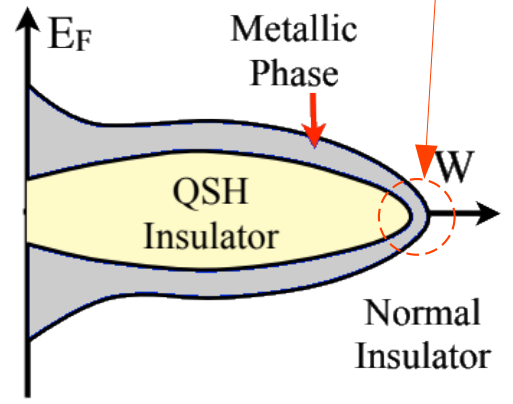


"Topological Anderson Insulator"

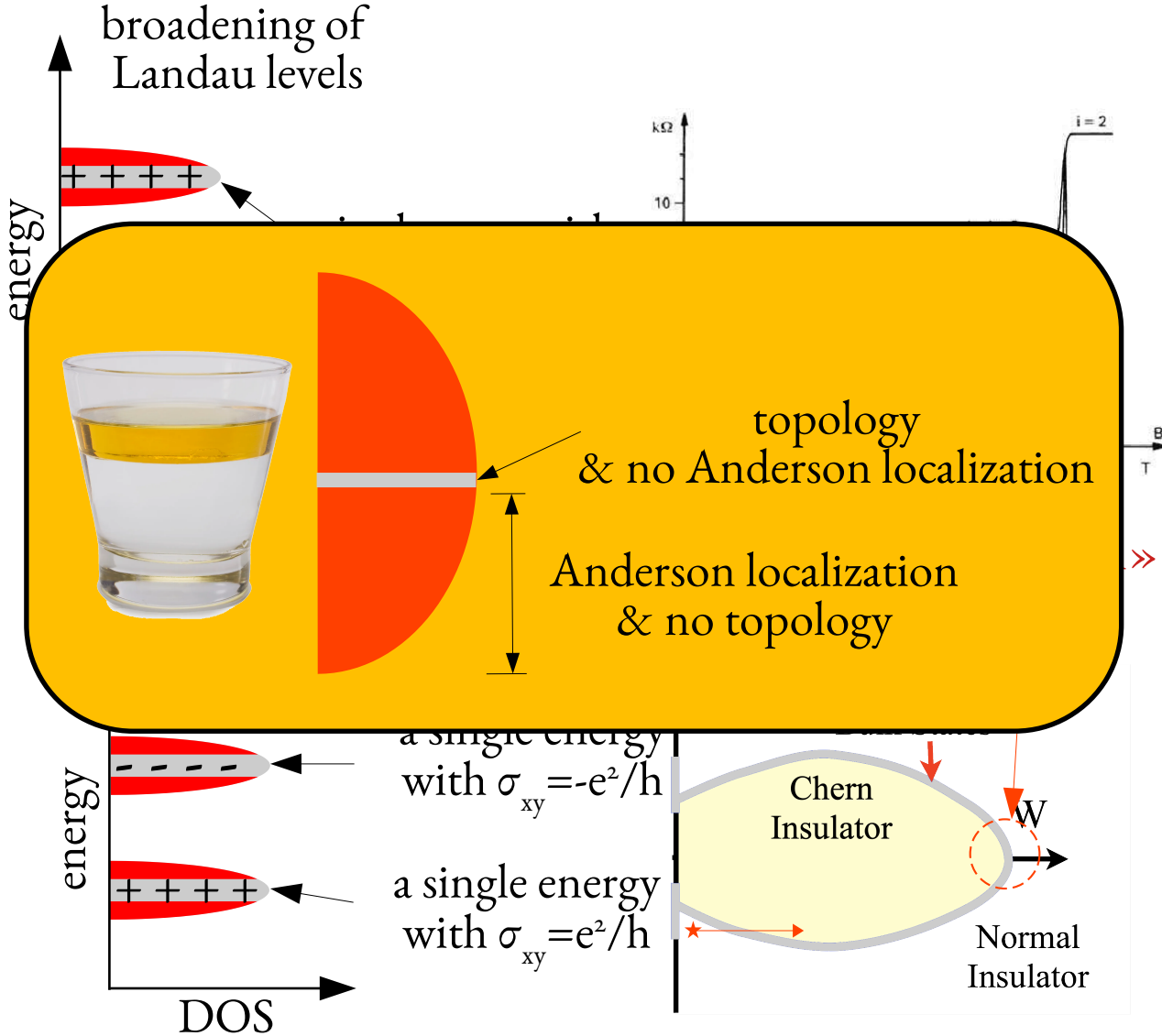
Quantum Spin Hall insulator



«pair annihilation»

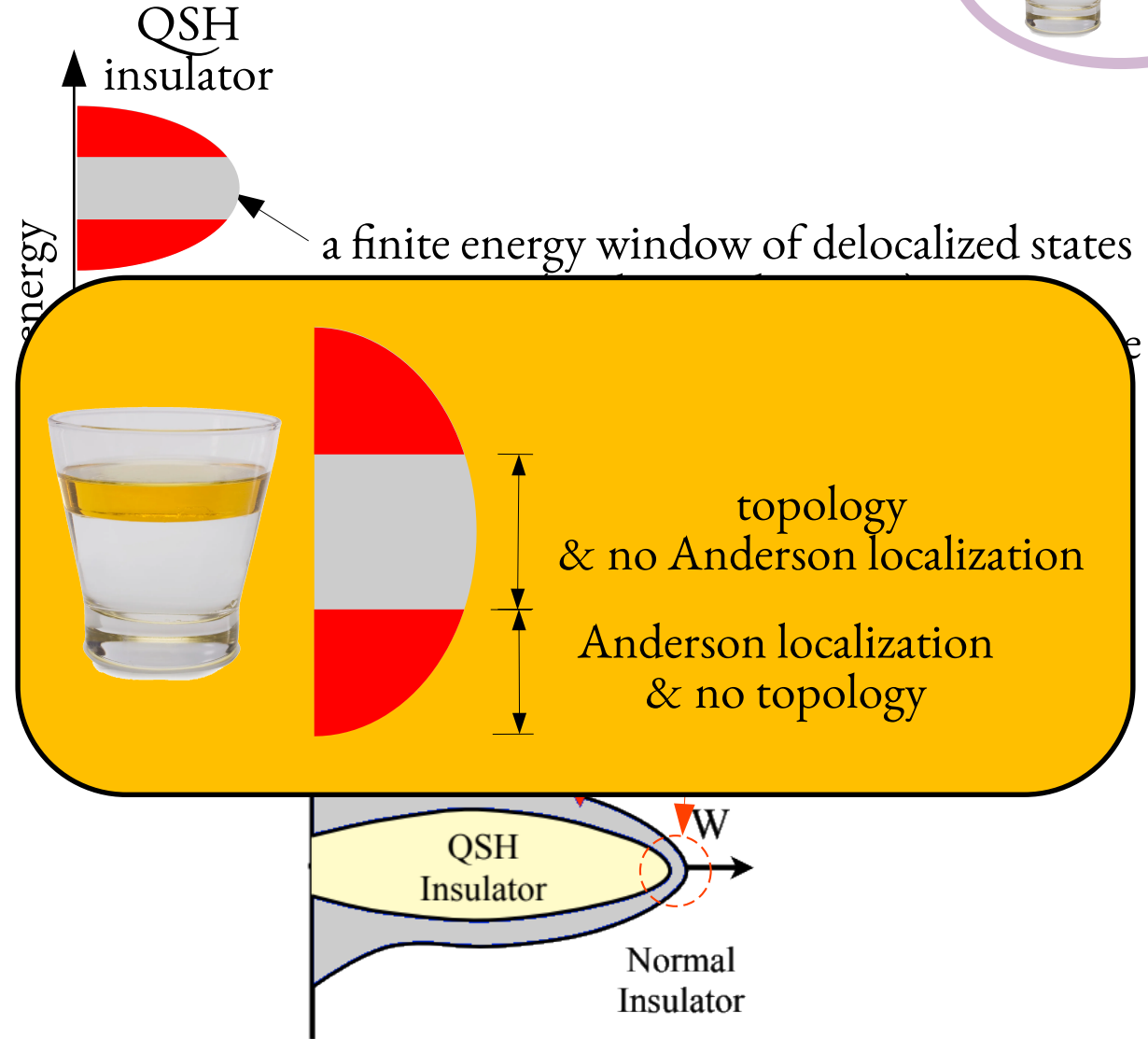


Quantum Hall/Chern insulator



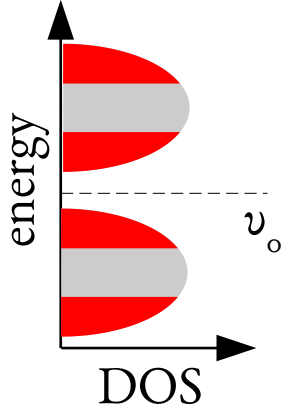
"Topological Anderson Insulator"

Quantum Spin Hall insulator



Tenfold-way

class	\mathcal{T}	\mathcal{P}	\mathcal{C}	$d=0$	$d=1$	$d=2$	$d=3$
A	0	0	0	\mathbb{Z}	0	\mathbb{Z}	0
AIII	0	0	1	0	\mathbb{Z}	0	\mathbb{Z}
AI	+	0	0	\mathbb{Z}	0	0	0
BDI	+	-	1	\mathbb{Z}_2	\mathbb{Z}	0	0



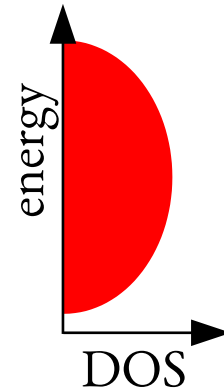
Applies to insulators, i.e., $\sigma_{xx} = 0$
(for fixed filling $\nu = \nu_0$)

Phase transition: $\sigma_{xx} \neq 0$
(for fixed filling $\nu = \nu_0$)

Topologically localized insulators



class	\mathcal{T}	\mathcal{P}	\mathcal{C}	$d=0$	$d=1$	$d=2$	$d=3$
A	0	0	0	?	?	?	\mathbb{Z}
AIII	0	0	1	?	?	?	?
AI	+	0	0	?	?	?	?
BDI	+	-	1	?	?	?	?



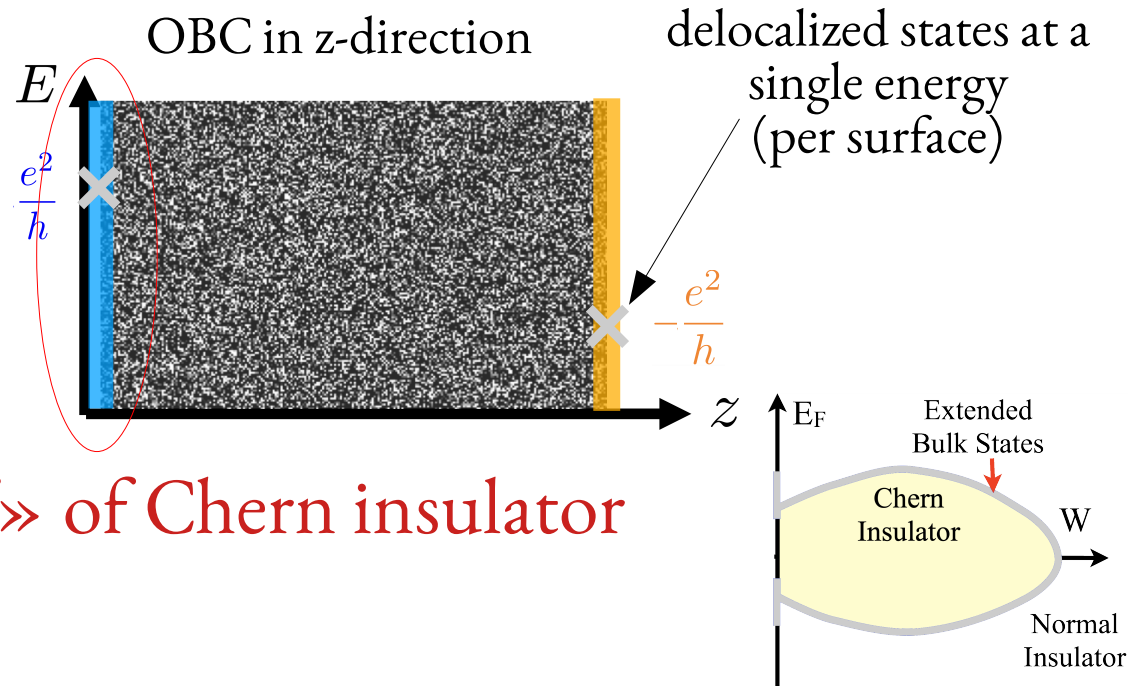
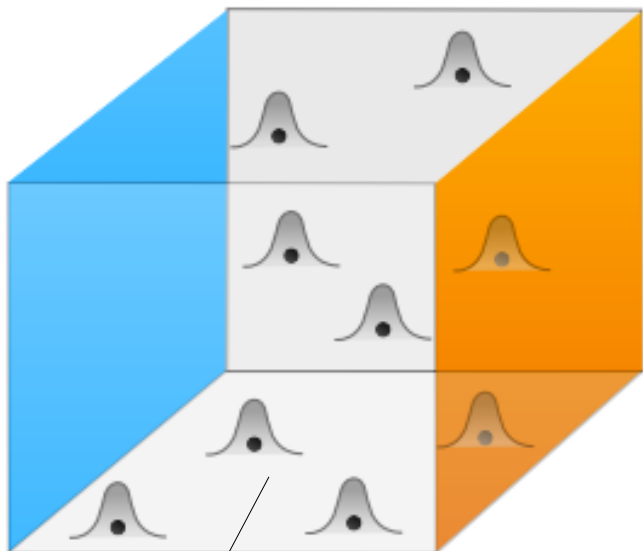
Anderson localization
& topology



Applies to fully localized insulators,
i.e., $\sigma_{xx} = 0$ (for **all** ν)

Phase transition: $\sigma_{xx} \neq 0$ (for **any** ν)

Topologically localized insulators



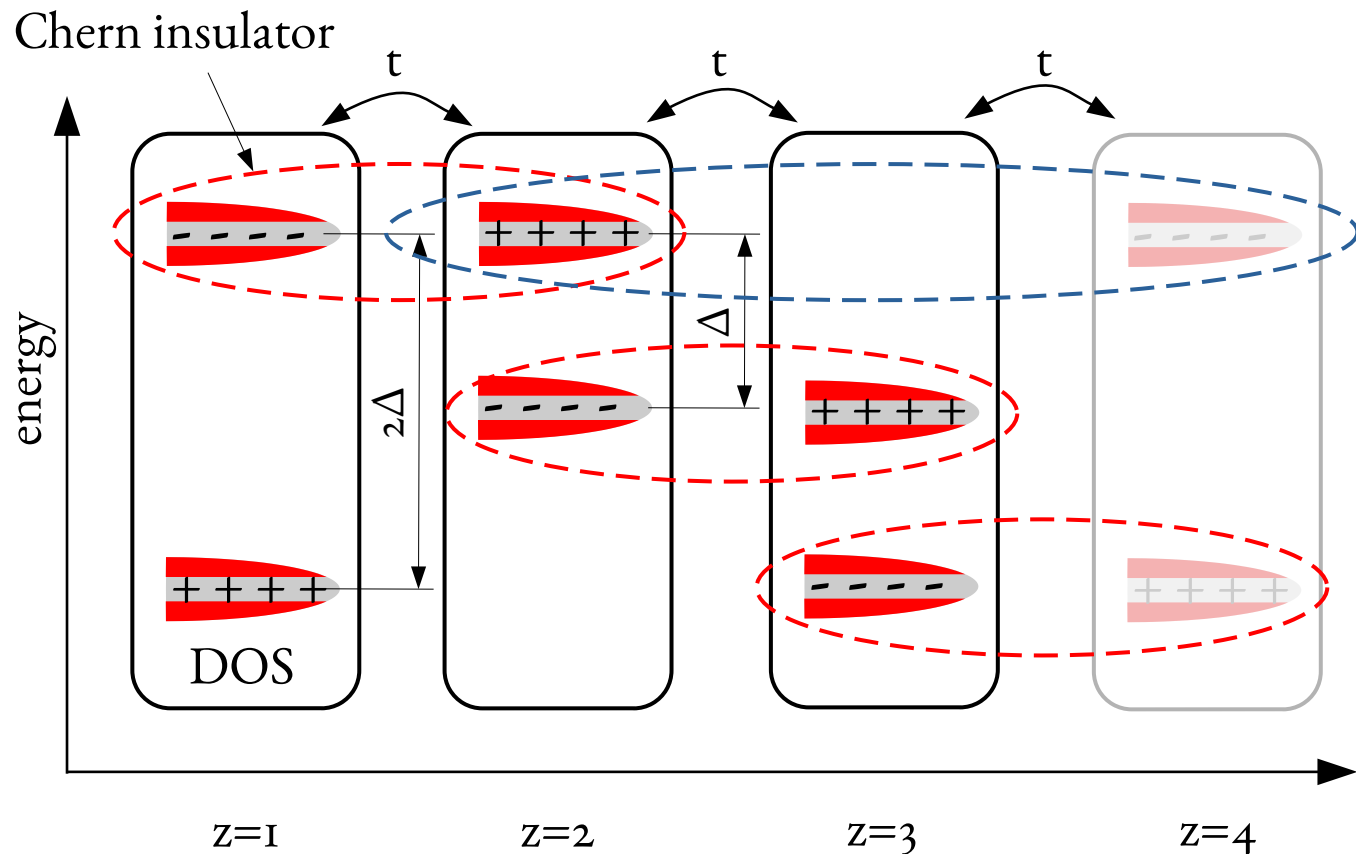
«half» of Chern insulator

⇒ no full Anderson localization for an arbitrary W !

$$\vec{P} = \alpha_{ME} \vec{B}$$

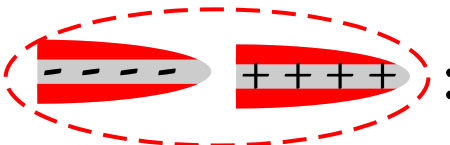
Bulk α_{ME} quantized to an integer value

	Tenfold way	TLI
Bulk	Not fully localized	Fully localized
Boundary	Delocalized	Not fully localized



Interlayer coupling t :

- on-resonant coupling 
- off-resonant coupling 

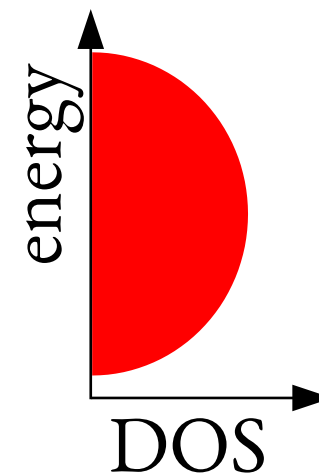
«dimers»  :

- two-dimensional
- zero net Hall conductance $\sigma_{xy} = 0$

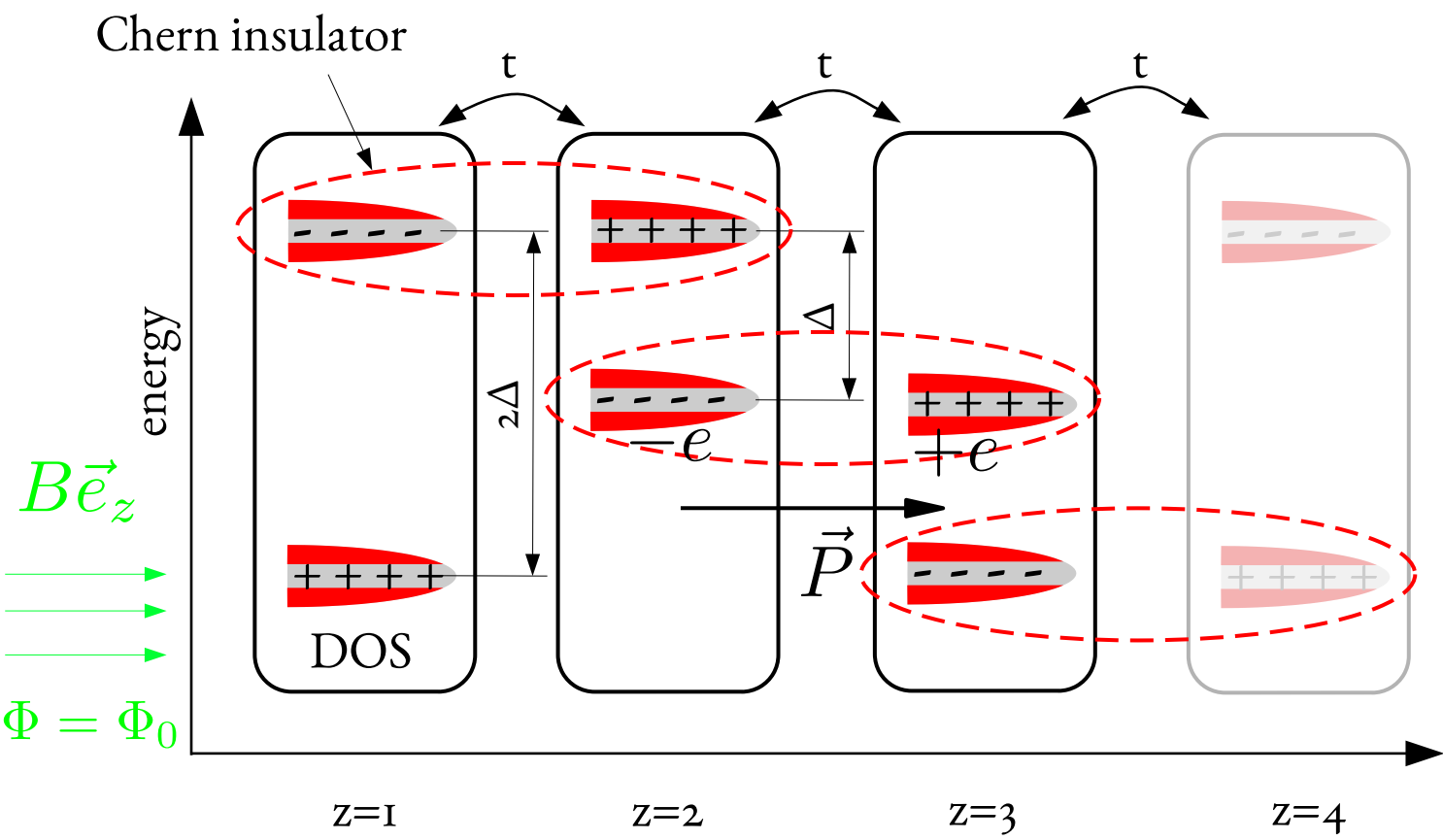
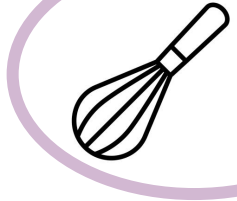
pair-annihilation



1)



Bulk is fully localized for arbitrarily weak disorder W !

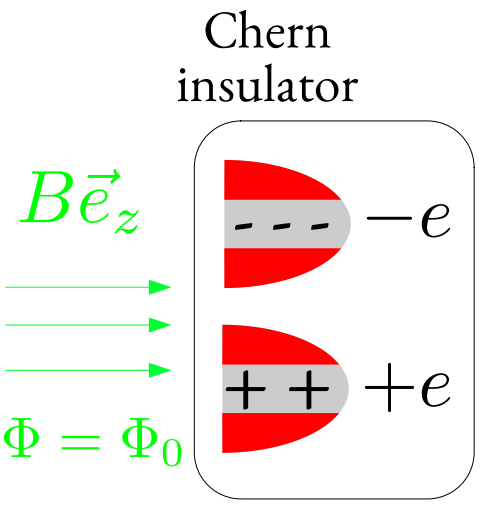


Interlayer coupling t:

- on-resonant coupling
- off-resonant coupling

2) Quantized bulk magnetoelectric polarizability!

$$\vec{P} = \hat{\alpha}_{ME} \vec{B}$$

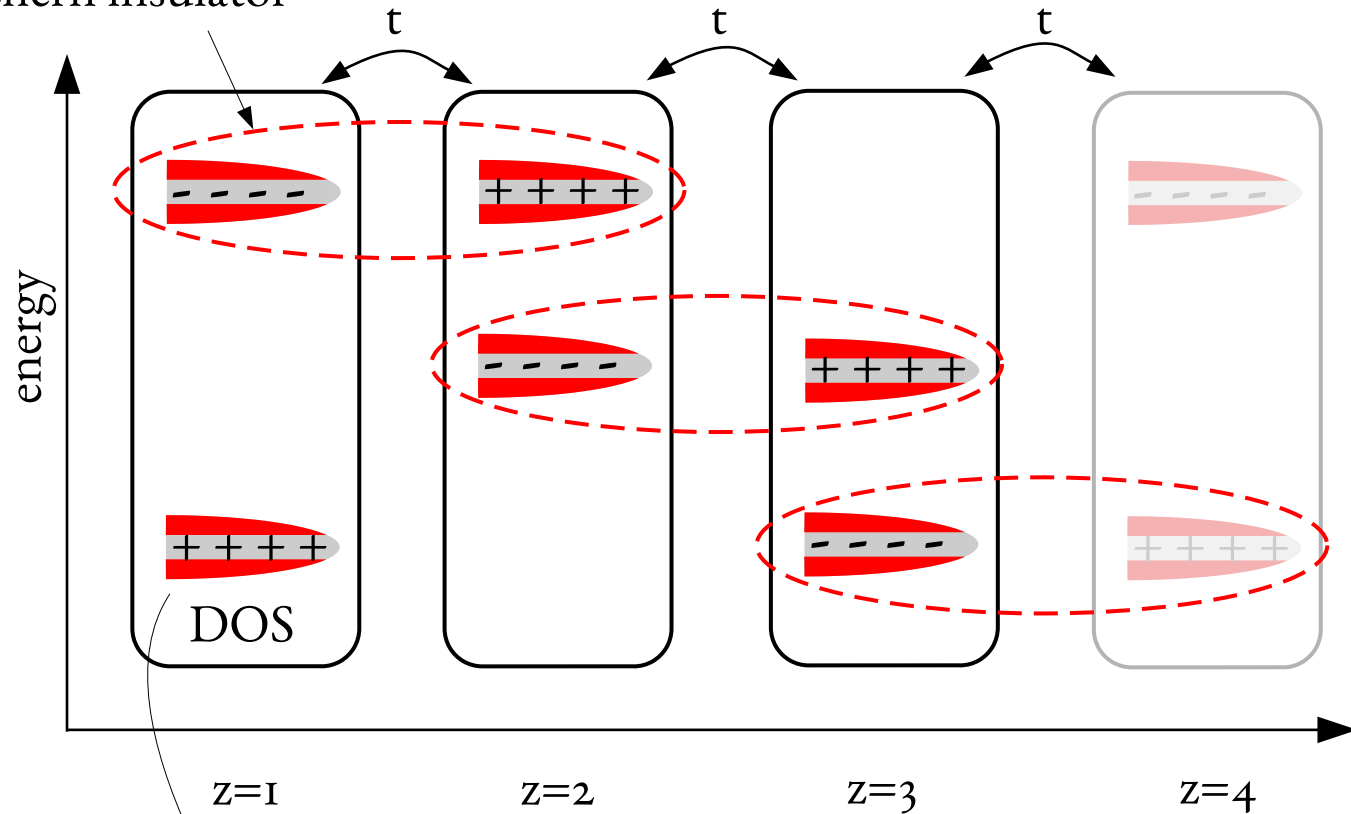


Streda formula:

$$\sigma_{xy} = -e \left(\frac{\partial n}{\partial B} \right)_{\mu} \quad \text{9}$$

$$\hat{\alpha}_{ME} = \begin{pmatrix} \alpha_{xx} & \alpha_{xy} & 0 \\ \alpha_{yx} & \alpha_{yy} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Chern insulator

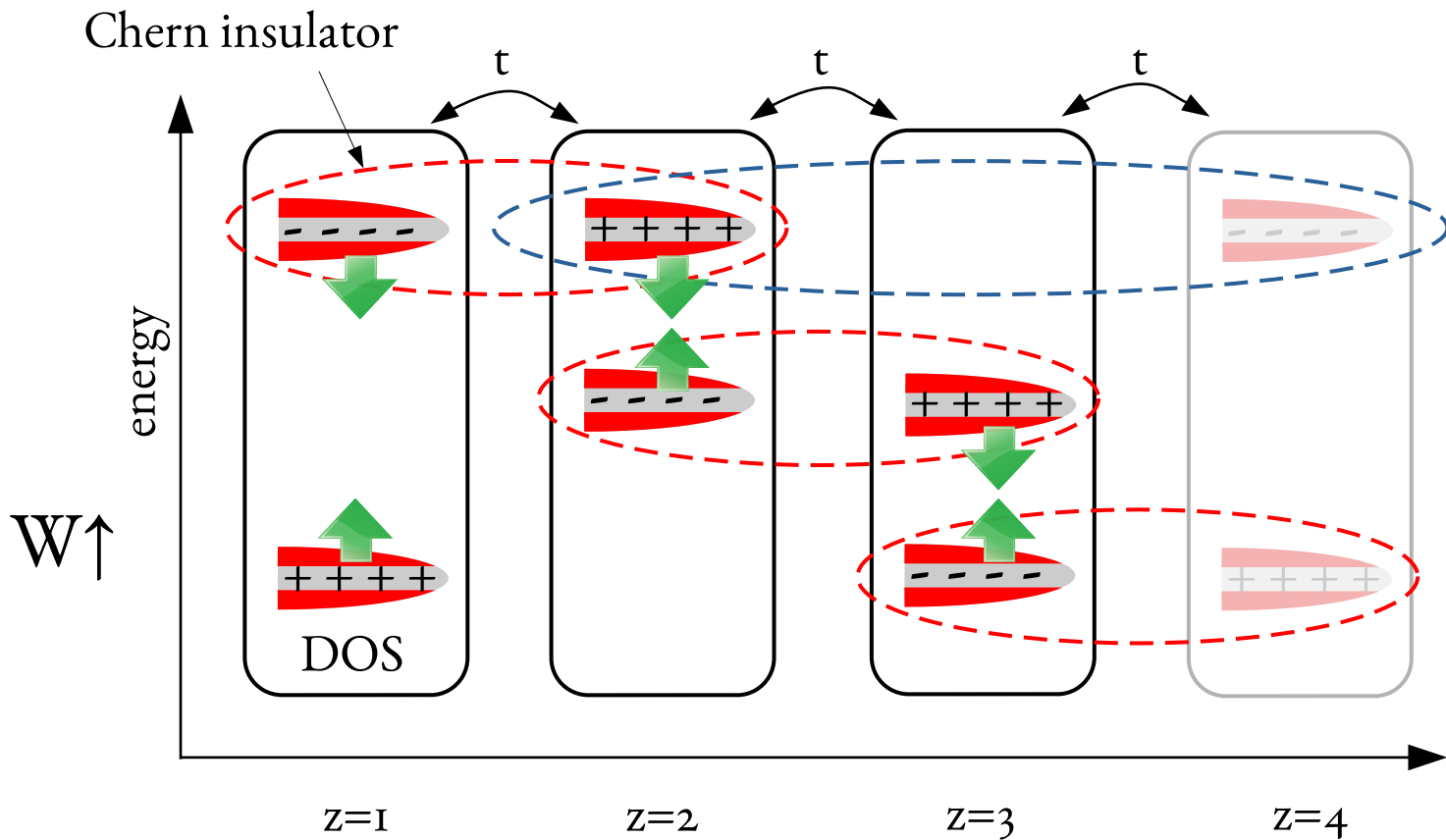


Interlayer coupling t :

- on-resonant coupling 
- off-resonant coupling 

Remove bulk!

3) Delocalized states at a single energy with quantized surface Hall conductance: $\sigma_{xy}^{\text{surf}} = e^2/h$



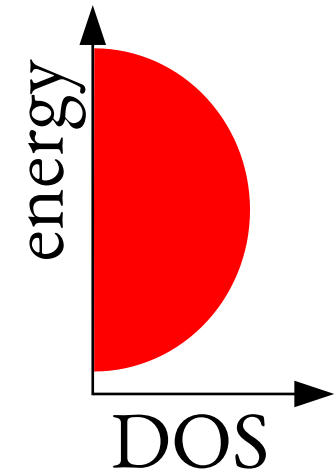
Interlayer coupling t :

- on-resonant coupling
- off-resonant coupling



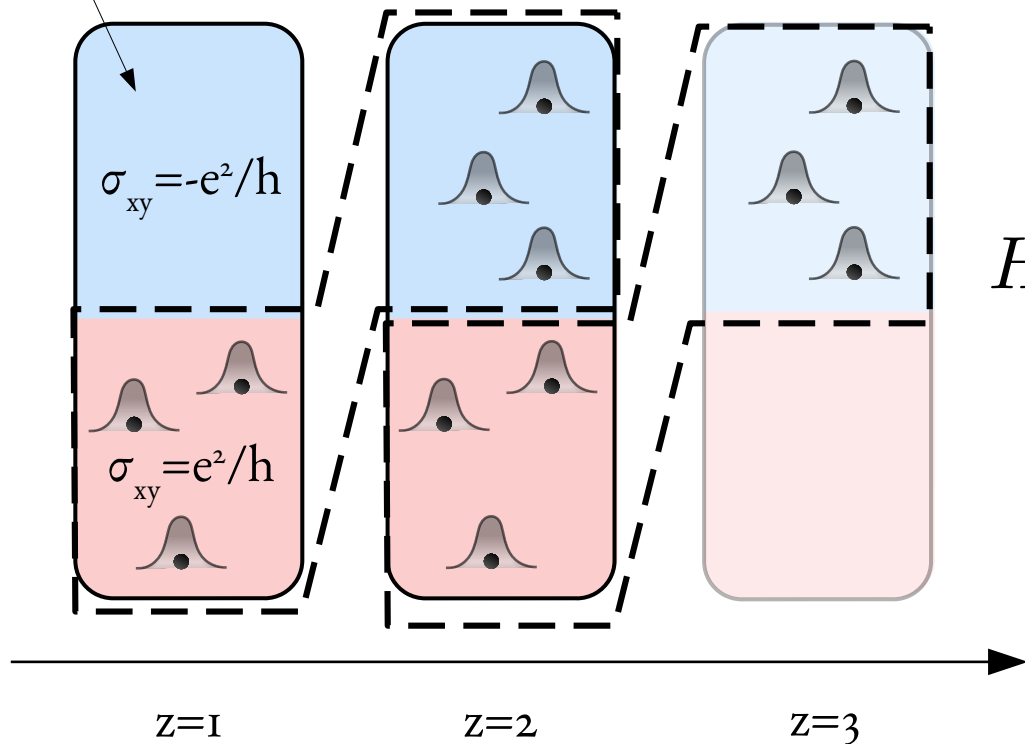
Numerically difficult problem


I)

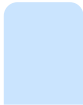


Bulk is fully localized only for finite disorder W !

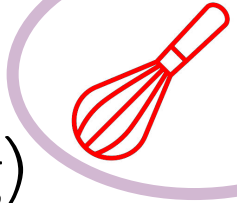
Hilbert space of
2d lattice



P_z^+ projector onto 

P_z^- projector onto 

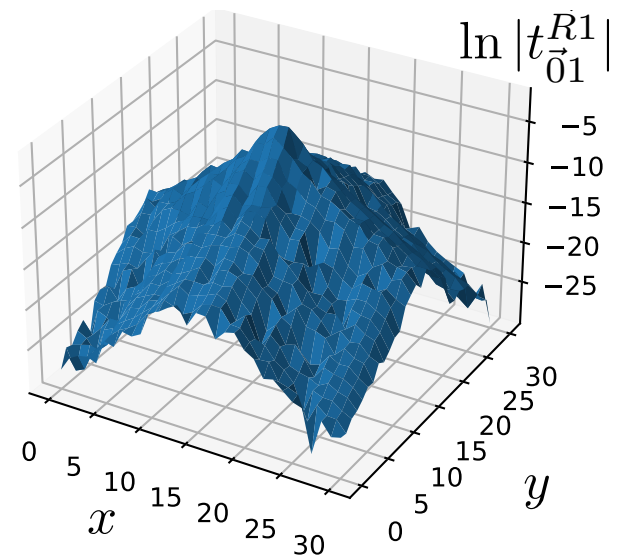
disorder
(onsite + interlayer hopping)



$$H = \sum_z (P_z^+ + P_{z+1}^-) H_W (P_z^+ + P_{z+1}^-)$$

$$H_W = \sum_{\vec{R}, \alpha} W_{\vec{R}\alpha} |g_{\vec{R}\alpha}\rangle \langle g_{\vec{R}\alpha}|$$

with $|g_{\vec{R}\alpha}\rangle = |\vec{R}\alpha\rangle + |(\vec{R} + \hat{e}_z)\alpha\rangle$



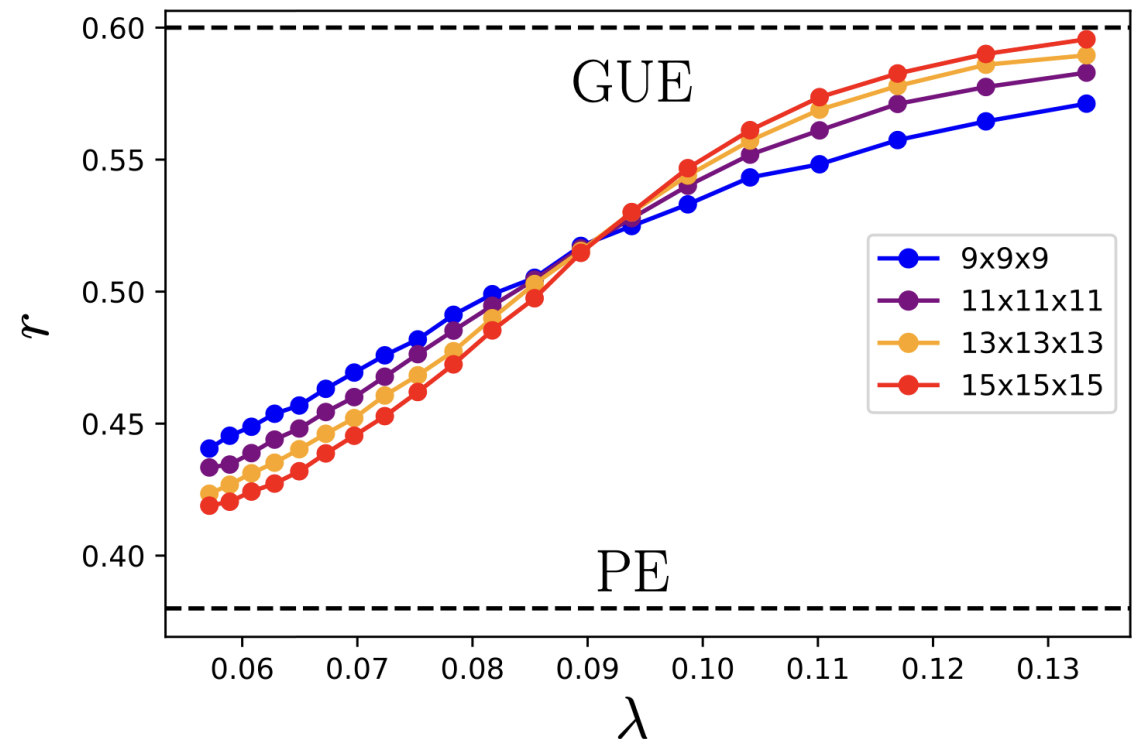
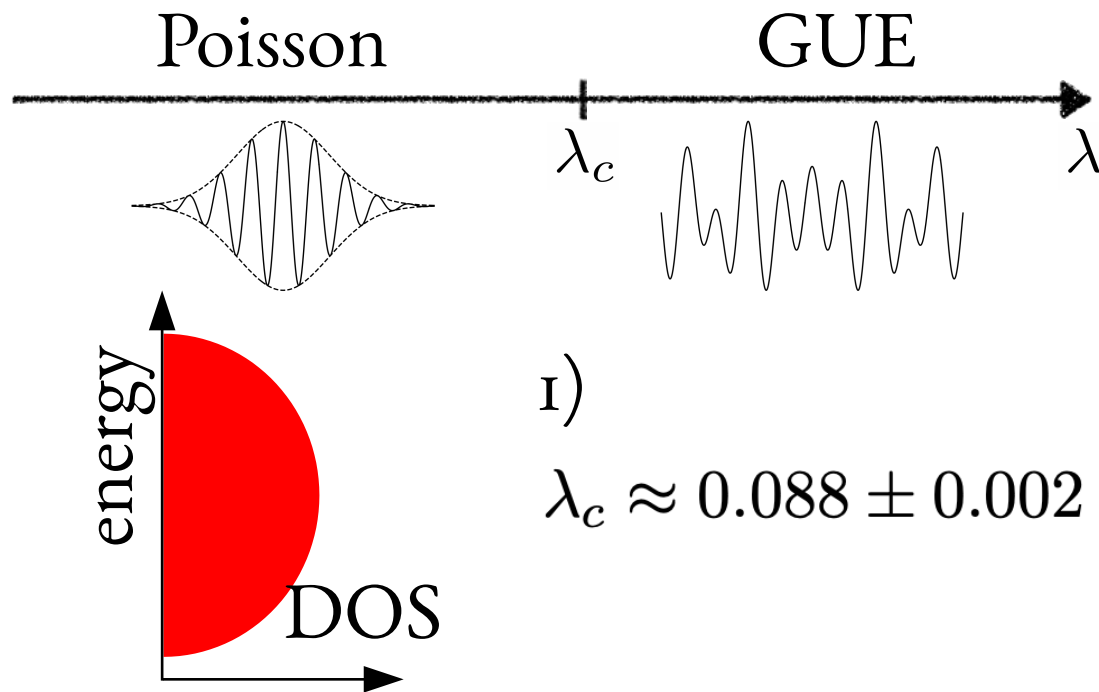
TLI to metal transition: $H_V(\lambda) = H + \lambda V$ n.n. hopping



Level spacing statistics (localization/delocalization):

$$s_n = E_{n+1} - E_n \longrightarrow \text{Level spacing ratio}$$

$$r_n = \min\{s_n, s_{n+1}\} / \max\{s_n, s_{n+1}\}$$





Bulk & boundary topological invariants $\alpha_{ME}, \sigma_{xy}^{\text{surf}}$

Fully localized phase: eigenstates $|\psi_n\rangle$ can be labeled by lattice vectors $|\psi_{\mathbf{R}\alpha}\rangle$

Unitary: $U|\mathbf{R}\alpha\rangle = |\psi_{\mathbf{R}\alpha}\rangle$

2) bulk topological invariant is the third winding number of U , $\alpha_{ME} = \nu[U]$

$$\nu[U] = \frac{i\pi}{3} \frac{1}{N_x N_y N_z} \epsilon^{ijk} \text{tr} \left(U^{-1} [\hat{X}_i, U] U^{-1} [\hat{X}_j, U] U^{-1} [\hat{X}_k, U] \right) \in \mathbb{Z} \quad \text{if } \langle \vec{R}'\alpha' | U | \vec{R}\alpha \rangle \sim e^{-\gamma|\vec{R}-\vec{R}'|}$$

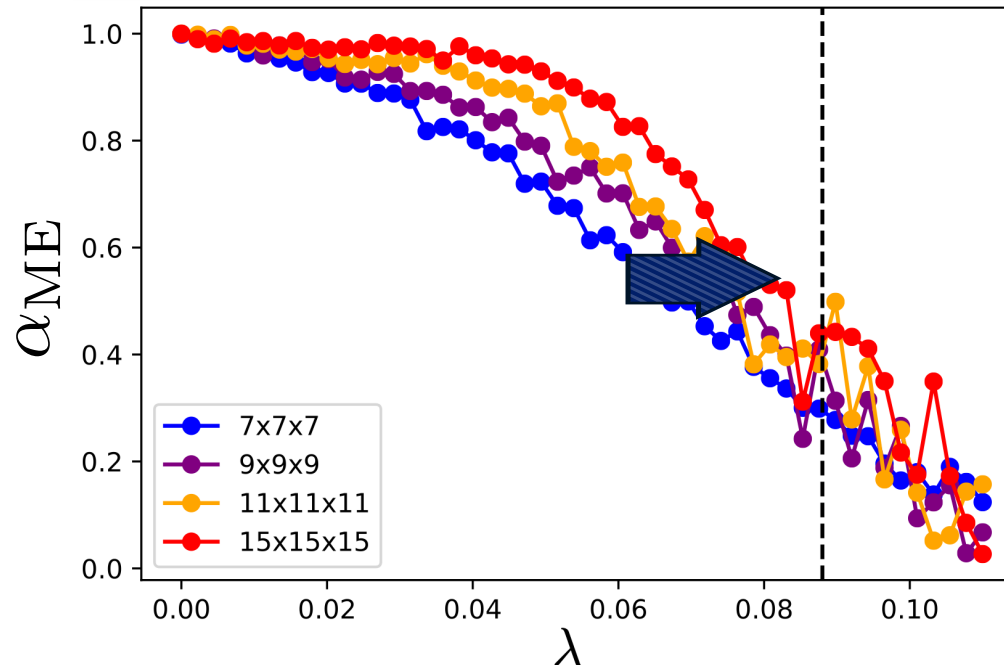
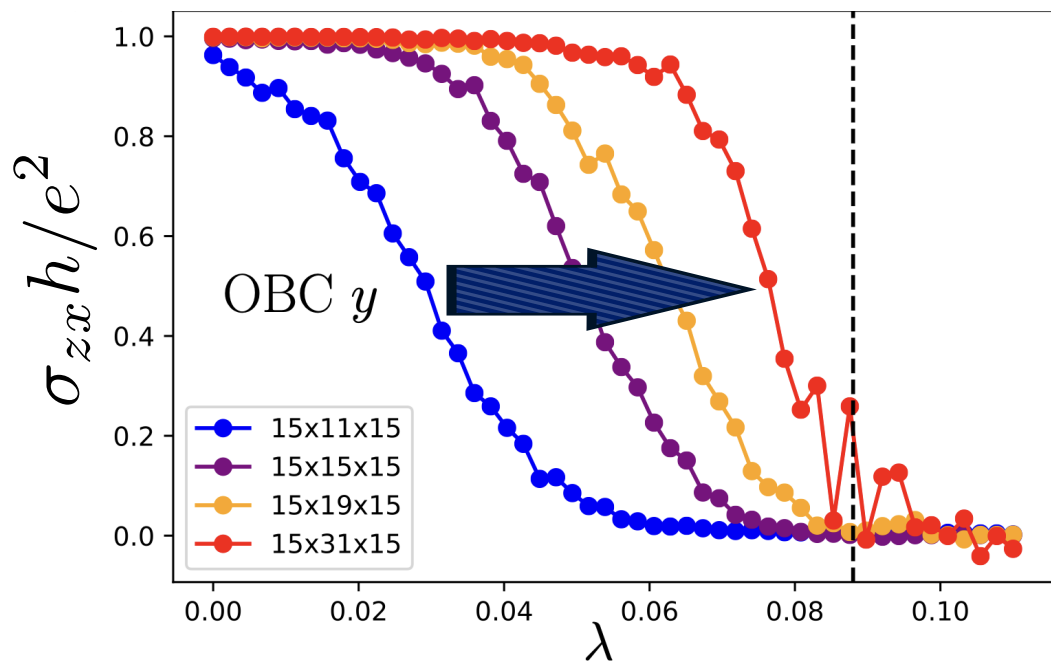
✓ if fully localized

3) boundary topological invariant is given by Chern number, $\sigma_{xy}^{\text{surf}} = \text{Ch}[P] e^2/h$

$$\text{Ch}[P] = \frac{2\pi i}{N_x N_y} \text{Tr} \left(P \left[[\hat{X}_1, P], [\hat{X}_2, P] \right] \right)$$

P projector onto d.o.f. on the surface

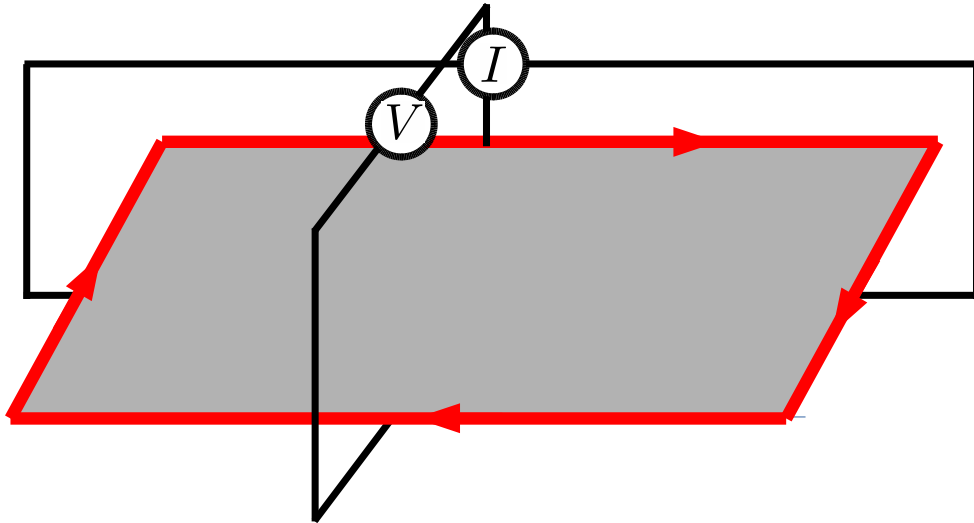
TLI to metal transition: $H_V(\lambda) = H + \lambda V$ ↘ n.n. hopping





How to measure Hall conductance?

Four-terminal measurement

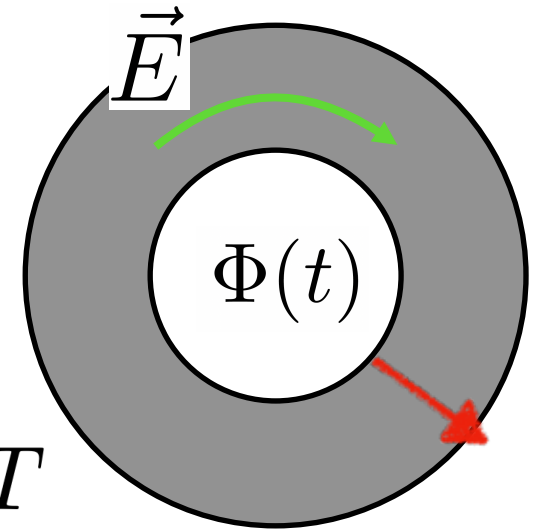


Only edges contribute to the transport

Corbino disk (Laughlin's pump)

Adiabatic
flux insertion:

$$\Phi(t) = \Phi_0 t / T$$

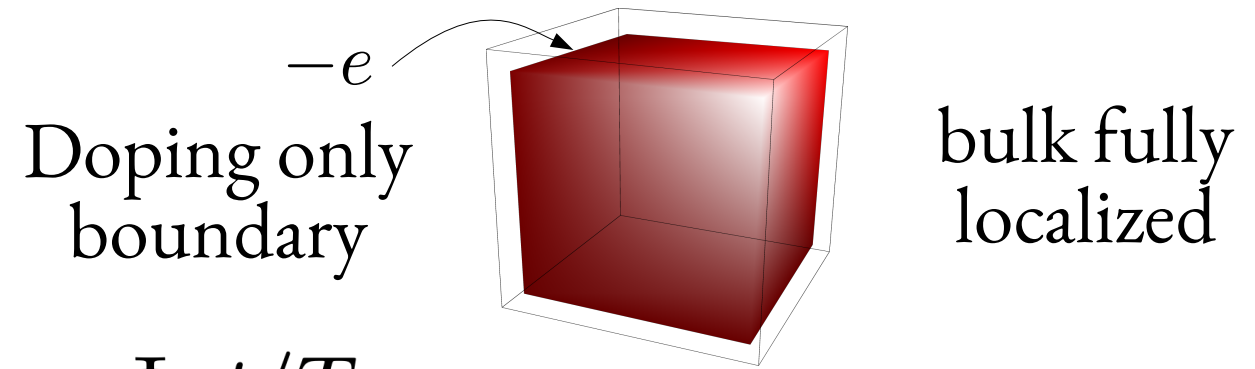


Quantized charge transfer (via bulk)

$$\Delta Q = \sigma_H \frac{h}{e} = ne$$

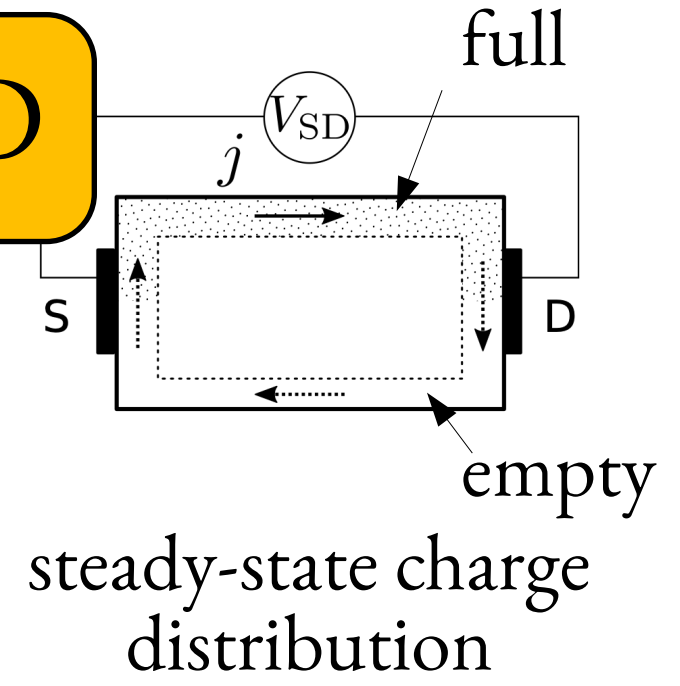
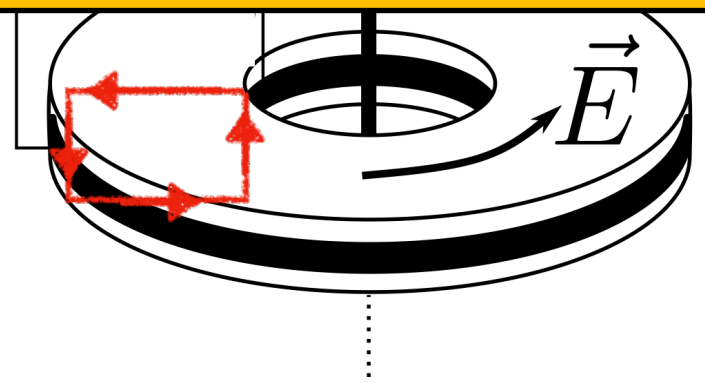


TLL: $\sigma_{xy}^{\text{surf}}$ is guaranteed to be quantized only for fully filled boundary



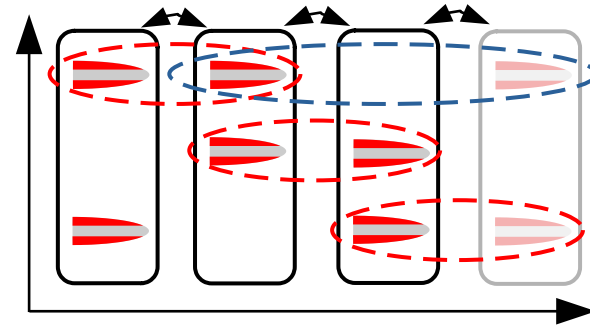
Adiabatic flux insertion $\Phi(t) = \Phi_0 t / T$

Charge transfer of $e^* \sigma_{xy}^{\text{surf}}$ from S to D



Outlook

1) Better understanding of the model



2) Field theory for TLI? Critical exponents for TLI to metal transition?

3)

class	\mathcal{T}	\mathcal{P}	\mathcal{C}	$d = 0$	$d = 1$	$d = 2$	$d = 3$
A	0	0	0	?	?	?	\mathbb{Z}
AIII	0	0	1	?	?	?	?
AI	+	0	0	?	?	?	?
BDI	+	-	1	?	?	?	?

